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q -Fractional calculus for Rubin's q -difference operator

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Abstract

In this paper we introduce a fractional q -integral operator and derivative as a generalization of Rubin's q -difference operator. We also reformulate the definition of the q^2 -Fourier transform and the q -analogue of the Fourier multiplier introduced by Rubin in (*J. Math. Anal. Appl.* 212(2):571–582, 1997; *Proc. Am. Math. Soc.* 135(3):777–785, 2007). As applications, we give summation formulas for ${}_2\phi_1$ finite series, we also use the q^2 -Fourier transform and Hahn q -Laplace transform to solve a fractional q -diffusion equation.

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1 Introduction and preliminaries

Let q be a positive number, $0 < q < 1$. In the following, we follow the notations and notations of q -hypergeometric functions, the q -gamma function $\Gamma_q(x)$, Jackson q -exponential functions $E_q(x)$, and the q -shifted factorial as in [1, 2]. The q -difference operator is defined by

$$D_q f(z) := \frac{f(z) - f(qz)}{z - qz} \quad (z \neq 0). \quad (1.1)$$

Jackson [3] introduced an integral denoted by

$$\int_a^b f(x) d_q x$$

as a right inverse of the q -derivative. It is defined by

$$\int_a^b f(t) d_q t := \int_0^b f(t) d_q t - \int_0^a f(t) d_q t \quad (a, b \in \mathbb{C}), \quad (1.2)$$

where

$$\int_0^x f(t) d_q t := (1 - q) \sum_{n=0}^{\infty} x q^n f(x q^n) \quad (x \in \mathbb{C}), \quad (1.3)$$

provided that the series on the right-hand side of (1.3) converges at $x = a$ and b .

There is no unique canonical choice for the q -integration over $[0, \infty)$. In [4], Hahn defined the q -integration for a function f over $[0, \infty)$ by

$$\int_0^\infty f(t) d_q t = (1-q) \sum_{n=-\infty}^{\infty} q^n f(q^n),$$

while in [5] Matsuo defined q -integrations on the interval $[0, \infty)$ and $(-\infty, \infty)$ by

$$\int_0^{\infty/b} f(t) d_q t := \frac{1-q}{b} \sum_{n=-\infty}^{\infty} q^n f(q^n/b) \quad (b > 0), \quad (1.4)$$

$$\int_{-\infty/b}^{\infty/b} f(t) d_q t = \frac{1-q}{b} \sum_{n=-\infty}^{\infty} q^n (f(q^n/b) + f(-q^n/b)), \quad (1.5)$$

respectively, provided that the series converges absolutely. For any $q \in (0, 1)$ and $0 < b < \infty$, we define the spaces

$$L^p(\mathbb{R}_{b,q}) := \left\{ f : \int_{-\infty/b}^{\infty/b} |f(x)|^p d_q x < \infty, p \geq 1 \right\}, \quad \mathbb{R}_{b,q} := \{ \pm q^n/b : n \in \mathbb{Z} \},$$

$$L^\infty(\mathbb{R}_{b,q}) := \{ f : \|f\|_\infty := \sup \{ f(\pm q^n/b) \mid n \in \mathbb{Z} \} < \infty \}.$$

We shall use the particular notation \mathbb{R}_q , $\widetilde{\mathbb{R}}_q$ and $\widetilde{\mathbb{R}}_{q,+}$ to denote $\mathbb{R}_{1,q}$, $\mathbb{R}_{\sqrt{1-q},q}$ and $\{ \frac{q^k}{\sqrt{1-q}}, k \in \mathbb{Z} \}$, respectively. One can verify that $L^2(\mathbb{R}_{b,q})$ associated with the inner product

$$\langle f, g \rangle := \int_{-\infty/b}^{\infty/b} f(t) \overline{g(t)} d_q t, \quad f, g \in L^2(\mathbb{R}_{b,q}),$$

is a Hilbert space. The Riemann-Liouville fractional q -integral operator is introduced by Al-Salam in [6] and later by Agarwal in [7] and defined by

$$I_q^\alpha f(x) := \frac{x^{\alpha-1}}{\Gamma_q(\alpha)} \int_0^x (qt/x; q)_{\alpha-1} f(t) d_q t, \quad \alpha \notin \{-1, -2, \dots\}. \quad (1.6)$$

Using (1.3), (1.6) reduces to

$$I_q^\alpha f(x) = x^\alpha (1-q)^\alpha \sum_{n=0}^{\infty} q^n \frac{(q^\alpha; q)_n}{(q; q)_n} f(xq^n), \quad (1.7)$$

which is valid for all α . The Riemann-Liouville fractional q -derivative of order α , $\alpha > 0$, is defined by

$$D_q^\alpha = D_q^k I_q^{k-\alpha} \quad (k = \lceil \alpha \rceil).$$

Rubin in [8, 9] introduced the q -difference operator

$$\partial_q f(z) = \frac{f(q^{-1}z) + f(-q^{-1}z) - f(qz) + f(-qz) - 2f(-z)}{2(1-q)z} \quad (z \neq 0). \quad (1.8)$$

It is straightforward to prove that if a function f is differentiable at a point z , then

$$\lim_{q \rightarrow 1^-} \partial_q f(z) = f'(z).$$

Also,

$$\delta_q f(z) = \begin{cases} D_q f(z) & \text{if } f \text{ is odd,} \\ \frac{1}{q} D_{q^{-1}} f(z) & \text{if } f \text{ is even.} \end{cases}$$

Let f and g be functions defined on a set A , where A satisfies

$$z \in A \rightarrow \pm q^{\pm 1} z \in A,$$

and let f_e and f_o be the even and odd parts of f , respectively. The following properties of the ∂_q operator are from [9, 10] and hold for all $z \in A \setminus \{0\}$.

- (i) $\partial_q f(z) = \frac{1}{q} D_{q^{-1}} f_e(z) + D_q f_o(z)$.
- (ii) For two functions f and g ,
 - if f is even and g is odd, then

$$\partial_q(fg)(z) = qg(z)(\partial_q f)(qz) + f(qz)\partial_q g(z);$$

- if f and g are even, then

$$\partial_q(fg)(z) = \partial_q f(z)g(z) + f(z/q)\partial_q g(z);$$

- if f and g are odd, then

$$\partial_q(fg)(z) = \frac{1}{q} \left(f(z)(\partial_q g)(z/q) + (\partial_q f)(z/q)g(z/q) \right).$$

The q -translation ε^y is introduced by Ismail in [2] and is defined on monomials by

$$\varepsilon^y x^n := x^n(-y/x; q)_n, \quad (1.9)$$

and it is extended to polynomials as a linear operator. Thus

$$\varepsilon^y \left(\sum_{n=0}^m f_n x^n \right) := \sum_{n=0}^m f_n \varepsilon^y x^n (-y/x; q)_n. \quad (1.10)$$

The q -translation operator is defined for x^α , $\alpha > 0$, to be

$$\varepsilon^y x^\alpha := x^\alpha (-y/x; q)_\alpha. \quad (1.11)$$

In [4], Hahn defined the following q -analogue of the Laplace transform:

$${}_q L_s f(x) = \phi(s) = \frac{1}{1-q} \int_0^{s^{-1}} E_q(-qsx) f(x) d_q x \quad (\operatorname{Re}(s) > 0). \quad (1.12)$$

Abdi [11] studied certain properties of these q -transforms. In [12], he used these analogues to solve linear q -difference equations with constant coefficients and certain allied equations. In [4, equation (9.5)], Hahn defined the convolution of two functions F, G to be

$$(F * G)(x) = \frac{x}{1-q} \int_0^1 F(tx)G[x-tqx] d_q t, \quad (1.13)$$

where $G[x-y]$, for

$$G(x) := \sum_{n=0}^{\infty} a_n x^n,$$

is defined to be

$$G[x-y] := \sum_{n=0}^{\infty} a_n [x-y]_n, \quad \text{with } [x-y]_n := x^n (y/x; q)_n.$$

Using the definition of q -integration, $(F * G)$ is nothing but

$$(F * G)(x) = \frac{1}{1-q} \int_0^x F(t) \varepsilon^{-qt} G(x) d_q t, \quad (1.14)$$

where ε is the translation operator (1.10). It is remarked by Hahn [4, p.373] that the convolution theorem

$${}_q L_s(F * G) = {}_q L_s F {}_q L_s G \quad (1.15)$$

holds. One can verify that if $\Phi(s) := {}_q L_s F(x)$ and $0 < \alpha < 1$, then

$${}_q L_s D_q^\alpha F(x) = \frac{s^\alpha}{(1-q)^\alpha} \Phi(s) - I_q^{1-\alpha} F(0^+) \frac{1}{(1-q)}; \quad (1.16)$$

see [13].

2 Orthogonality relations and completeness criteria

Koornwinder and Swarttouw introduced a q -analogue of the cosine and sine Fourier transform in [14] with the functions $\text{Cos}(z; q^2)$ and $\text{Sin}(z; q^2)$ defined by

$$\begin{aligned} \text{Cos}(z; q^2) &= \sum_{k=0}^{\infty} \frac{(-1)^k q^{k(k+1)} (z(1-q))^{2k}}{(q; q)_{2k}} \\ &= {}_1\phi_1(0; q; q^2, q^2 z^2 (1-q)^2), \\ \text{Sin}(z; q^2) &= \sum_{k=0}^{\infty} \frac{(-1)^k q^{k(k+1)} (z(1-q))^{2k+1}}{(q; q)_{2k+1}} \\ &= z {}_1\phi_1(0; q^3; q^2, q^2 z^2 (1-q)^2). \end{aligned} \quad (2.1)$$

A q -analogue of the exponential function is introduced in [8, 9] and defined by

$$e(z; q^2) := \text{Cos}(iz; q^2) - i \text{Sin}(iz; q^2).$$

Straightforward calculations give

$$\delta_q \cos(\lambda x; q^2) = -\lambda \sin(\lambda x; q^2), \quad \delta_q \sin(\lambda x; q^2) = \lambda \cos(\lambda x; q^2)$$

and

$$\delta_q e(\lambda x; q^2) = \lambda e(\lambda x; q^2),$$

where $x \in \mathbb{C}$ and λ is a fixed complex number. Fitouhi *et al.* in [15] proved that

$$\left| \frac{(z; q)_\infty}{(q; q)_\infty} {}_1\phi_1(0; z, q, q^{1+n}) \right| \leq \frac{(-|z|, -q; q)_\infty}{(q; q)_\infty} \begin{cases} 1 & \text{if } n \geq 0, \\ |z|^{-n} q^{\frac{n(n+1)}{2}} & \text{if } n < 0. \end{cases}$$

Hence,

$$\left| \sin\left(\frac{q^n}{1-q}; q^2\right) \right| \leq \frac{(-q^2; q^2)_\infty}{(q; q^2)_\infty} \begin{cases} 1 & \text{if } n \geq 0, \\ q^{n^2} & \text{if } n < 0 \end{cases} \quad (2.2)$$

and

$$\left| \cos\left(\frac{q^n}{1-q}; q^2\right) \right| \leq \frac{(-q^2; q^2)_\infty}{(q; q^2)_\infty} \begin{cases} 1 & \text{if } n \geq 0, \\ q^{n^2-2n} & \text{if } n < 0. \end{cases} \quad (2.3)$$

Consequently,

$$\left| e\left(\frac{q^n}{1-q}; q^2\right) \right| \leq 2 \frac{(-q^2; q^2)_\infty}{(q; q^2)_\infty} \begin{cases} 1, & n \geq 0, \\ q^{n^2}, & n < 0. \end{cases} \quad (2.4)$$

The following orthogonality relation is proved in [14].

Theorem 2.1 Let $|z| < 1$ and n, m be integers. Then

$$\delta_{mn} = \sum_{k=-\infty}^{\infty} z^{k+n} \frac{(z^2; q)_\infty}{(q; q)_\infty} {}_1\phi_1(0; z^2; q, q^{n+k+1}) z^{k+m} \frac{(z^2; q)_\infty}{(q; q)_\infty} {}_1\phi_1(0; z^2; q, q^{m+k+1}), \quad (2.5)$$

where the sum converges absolutely and uniformly on compact subsets of the open unit disc.

The following identity, which follows from (2.5) when we replace q by q^2 and z by q^α , $\alpha > 0$, is essential in our investigations.

$$\begin{aligned} & q^{-\alpha(n+m)} \frac{(q^2; q^2)_\infty^2}{(q^{2\alpha}; q^2)_\infty^2} \delta_{mn} \\ &= \sum_{k=-\infty}^{\infty} q^{2k\alpha} {}_1\phi_1(0; q^{2\alpha}; q^2, q^{2n+2k+2}) {}_1\phi_1(0; q^{2\alpha}; q^2, q^{2m+2k+2}). \end{aligned} \quad (2.6)$$

Theorem 2.2 For $0 < q < 1$,

$$\int_0^{\infty/\sqrt{1-q}} \text{Sin}\left(\frac{q^n z}{\sqrt{1-q}}; q^2\right) \text{Sin}\left(\frac{q^m z}{\sqrt{1-q}}; q^2\right) d_q z = \frac{\sqrt{1-q}(q^2; q^2)_\infty^2}{(q; q^2)_\infty^2} q^{-n} \delta_{n,m}, \quad (2.7)$$

$$\int_0^{\infty/\sqrt{1-q}} \text{Cos}\left(\frac{q^n z}{\sqrt{1-q}}; q^2\right) \text{Cos}\left(\frac{q^m z}{\sqrt{1-q}}; q^2\right) d_q z = \frac{\sqrt{1-q}(q^2; q^2)_\infty^2}{(q; q^2)_\infty^2} q^{-n} \delta_{n,m} \quad (2.8)$$

and

$$\int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} e\left(i \frac{q^n z}{\sqrt{1-q}}; q^2\right) e\left(-i \frac{q^m z}{\sqrt{1-q}}; q^2\right) d_q z = 4 \frac{\sqrt{1-q}(q^2; q^2)_\infty^2}{(q; q^2)_\infty^2} q^{-n} \delta_{n,m}. \quad (2.9)$$

Proof We start with proving (2.7). Since

$$\begin{aligned} & \int_0^{\infty/\sqrt{1-q}} \text{Sin}\left(\frac{q^n z}{\sqrt{1-q}}; q^2\right) \text{Sin}\left(\frac{q^m z}{\sqrt{1-q}}; q^2\right) d_q z \\ &= \sum_{k=-\infty}^{\infty} q^k \sqrt{1-q} \text{Sin}\left(\frac{q^{n+k}}{1-q}; q^2\right) \text{Sin}\left(\frac{q^{m+k}}{1-q}; q^2\right) \\ &= \frac{q^{n+m}}{(1-q)^{3/2}} \sum_{-\infty}^{\infty} q^{3k} {}_1\phi_1(0; q^3, q^2; q^{2+2n+2k}) {}_1\phi_1(0; q^3, q^2; q^{2+2m+2k}). \end{aligned} \quad (2.10)$$

In (2.6), set $\alpha = 3/2$ to obtain

$$\sum_{-\infty}^{\infty} q^{3k} {}_1\phi_1(0; q^3, q^2, q^{2+2n+2k}) {}_1\phi_1(0; q^3, q^2, q^{2+2m+2k}) = q^{-3(n+m)/2} \frac{(q^2; q^2)_\infty^2}{(q^3; q^2)_\infty^2} \delta_{n,m}. \quad (2.11)$$

Combining (2.10) and (2.11) yields (2.7). The proof of (2.8) follows similarly and the proof of (2.9) follows by combining (2.7) and (2.8). \square

Theorem 2.3 For any $q \in (0, 1)$,

- (a) the set $\{e(\pm \frac{q^n}{\sqrt{1-q}} x; q^2), n \in \mathbb{Z}\}$ is a complete orthogonal set in $L_q^2(\widetilde{\mathbb{R}}_q)$,
- (b) both of the sets $\{\text{Cos}(\frac{xq^n}{\sqrt{1-q}}; q^2), n \in \mathbb{Z}\}$ and $\{\text{Sin}(\frac{xq^n}{\sqrt{1-q}}; q^2), n \in \mathbb{Z}\}$ are complete orthogonal sets in $L_q^2(\widetilde{\mathbb{R}}_{q,+})$.

Proof We only prove (a). The proof of (b) is similar and is omitted. From Theorem 2.2, it remains only to prove that the set $\{e(\pm \frac{q^n}{\sqrt{1-q}} x; q^2), n \in \mathbb{Z}\}$ is complete in $L_q^2(\widetilde{\mathbb{R}}_q)$. This is equivalent to proving that if there exists a function $f \in L^2(\widetilde{\mathbb{R}}_q)$ such that

$$\left\langle f, e\left(\pm \frac{q^n}{\sqrt{1-q}} x; q^2\right) \right\rangle = 0 \quad (n \in \mathbb{Z}), \quad (2.12)$$

then

$$f\left(\pm \frac{q^n}{1-q} x\right) = 0 \quad (n \in \mathbb{Z}).$$

From (2.12) we deduce

$$\int_0^{\infty/\sqrt{1-q}} f_e(t) \cos\left(\frac{q^n}{\sqrt{1-q}}t; q^2\right) d_q t = 0 \quad (n \in \mathbb{Z}),$$

$$\int_0^{\infty/\sqrt{1-q}} f_o(t) \sin\left(\frac{q^n}{\sqrt{1-q}}t; q^2\right) d_q t = 0 \quad (n \in \mathbb{Z}),$$

where f_e and f_o are the even and odd parts of the function f . Then from (3.8)-(3.9) we obtain $f_e(t) = f_o(t) = f(t) = 0$ for all $t \in \widetilde{\mathbb{R}}_q$. Hence $\{e(\pm \frac{q^n}{\sqrt{1-q}}x; q^2), n \in \mathbb{Z}\}$ is a complete orthogonal set in $L^2(\widetilde{\mathbb{R}}_q)$. \square

Theorem 2.4

(1) If $f \in L^2(\widetilde{\mathbb{R}}_q)$, then

$$f(x) = \sum_{-\infty}^{\infty} c_n e\left(ix \frac{q^n}{\sqrt{1-q}}; q^2\right) + \sum_{-\infty}^{\infty} d_n e\left(-ix \frac{q^n}{\sqrt{1-q}}; q^2\right) \quad (x \in \widetilde{\mathbb{R}}_q), \quad (2.13)$$

where

$$c_n = \frac{q^n}{C} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) e\left(-it \frac{q^n}{\sqrt{1-q}}; q^2\right) d_q t \quad (n \in \mathbb{Z}),$$

$$d_n = \frac{q^n}{C} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) e\left(it \frac{q^n}{\sqrt{1-q}}; q^2\right) d_q t \quad (n \in \mathbb{Z}),$$

and $C := \frac{4\sqrt{1-q}(q^2;q^2)_\infty^2}{(q;q^2)_\infty^2}$.

(2) If $f \in L^2(\widetilde{\mathbb{R}}_q)$, then

$$f(x) = \sum_{n=-\infty}^{\infty} a_n \cos\left(\frac{xq^n}{\sqrt{1-q}}; q^2\right) + \sum_{n=-\infty}^{\infty} b_n \sin\left(\frac{xq^n}{\sqrt{1-q}}; q^2\right) \quad (x \in \widetilde{\mathbb{R}}_q), \quad (2.14)$$

where

$$a_n = \frac{4}{C} \int_0^{\infty/\sqrt{1-q}} f_e(t) \cos\left(\frac{tq^n}{\sqrt{1-q}}; q^2\right) d_q t \quad (n \in \mathbb{Z})$$

and

$$b_n = \frac{4}{C} \int_0^{\infty/\sqrt{1-q}} f_o(t) \sin\left(\frac{tq^n}{\sqrt{1-q}}; q^2\right) d_q t \quad (n \in \mathbb{Z}).$$

Proof The proof of (1) follows directly from Theorem 2.3 and the orthogonality relations (2.9). In the following we give in detail the proof of (2). Let $f = f_e + f_o$ be any function in $L^2(\widetilde{\mathbb{R}}_q)$. Clearly both f_e and f_o belong to $L^2(\widetilde{\mathbb{R}}_q)$. The restriction of f_e to $\widetilde{\mathbb{R}}_{q,+}$ can be represented in the complete orthogonal set $\{\cos(\frac{xq^n}{\sqrt{1-q}}), n \in \mathbb{Z}\}$ as

$$f_e(x) = \sum_{n=-\infty}^{\infty} a_n \cos\left(\frac{xq^n}{\sqrt{1-q}}; q^2\right) \quad (x \in \widetilde{\mathbb{R}}_{q,+}), \quad (2.15)$$

where

$$a_n = \frac{4}{C} \int_0^{\infty/\sqrt{1-q}} f_e(t) \cos\left(\frac{tq^n}{\sqrt{1-q}}; q^2\right) d_q t \quad (n \in \mathbb{Z}).$$

The orthogonal set $\{\sin(\frac{xq^n}{\sqrt{1-q}}), n \in \mathbb{Z}\}$ also spans $L^2(\tilde{\mathbb{R}}_{q,+})$, hence

$$f_o(x) = \sum_{n=-\infty}^{\infty} b_n \sin\left(\frac{xq^n}{\sqrt{1-q}}; q^2\right) \quad (x \in \tilde{\mathbb{R}}_{q,+}), \quad (2.16)$$

where

$$b_n = \frac{4}{C} \int_0^{\infty/\sqrt{1-q}} f_o(t) \sin\left(\frac{tq^n}{\sqrt{1-q}}; q^2\right) d_q t \quad (n \in \mathbb{Z}).$$

Because both sides of (2.15) are even functions on $\tilde{\mathbb{R}}_q$, the equality extends on $\tilde{\mathbb{R}}_q$; and similarly the two sides of (2.16). Hence we have the representation (2.14) of any $f \in L^2(\tilde{\mathbb{R}}_q)$. \square

3 Rubin's q^2 -Fourier transform

Koornwinder and Swarttouw [14] introduced the pair of q -transforms

$$\begin{aligned} g(q^n) &= \frac{(q;q^2)_\infty}{(q^2;q^2)_\infty} \sum_{k=-\infty}^{\infty} q^k \begin{cases} \cos\left(\frac{q^{k+n}}{1-q}; q^2\right) \\ \text{or} \\ \sin\left(\frac{q^{k+n}}{1-q}; q^2\right) \end{cases} f(q^k), \\ f(q^k) &= \frac{(q;q^2)_\infty}{(q^2;q^2)_\infty} \sum_{n=-\infty}^{\infty} q^n \begin{cases} \cos\left(\frac{q^{k+n}}{1-q}; q^2\right) \\ \text{or} \\ \sin\left(\frac{q^{k+n}}{1-q}; q^2\right) \end{cases} g(q^n), \end{aligned} \quad (3.1)$$

where $0 < q < 1$ and f, g are in the space $L^2(\mathbb{R}_q)$. Now assume that $\frac{\log(1-q)}{\log q} \in 2\mathbb{Z}$ or, equivalently,

$$q \in \{q \in (0, 1) : 1 - q = q^{2m} \text{ for some integer } m\}. \quad (3.2)$$

Then, by replacing q^k and q^n in (3.1) by $q^k\sqrt{1-q}$ and $q^n\sqrt{1-q}$, and then $f(q^k\sqrt{1-q})$ and $g(q^n\sqrt{1-q})$ by $f(q^k)$ and $g(q^n)$, Koornwinder and Swarttouw obtained the following q -analogue of the cosine and sine Fourier transforms:

$$\begin{aligned} g(\lambda) &= \frac{\sqrt{1+q}}{\Gamma_q^2(1/2)} \int_0^\infty f(t) \begin{cases} \cos(t\lambda; q^2) \\ \text{or} \\ \sin(t\lambda; q^2) \end{cases} d_q t, \\ f(x) &= \frac{\sqrt{1+q}}{\Gamma_q^2(1/2)} \int_0^\infty f(t) \begin{cases} \cos(x\lambda; q^2) \\ \text{or} \\ \sin(x\lambda; q^2) \end{cases} d_q t. \end{aligned} \quad (3.3)$$

Therefore, if we let $q \rightarrow 1^-$ for such q 's that satisfy (3.2), we obtain the cosine and sine Fourier transforms

$$g(\lambda) = \sqrt{\frac{2}{\pi}} \int_0^\infty f(t) \cos(\lambda t) dt, \quad f(t) = \sqrt{\frac{2}{\pi}} \int_0^\infty g(\lambda) \cos(t\lambda) d\lambda, \quad (3.4)$$

$$g(\lambda) = \sqrt{\frac{2}{\pi}} \int_0^\infty f(t) \sin(\lambda t) dt, \quad f(t) = \sqrt{\frac{2}{\pi}} \int_0^\infty g(\lambda) \sin(t\lambda) d\lambda. \quad (3.5)$$

The pair of functions $\text{Cos}(\lambda x; q^2)$ and $\text{Sin}(\lambda x; q^2)$ satisfy

$$-\frac{1}{q} D_{q^{-1}} D_q y(x) = \begin{cases} \lambda^2 y(x) & \text{if } y(x) = \text{Sin}(\lambda x; q^2), \\ q\lambda^2 y(x) & \text{if } y(x) = \text{Cos}(\lambda x; q^2). \end{cases}$$

Therefore, the eigenfunctions $\{\text{Cos}(\lambda x; q^2), \text{Sin}(\lambda x; q^2)\}$ have two different eigenvalues. Consequently, as remarked by Koornwinder and Swarttouw in [14], no q -exponential functions built from $\{\text{Cos}(x; q^2), \text{Sin}(x; q^2)\}$ will satisfy an eigenfunction problem. This motivated Rubin [8] to define the q -difference operator (1.8) since for this operator, the functions $\{\text{Cos}(\lambda x; q^2), \text{Sin}(\lambda x; q^2)\}$ are solutions of the eigenvalue problem

$$-\delta_q^2 y(x) = \lambda^2 y(x).$$

Rubin [8] introduced a q^2 -analogue of the Fourier transform in the form

$$\hat{f}(x; q^2) := \mathcal{F}_q(f)(x) = \frac{\sqrt{1+q}}{2\Gamma_{q^2}(1/2)} \int_{-\infty}^\infty f(t) e(-itx; q^2) d_q t, \quad (3.6)$$

where $f \in L^1(\mathbb{R}_q)$ and q satisfies condition (3.2).

Remark 3.1 Rubin [9] proved that

- (1) the q^2 -Fourier transform defines a bounded linear operator from $L^1(\mathbb{R}_q)$ to $L^\infty(\mathbb{R}_q)$,
- (2) the q^2 -Fourier transform is defined and bounded on $L^1(\mathbb{R}_q) \cap L^2(\mathbb{R}_q)$,
- (3) $L^1(\mathbb{R}_q) \cap L^2(\mathbb{R}_q)$ is dense in $L^2(\mathbb{R}_q)$ (consider the functions with finite support).

Consequently, the q^2 -Fourier transform defines a bounded extension to $L^2(\mathbb{R}_q)$.

Koornwinder and Swarttouw introduced the q -Hankel transforms (3.1) which can be written in the form (3.3) only if q satisfies condition (3.2). In fact, we can write the q -transforms in (3.1) as q -integral on $(-\infty, \infty)$ by using Matsuo definition (1.4) as in the following. Rewrite the transform pair in (3.1) as

$$g\left(\frac{q^n}{\sqrt{1-q}}\right) = \frac{(q; q^2)_\infty}{(q^2; q^2)_\infty} \sum_{k=-\infty}^\infty q^k \begin{cases} \text{Cos}\left(\frac{q^{k+n}}{1-q}; q^2\right) \\ \text{or} \\ \text{Sin}\left(\frac{q^{k+n}}{1-q}; q^2\right) \end{cases} f\left(\frac{q^k}{\sqrt{1-q}}\right), \quad (3.7)$$

$$f\left(\frac{q^k}{\sqrt{1-q}}\right) = \frac{(q; q^2)_\infty}{(q^2; q^2)_\infty} \sum_{n=-\infty}^\infty q^n \begin{cases} \text{Cos}\left(\frac{q^{k+n}}{1-q}; q^2\right) \\ \text{or} \\ \text{Sin}\left(\frac{q^{k+n}}{1-q}; q^2\right) \end{cases} g\left(\frac{q^n}{\sqrt{1-q}}\right),$$

where we assume that the functions f and g are in the space $L^1(\widetilde{\mathbb{R}}_q) \cap L^2(\widetilde{\mathbb{R}}_q)$. Using Matsuo definition of the q -integration on $(0, \infty)$, (1.4) with $b = \sqrt{1-q}$, the transformations in (3.7) can be written as

$$\begin{aligned} g(x) &= \frac{\sqrt{1+q}}{\Gamma_{q^2}(1/2)} \int_0^{\infty/\sqrt{1-q}} f(t) \cos(xt; q^2) d_q t, \\ f(t) &= \frac{\sqrt{1+q}}{\Gamma_{q^2}(1/2)} \int_0^{\infty/\sqrt{1-q}} g(x) \cos(xt; q^2) d_q x \end{aligned} \quad (3.8)$$

and

$$\begin{aligned} g(x) &= \frac{\sqrt{1+q}}{\Gamma_{q^2}(1/2)} \int_0^{\infty/\sqrt{1-q}} f(t) \sin(xt; q^2) d_q t, \\ f(t) &= \frac{\sqrt{1+q}}{\Gamma_{q^2}(1/2)} \int_0^{\infty/\sqrt{1-q}} g(x) \sin(xt; q^2) d_q x, \end{aligned} \quad (3.9)$$

where $x, t \in \widetilde{\mathbb{R}}_q$ and f, g are in $L^1(\widetilde{\mathbb{R}}_q) \cap L^2(\widetilde{\mathbb{R}}_q)$. This is similar to Rubin's work in [9]. Consequently, we set the following reformulation of Rubin's definition of the q^2 -Fourier transform (3.6).

Definition 3.2 Let $0 < q < 1$. We define the q^2 -Fourier transform for any function $f \in L^1(\widetilde{\mathbb{R}}_q)$ to be

$$\hat{f}(x; q^2) := \mathcal{F}_q(f)(x) = \frac{\sqrt{1+q}}{2\Gamma_{q^2}(1/2)} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) e(-itx; q^2) d_q t. \quad (3.10)$$

It is clear that Rubin's definition of the q^2 -Fourier transform is a special case of (3.2) because if $1-q = q^{2m}$ for some $m \in \mathbb{Z}$, then

$$\int_0^{\infty/\sqrt{1-q}} f(t) d_q t = \int_0^{\infty/q^m} f(t) d_q t = \int_0^{\infty} f(t) d_q t.$$

However, we get the classical Fourier transform only when $q \rightarrow 1^-$ and q satisfies (3.2). Similar to Rubin's results mentioned in Remark 3.1, we can prove that the q^2 -Fourier transform defines a bounded linear operator from $L^1(\widetilde{\mathbb{R}}_q)$ to $L^\infty(\widetilde{\mathbb{R}}_q)$, and $L^1(\widetilde{\mathbb{R}}_q) \cap L^2(\widetilde{\mathbb{R}}_q)$ is dense in $L^2(\widetilde{\mathbb{R}}_q)$. Therefore, the q^2 -Fourier transform in (3.2) defines a bounded extension to $L^2(\widetilde{\mathbb{R}}_q)$.

The proofs of the following results, which are valid for any $q \in (0, 1)$, are similar to the proofs in [8]. Therefore, we state them without proofs.

(1) If $f \in L^2(\widetilde{\mathbb{R}}_q)$, then

$$f(t) = \frac{\sqrt{1+q}}{2\Gamma_{q^2}(1/2)} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} \mathcal{F}(f)(x) e(itx(1-q); q^2) d_q x, \quad t \in \widetilde{\mathbb{R}}_q. \quad (3.11)$$

(2) If $f(u)$ and $uf(u) \in L^1(\widetilde{\mathbb{R}}_q)$, then

$$\partial_q(\mathcal{F}_q f)(x) = \mathcal{F}_q(-iuf(u))(x).$$

(3) If f and $\partial_q f \in L_q^1(\widetilde{\mathbb{R}}_q)$, then

$$\mathcal{F}_q(\partial_q f)(x) = ix\mathcal{F}_q(f)(x). \quad (3.12)$$

We reformulate the definitions of q^2 -Fourier multiplier and the q^2 -Fourier convolution formula introduced by Rubin in [9] with the restriction (3.2) to any $q \in (0, 1)$.

Definition 3.3 Let $q \in (0, 1)$. We define the q^2 -Fourier multiplier operator corresponding to translation by y to be

$$(T_y f)(x) = \frac{\sqrt{1+q}}{2\Gamma_{q^2}(1/2)} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} e(-ity; q^2) (\mathcal{F}_q f)(t) e(itx; q^2) d_q t, \quad (3.13)$$

whenever the q -integral makes sense. If $f \in L^2(\widetilde{\mathbb{R}}_q)$ and $g \in L^1(\widetilde{\mathbb{R}}_q)$, we define the multiplier corresponding to Fourier convolution of f with g to be

$$(f * g)(z) = \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} [T_\lambda f](z) g(\lambda) d_q \lambda. \quad (3.14)$$

Theorem 3.4 Let f and g be two functions in $(L^1 \cap L^2)(\widetilde{\mathbb{R}}_q)$. Then

$$\mathcal{F}_q(f * g)(x) = \mathcal{F}_q(f)(x) \mathcal{F}_q(g)(x) \quad (x \in \widetilde{\mathbb{R}}_q). \quad (3.15)$$

Proof The proof of (3.15) is completely similar to the proof of [9, Theorem 8] and is omitted. \square

4 Fractional q -operator as a generalization of a q -difference operator

Let f be an integrable function of period 2π . Weyl, see Zygmund's book [16], introduced a fractional operator which is more convenient for trigonometric series than the Riemann-Liouville fractional operator. This operator is defined by

$$(I_\alpha f)(x) \sim \sum_{n=-\infty}^{\infty} c_n \frac{e^{inx}}{(in)^\alpha} \quad \text{if } f(x) \sim \sum_{-\infty}^{\infty} c_n e^{inx}, c_0 = 0, \quad (4.1)$$

where $i^\alpha = e^{i\alpha\pi/2}$. Zygmund [16, p.133] pointed out that

$$I_\alpha f(x) = \frac{1}{2\pi} \int_0^{2\pi} f(t) \Psi_\alpha(x-t) dt, \quad \Psi_\alpha(x) = \sum_{n \neq 0} \frac{e^{inx}}{(in)^\alpha}.$$

He also proved the semigroup identity

$$I_\alpha I_\beta = I_{\alpha+\beta}, \quad \alpha, \beta > 0.$$

In [17], Ismail and Rahman defined a q -analogue of the fractional operator I_α , so that $\alpha = 1$ represents a right inverse of the Askey-Wilson operator \mathcal{D}_q which is defined by

$$(\mathcal{D}_q f)(x) := \frac{\check{f}(q^{1/2}e^{i\theta}) - \check{f}(q^{-1/2}e^{i\theta})}{(q^{1/2} - q^{-1/2}) \sin \theta}, \quad x = \cos \theta,$$

where $f(x) = \check{f}(z)$ with $x = (z + 1/z)/2$.

In this section, we introduce a q -analogue of the fractional operator (4.1) as a generalization of the q -difference operator defined by Rubin in [8]. From Theorem 2.3, consequently,

$$f(x) = \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) \Psi_0(x, t) d_q t,$$

where

$$\begin{aligned} \Psi_0(x, t) &= \sum_{-\infty}^{\infty} q^n e\left(ix \frac{q^n}{\sqrt{1-q}}\right) e\left(-it \frac{q^n}{\sqrt{1-q}}\right) + q^n e\left(-ix \frac{q^n}{\sqrt{1-q}}\right) e\left(it \frac{q^n}{\sqrt{1-q}}\right) \\ &= \sum_{-\infty}^{\infty} q^n \cos\left(x \frac{q^n}{\sqrt{1-q}}; q^2\right) \cos\left(t \frac{q^n}{\sqrt{1-q}}\right) \\ &\quad + q^n \sin\left(x \frac{q^n}{\sqrt{1-q}}; q^2\right) \sin\left(t \frac{q^n}{\sqrt{1-q}}; q^2\right). \end{aligned}$$

Lemma 4.1 *The series*

$$\sum_{n=-\infty}^{\infty} q^{n(1-\alpha)} e\left(ix \frac{q^n}{\sqrt{1-q}}; q^2\right) e\left(-it \frac{q^n}{\sqrt{1-q}}; q^2\right) \quad (x, t \in \widetilde{\mathbb{R}}_q) \quad (4.2)$$

is absolutely convergent only when $\operatorname{Re} \alpha < 1$.

Proof The series in (4.2) can be written as

$$\left(\sum_{n=-\infty}^{-1} + \sum_{n=0}^{\infty} \right) q^{n(1-\alpha)} e\left(ix \frac{q^n}{\sqrt{1-q}}; q^2\right) e\left(-it \frac{q^n}{\sqrt{1-q}}; q^2\right).$$

From (2.4), the series $\sum_{n=0}^{\infty} q^{n(1-\alpha)} e\left(ix \frac{q^n}{\sqrt{1-q}}; q^2\right) e\left(-it \frac{q^n}{\sqrt{1-q}}; q^2\right)$ is absolutely convergent for $\operatorname{Re} \alpha < 1$ and diverges for $\operatorname{Re} \alpha \geq 1$, while the series $\sum_{n=-\infty}^{-1} q^{n(1-\alpha)} e\left(ix \frac{q^n}{\sqrt{1-q}}; q^2\right) e\left(-it \frac{q^n}{\sqrt{1-q}}; q^2\right)$ is absolutely convergent for all $\alpha \in \mathbb{C}$. \square

Set

$$\begin{aligned} \Psi_{\alpha}(x, t) &:= (1-q)^{\alpha/2} \sum_{n=-\infty}^{\infty} q^n \frac{e\left(ix \frac{q^n}{\sqrt{1-q}}; q^2\right) e\left(-it \frac{q^n}{\sqrt{1-q}}; q^2\right)}{(iq^n)^{\alpha}} \\ &\quad + (1-q)^{\alpha/2} \sum_{n=-\infty}^{\infty} q^n \frac{e\left(-ix \frac{q^n}{\sqrt{1-q}}; q^2\right) e\left(it \frac{q^n}{\sqrt{1-q}}; q^2\right)}{(-iq^n)^{\alpha}}, \end{aligned}$$

where $x, t \in \widetilde{\mathbb{R}}_q$ and i^{α} is defined with respect to the principal branch, i.e., $i^{\alpha} = e^{i\frac{\pi}{2}\alpha}$.

Lemma 4.2 *For $k \in \mathbb{N}$ and $\operatorname{Re} \alpha < 1$,*

$$\delta_{q,x}^k \Psi_{\alpha}(x, t) = \Psi_{\alpha-k}(x, t).$$

Proof The proof follows directly by using that

$$\delta_{q,x}^k e\left(ix \frac{q^n}{\sqrt{1-q}}; q^2\right) = \left(\frac{iq^n}{\sqrt{1-q}}\right)^k e\left(ix \frac{q^n}{\sqrt{1-q}}; q^2\right). \quad \square$$

A direct calculation yields the following identity, which holds for $\operatorname{Re} \alpha < 1$,

$$(1-q)^{-\alpha/2} \Psi_\alpha(x, t) = 2 \cos\left(\frac{\pi}{2}\alpha\right) A_\alpha(x, t) + 2 \sin\left(\frac{\pi}{2}\alpha\right) B_\alpha(x, t), \quad (4.3)$$

where

$$A_\alpha(x, t) := \sum_{k=-\infty}^{\infty} q^{k(1-\alpha)} \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) \cos\left(\frac{q^{m+k}}{1-q}; q^2\right) + q^{k(1-\alpha)} \sin\left(\frac{q^{n+k}}{1-q}; q^2\right) \sin\left(\frac{q^{m+k}}{1-q}; q^2\right) \quad (4.4)$$

and

$$B_\alpha(x, t) := \sum_{k=-\infty}^{\infty} q^{k(1-\alpha)} \cos\left(\frac{q^{m+k}}{1-q}; q^2\right) \sin\left(\frac{q^{n+k}}{1-q}; q^2\right) - q^{k(1-\alpha)} \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) \sin\left(\frac{q^{m+k}}{1-q}; q^2\right). \quad (4.5)$$

Theorem 4.3 For $\operatorname{Re}(\alpha) < 1$,

$$(1-q)^{-\alpha/2} \Psi_\alpha\left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}}\right) = \begin{cases} 2 \cos\left(\frac{\pi}{2}\alpha\right) q^{-m(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r} \\ \quad - 2 \sin\left(\frac{\pi}{2}\alpha\right) q^{-m(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r} \\ \quad \text{if } n > m + [\frac{1-\operatorname{Re} \alpha}{2}], \\ 2 \cos\left(\frac{\pi}{2}\alpha\right) q^{-n(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r} \\ \quad - 2 \sin\left(\frac{\pi}{2}\alpha\right) q^{-n(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r} \\ \quad \text{if } m > n + [\frac{1-\operatorname{Re} \alpha}{2}], \end{cases} \quad (4.6)$$

where $\tau_r = \begin{cases} q^{-\frac{r}{2}} q^{-\frac{1-\alpha}{2}}, & r \text{ is odd}, \\ q^{\frac{r}{2}}, & r \text{ is even}. \end{cases}$

Moreover,

$$(1-q)^{-\alpha/2} \Psi_\alpha\left(\frac{q^n}{\sqrt{1-q}}, -\frac{q^m}{\sqrt{1-q}}\right) = \begin{cases} 2 \cos\left(\frac{\pi}{2}\alpha\right) q^{-m(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r} \\ \quad + 2 \sin\left(\frac{\pi}{2}\alpha\right) q^{-m(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r}, \\ \quad n > m + [\frac{1-\operatorname{Re} \alpha}{2}], \\ 2 \cos\left(\frac{\pi}{2}\alpha\right) q^{-n(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r} \\ \quad + 2 \sin\left(\frac{\pi}{2}\alpha\right) q^{-n(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r}, \\ \quad m > n + [\frac{1-\operatorname{Re} \alpha}{2}], \end{cases} \quad (4.7)$$

$$(1-q)^{-\alpha/2} \Psi_{\alpha} \left(-\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}} \right) = \begin{cases} 2 \cos(\frac{\pi}{2}\alpha) q^{-m(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r} \\ \quad - 2 \sin(\frac{\pi}{2}\alpha) q^{-m(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r} \\ \quad \text{if } n > m + [\frac{1-\operatorname{Re}\alpha}{2}], \\ 2 \cos(\frac{\pi}{2}\alpha) q^{-n(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r} \\ \quad - 2 \sin(\frac{\pi}{2}\alpha) q^{-n(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r} \\ \quad \text{if } m > n + [\frac{1-\operatorname{Re}\alpha}{2}], \end{cases} \quad (4.8)$$

$$(1-q)^{-\alpha/2} \Psi_{\alpha} \left(-\frac{q^n}{\sqrt{1-q}}, -\frac{q^m}{\sqrt{1-q}} \right) = \begin{cases} 2 \cos(\frac{\pi}{2}\alpha) q^{-m(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r} \\ \quad + 2 \sin(\frac{\pi}{2}\alpha) q^{-m(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r} \\ \quad \text{if } n > m + [\frac{1-\operatorname{Re}\alpha}{2}], \\ 2 \cos(\frac{\pi}{2}\alpha) q^{-n(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r} \\ \quad + 2 \sin(\frac{\pi}{2}\alpha) q^{-n(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r} \\ \quad \text{if } m > n + [\frac{1-\operatorname{Re}\alpha}{2}]. \end{cases} \quad (4.9)$$

Proof Using the following formula from [14, p.455]

$$\begin{aligned} & \sum_{k=-\infty}^{\infty} s^k y^{n+k} \frac{(y^2; q^2)_\infty}{(q^2; q^2)_\infty} {}_1\phi_1(0; y^2; q^2, q^{2n+2k+2}) x^{m+k} \frac{(x^2; q^2)_\infty}{(q^2; q^2)_\infty} {}_1\phi_1(0; x^2; q^2, q^{2m+2k+2}) \\ &= s^{-m} y^{n-m} \frac{(s^{-1}xy^{-1}, y^2; q^2)_\infty}{(sxy, q^2; q^2)_\infty} {}_2\phi_1(q^2sx^{-1}y, sxy; y^2; q^2, q^{2n-2m}s^{-1}xy^{-1}) \\ &= s^{-n} x^{m-n} \frac{(s^{-1}yx^{-1}, x^2; q^2)_\infty}{(sxy, q^2; q^2)_\infty} {}_2\phi_1(q^2sxy^{-1}, sxy; x^2; q^2, q^{2m-2n}s^{-1}yx^{-1}), \end{aligned}$$

where $|sxy| < 1$, we can prove that

$$\begin{aligned} & \sum_{k=-\infty}^{\infty} q^{k(1-\alpha)} \cos\left(\frac{q^{m+k}}{1-q}; q^2\right) \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) \\ &= \begin{cases} q^{-m(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} {}_2\phi_1(q^{2-\alpha}, q^{1-\alpha}; q; q^2, q^{2n-2m+\alpha}), & n \geq m + [\operatorname{Re}\alpha/2], \\ q^{-n(1-\alpha)} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} {}_2\phi_1(q^{2-\alpha}, q^{1-\alpha}; q; q^2, q^{2m-2n+\alpha}), & m \geq n + [\operatorname{Re}\alpha/2], \end{cases} \end{aligned}$$

where $\operatorname{Re}(1-\alpha) > 0$ and

$$\begin{aligned} & \sum_{k=-\infty}^{\infty} q^{k(1-\alpha)} \sin\left(\frac{q^{m+k}}{1-q}; q^2\right) \sin\left(\frac{q^{n+k}}{1-q}; q^2\right) \\ &= \begin{cases} q^{n-m} \frac{q^{-m(1-\alpha)}}{1-q} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{3-\alpha}, q; q^2)_\infty} {}_2\phi_1(q^{2-\alpha}, q^{3-\alpha}; q^3; q^2, q^{2n-2m+\alpha}), & n > m + [\operatorname{Re}\alpha/2], \\ q^{m-n} \frac{q^{-n(1-\alpha)}}{1-q} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{3-\alpha}, q; q^2)_\infty} {}_2\phi_1(q^{2-\alpha}, q^{3-\alpha}; q^3; q^2, q^{2m-2n+\alpha}), & m > n + [\operatorname{Re}\alpha/2]. \end{cases} \end{aligned}$$

Also,

$$\begin{aligned} & \sum_{k=-\infty}^{\infty} q^{k(1-\alpha)} \operatorname{Sin}\left(\frac{q^{m+k}}{1-q}; q^2\right) \operatorname{Cos}\left(\frac{q^{n+k}}{1-q}; q^2\right) \\ &= \begin{cases} q^{-m(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} {}_2\phi_1(q^{1-\alpha}, q^{2-\alpha}; q; q^2, q^{2n-2m+\alpha+1}), \\ q^{m-n} \frac{q^{-n(1-\alpha)}}{1-q} \frac{(q^{\alpha-1}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} {}_2\phi_1(q^{2-\alpha}, q^{3-\alpha}; q^3; q^2, q^{2m-2n+\alpha-1}). \end{cases} \end{aligned}$$

Hence, if $x := \frac{q^n}{\sqrt{1-q}}$ and $t := \frac{q^m}{\sqrt{1-q}}$, then

$$A_\alpha(x, t) = \begin{cases} t^{-(1-\alpha)} (1-q)^{-(1-\alpha)/2} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} \left(\frac{x}{t}\right)^r, & n > m + [-\operatorname{Re} \alpha/2], \\ x^{-(1-\alpha)} (1-q)^{-(1-\alpha)/2} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} \left(\frac{t}{x}\right)^r, & m > n + [-\operatorname{Re} \alpha/2]. \end{cases} \quad (4.10)$$

Also,

$$A_\alpha(x, -t) = \begin{cases} t^{-(1-\alpha)} (1-q)^{-(1-\alpha)/2} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} \left(\frac{x}{t}\right)^r, & n > m + [-\operatorname{Re} \alpha/2], \\ x^{-(1-\alpha)} (1-q)^{-(1-\alpha)/2} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\alpha[\frac{r}{2}]} \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} \left(\frac{t}{x}\right)^r, & m > n + [-\operatorname{Re} \alpha/2], \end{cases} \quad (4.11)$$

$$B_\alpha(x, t) = - \begin{cases} q^{-m(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r}, & n \geq m + [\frac{1-\operatorname{Re} \alpha}{2}], \\ q^{-n(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r}, & m > n + [\frac{1-\operatorname{Re} \alpha}{2}], \end{cases} \quad (4.12)$$

$$B_\alpha(x, -t) = \begin{cases} q^{-m(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(n-m)r}, & n \geq m + [\frac{1-\operatorname{Re} \alpha}{2}], \\ q^{-n(1-\alpha)} \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{r=0}^{\infty} (-1)^r q^{\frac{\alpha r}{2}} \tau_r \frac{(q^{1-\alpha}; q)_r}{(q; q)_r} q^{(m-n)r}, & m > n + [\frac{1-\operatorname{Re} \alpha}{2}]. \end{cases} \quad (4.13)$$

Substituting from (4.10)-(4.13) into (4.3) yields the values $\Psi_\alpha(\pm x, \pm t)$ and the theorem follows. \square

Remark 4.4 In the previous theorem, we calculated the value of $\Psi_\alpha(x, t)$, $x, t \in \tilde{\mathbb{R}}_q$ and for specific values of x, t . We can calculate the values of $\Psi_\alpha(x, t)$ for all x, t by using the

identity

$$\begin{aligned} & \sum_{-\infty}^{\infty} s^{k+m} y^{k+n} \frac{(y^2; q^2)_{\infty}}{(q^2; q^2)_{\infty}} {}_1\phi_1(0; y^2; q^2, q^{2k+2n+2}) x^{k+m} \frac{(x^2; q^2)_{\infty}}{(q^2; q^2)_{\infty}} {}_1\phi_1(0; x^2; q^2, q^{2k+2m+2}) \\ &= y^{n-m} \frac{(s^{-1}xy^{-1}, q^{2n-2m+2}; q^2)_{\infty}}{(q^{2n-2m}s^{-1}xy^{-1}, q^2; q^2)_{\infty}} {}_2\phi_1(q^{2n-2m}s^{-1}xy^{-1}, s^{-1}yx^{-1}; q^{2n-2m+2}; q^2, sxy) \end{aligned}$$

for $|sxy| < 1$. See Proposition 4.1 of [14].

In this case we have

$$\begin{aligned} & CC_{\alpha}\left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}}\right) \\ &:= \sum_{r=-\infty}^{\infty} q^{r(1-\alpha)} \cos\left(\frac{q^{n+r}}{1-q}; q^2\right) \cos\left(\frac{q^{m+r}}{1-q}; q^2\right) \\ &= q^{-m(1-\alpha)} \frac{(q^{\alpha}, q^{2n-2m+2}, q^2; q^2)_{\infty}}{(q^{2n-2m+\alpha}, q, q; q^2)_{\infty}} {}_2\phi_1(q^{2n-2m+\alpha}, q^{\alpha}; q^{2n-2m+2}; q^2, q^{1-\alpha}), \quad (4.14) \end{aligned}$$

$$\begin{aligned} & SS_{\alpha}\left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}}\right) \\ &:= \sum_{r=-\infty}^{\infty} q^{r(1-\alpha)} \sin\left(\frac{q^{n+r}}{1-q}; q^2\right) \sin\left(\frac{q^{m+r}}{1-q}; q^2\right) \\ &= q^{n-m} q^{-m(1-\alpha)} \frac{(q^{\alpha}, q^{2n-2m+2}, q^2; q^2)_{\infty}}{(q^{2n-2m+\alpha}, q, q; q^2)_{\infty}} {}_2\phi_1(q^{2n-2m+\alpha}, q^{\alpha}; q^{2n-2m+2}; q^2, q^{3-\alpha}), \quad (4.15) \end{aligned}$$

$$\begin{aligned} & SC_{\alpha}\left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}}\right) \\ &:= \sum_{r=-\infty}^{\infty} q^{r(1-\alpha)} \sin\left(\frac{q^{n+r}}{1-q}; q^2\right) \cos\left(\frac{q^{m+r}}{1-q}; q^2\right) \\ &= q^{n-m} q^{-m(1-\alpha)} \frac{(q^{\alpha-1}, q^{2n-2m+2}, q^2; q^2)_{\infty}}{(q^{2n-2m+\alpha-1}, q, q; q^2)_{\infty}} {}_2\phi_1(q^{2n-2m+\alpha-1}, q^{\alpha+1}; q^{2n-2m+2}; q^2, q^{2-\alpha}), \\ & CS_{\alpha}\left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}}\right) \\ &:= \sum_{r=-\infty}^{\infty} q^{r(1-\alpha)} \sin\left(\frac{q^{m+r}}{1-q}; q^2\right) \cos\left(\frac{q^{n+r}}{1-q}; q^2\right) \\ &= q^{-m(1-\alpha)} \frac{(q^{\alpha+1}, q^{2n-2m+2}, q^2; q^2)_{\infty}}{(q^{2n-2m+\alpha+1}, q, q; q^2)_{\infty}} {}_2\phi_1(q^{2n-2m+\alpha+1}, q^{\alpha-1}; q^{2n-2m+2}; q^2, q^{2-\alpha}). \end{aligned}$$

Corollary 4.5 For each fixed $x, t \in \widetilde{\mathbb{R}}_q$, the function $\Psi_{\alpha}(x, t)$ as a function of α can be extended to an entire function on \mathbb{C} .

Proof If α is a positive integer, then the series on the right-hand sides of (4.6)-(4.9) are finite sums and hence are convergent. Since the zeros of the function $\cos(\frac{\pi}{2}\alpha)$ are the poles of the function $(q^{1-\alpha}; q^2)_{\infty}$ with the same orders. In fact

$$\lim_{\alpha \rightarrow (2j+1)} \frac{\cos \frac{\pi}{2}\alpha}{(q^{1-\alpha}; q^2)_{\infty}} = -\frac{\pi}{2 \ln q} \frac{q^{2+j}}{(q^2; q^2)_j (q^2; q^2)_{\infty}} \quad (j \in \mathbb{N}_0).$$

Similarly, the zeros of the function $\sin(\frac{\pi}{2}\alpha)$ are the poles of the function $(q^{2-\alpha}; q^2)_{\infty}$ with the same orders and

$$\lim_{\alpha \rightarrow (2j)} \frac{\sin \frac{\pi}{2}\alpha}{(q^{2-\alpha}; q^2)_{\infty}} = -\frac{\pi}{2 \ln q} \frac{q^{j^2-j}}{(q^2; q^2)_{j-1}(q^2; q^2)_{\infty}} \quad (j \in \mathbb{N}).$$

Then the left-hand sides of equations (4.6)-(4.9) are entire functions. Hence, as a function of α , the functions $\Psi_{-\alpha}(x, t)$ ($x, t \in \tilde{\mathbb{R}}_q$) can be analytically extended by defining its values when $\Re \alpha \geq 1$ by the left-hand sides of (4.6)-(4.9). \square

It also should be noted that for $\operatorname{Re}(\alpha) \geq 1$, the left-hand sides of (4.6) and (4.7) determine $\Psi_{-\alpha}(x, t)$ for all $x, t \in \tilde{\mathbb{R}}_q$, which is different from the case of $\operatorname{Re}(\alpha) < 1$; see Remark 4.4.

Definition 4.6 For $\operatorname{Re} \alpha > 0$, we define a fractional q -integral operator J_q^{α} on $L^2(\tilde{\mathbb{R}}_q) \cap L^1(\tilde{\mathbb{R}}_q)$ by

$$J_q^{\alpha} f(x) := \frac{1}{C} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) \Psi_{\alpha}(x, t) d_q t.$$

The following properties follow at once from (4.3) and their analytic continuation on \mathbb{C} .

- If $f \in L^2(\tilde{\mathbb{R}}_q) \cap L(\tilde{\mathbb{R}}_q)$ and f is even, then

$$J_q^{\alpha} f(x) = 2 \int_0^{\infty/\sqrt{1-q}} f(t) \varphi_{\alpha}(x, t) d_q t,$$

where

$$(1-q)^{-\alpha/2} \varphi_{\alpha} \left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}} \right) \\ = \begin{cases} 2q^{-m(1-\alpha)} \frac{(q^{\alpha}, q^2; q^2)_{\infty}}{(q^{1-\alpha}, q; q^2)_{\infty}} \cos \frac{\pi}{2}\alpha \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j}}{(q; q)_{2j}} q^{j(2n-2m+\alpha)} \\ \quad - 2q^{n-m-(m+1)(1-\alpha)} \sin \frac{\pi}{2}\alpha \frac{(q^{1+\alpha}, q^2; q^2)_{\infty}}{(q^{2-\alpha}, q; q^2)_{\infty}} \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j+1}}{(q; q)_{2j+1}} q^{j(2n-2m+\alpha-1)}, \\ 2q^{-n(1-\alpha)} \frac{(q^{\alpha}, q^2; q^2)_{\infty}}{(q^{1-\alpha}, q; q^2)_{\infty}} \cos \frac{\pi}{2}\alpha \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j}}{(q; q)_{2j}} q^{j(2m-2n+\alpha)} \\ \quad + 2q^{-n(1-\alpha)} \sin \frac{\pi}{2}\alpha \frac{(q^{1+\alpha}, q^2; q^2)_{\infty}}{(q^{2-\alpha}, q; q^2)_{\infty}} \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j}}{(q; q)_{2j}} q^{j(2m-2n+\alpha+1)} \end{cases}$$

and

$$(1-q)^{-\alpha/2} \varphi_{\alpha} \left(-\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}} \right) \\ = \begin{cases} 2q^{-m(1-\alpha)} \frac{(q^{\alpha}, q^2; q^2)_{\infty}}{(q^{1-\alpha}, q; q^2)_{\infty}} \cos \frac{\pi}{2}\alpha \sum_{j=0}^{\infty} \frac{(q^{\alpha+1}; q)_{2j}}{(q; q)_{2j}} q^{j(2n-2m+\alpha)} \\ \quad + 2q^{n-m-m(1-\alpha)} \sin \frac{\pi}{2}\alpha \frac{(q^{1+\alpha}, q^2; q^2)_{\infty}}{(q^{2-\alpha}, q; q^2)_{\infty}} \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j+1}}{(q; q)_{2j+1}} q^{j(2n-2m+\alpha-1)}, \\ 2q^{-n(1-\alpha)} \frac{(q^{\alpha}, q^2; q^2)_{\infty}}{(q^{1-\alpha}, q; q^2)_{\infty}} \cos \frac{\pi}{2}\alpha \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j}}{(q; q)_{2j}} q^{j(2m-2n+\alpha)} \\ \quad - 2q^{-n(1-\alpha)} \sin \frac{\pi}{2}\alpha \frac{(q^{1+\alpha}, q^2; q^2)_{\infty}}{(q^{2-\alpha}, q; q^2)_{\infty}} \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j}}{(q; q)_{2j}} q^{j(2m-2n+\alpha+1)}. \end{cases}$$

Hence, if f is an even function, then $\delta_q^{2k} f(x)$ is an even function and $\delta_q^{2k+1} f(x)$ is an odd function for all $k \in \mathbb{N}_0$.

- If $f \in L^2(\widetilde{\mathbb{R}}_q) \cap L(\widetilde{\mathbb{R}}_q)$ and f is odd, then

$$J_q^\alpha f(x) = 2 \int_0^{\infty/\sqrt{1-q}} f(t) \psi_\alpha(x, t) d_q t,$$

where

$$(1-q)^{-\alpha/2} \psi_\alpha \left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}} \right) = \begin{cases} 2q^{-m(1-\alpha)} q^{n-m} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \cos \frac{\pi}{2} \alpha \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j+1}}{(q; q)_{2j+1}} q^{j(2n-2m+\alpha)} \\ - 2q^{-m(1-\alpha)} \sin \frac{\pi}{2} \alpha \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j}}{(q; q)_{2j}} q^{j(2n-2m+\alpha+1)}, \\ 2q^{-n(1-\alpha)} q^{m-n} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \cos \frac{\pi}{2} \alpha \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j+1}}{(q; q)_{2j+1}} q^{j(2m-2n+\alpha)} \\ + 2q^{-(n+1)(1-\alpha)} q^{m-n} \sin \frac{\pi}{2} \alpha \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j+1}}{(q; q)_{2j+1}} q^{j(2m-2n+\alpha-1)} \end{cases}$$

and

$$(1-q)^{-\alpha/2} \psi_\alpha \left(-\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}} \right) = \begin{cases} -2q^{-m(1-\alpha)} q^{n-m} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \cos \frac{\pi}{2} \alpha \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j+1}}{(q; q)_{2j+1}} q^{j(2n-2m+\alpha)} \\ - 2q^{-m(1-\alpha)} \sin \frac{\pi}{2} \alpha \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j}}{(q; q)_{2j}} q^{j(2n-2m+\alpha+1)}, \\ 2q^{-n(1-\alpha)} q^{m-n} \frac{(q^\alpha, q^2; q^2)_\infty}{(q^{1-\alpha}, q; q^2)_\infty} \cos \frac{\pi}{2} \alpha \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j+1}}{(q; q)_{2j+1}} q^{j(2m-2n+\alpha)} \\ - 2q^{-(n+1)(1-\alpha)} q^{m-n} \sin \frac{\pi}{2} \alpha \frac{(q^{1+\alpha}, q^2; q^2)_\infty}{(q^{2-\alpha}, q; q^2)_\infty} \sum_{j=0}^{\infty} \frac{(q^{1-\alpha}; q)_{2j+1}}{(q; q)_{2j+1}} q^{j(2m-2n+\alpha-1)}. \end{cases}$$

Hence, if f is an odd function, then $\delta_q^{2k} f(x)$ is an odd function and $\delta_q^{2k+1} f(x)$ is an even function for all $k \in \mathbb{N}_0$.

Theorem 4.7 Let $f \in L^2(\widetilde{\mathbb{R}}_q)$. Then

$$\delta_q(J_q f)(x) = f(x) \quad \text{for all } x \in \widetilde{\mathbb{R}}_q. \quad (4.16)$$

Moreover, if f is q -regular at zero, then

$$J_q(\delta_q f)(x) = f(x) \quad \text{for all } x \in \widetilde{\mathbb{R}}_q. \quad (4.17)$$

Proof If $f \in L^2(\widetilde{\mathbb{R}}_q)$, then

$$J_q f(x) = \frac{2}{C} \int_0^{\infty/\sqrt{1-q}} f_e(t) \varphi_1(x, t) d_q t + \frac{2}{C} \int_0^{\infty/\sqrt{1-q}} f_e(t) \psi_1(x, t) d_q t,$$

where

$$\varphi_1 \left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}} \right) = \begin{cases} \frac{\pi \sqrt{1-q}}{\ln q}, & n > m, \\ \frac{\pi \sqrt{1-q}}{\ln q} + \frac{C}{2}, & m \geq n \end{cases}$$

and

$$\psi_1\left(\frac{q^n}{\sqrt{1-q}}, \frac{q^m}{\sqrt{1-q}}\right) = \begin{cases} -\frac{C}{2}, & n > m - 1, \\ 0, & m > n. \end{cases}$$

Hence,

$$\begin{aligned} J_q f(x) &= \left(1 + \frac{2\pi\sqrt{1-q}}{C \ln q}\right) \int_0^x f_e(t) d_q t + \frac{2\pi\sqrt{1-q}}{C \ln q} \int_x^\infty f_e(t) d_q t \\ &\quad - \int_{qx}^\infty f_0(t) d_q t. \end{aligned} \tag{4.18}$$

One can verify that the first and second q -integrals of (4.18) are odd functions, while the last q -integral is an even function. Consequently,

$$\begin{aligned} \delta_q J_q f(x) &= \left(1 + \frac{2\pi\sqrt{1-q}}{C \ln q}\right) D_{q,x} \int_0^x f_e(t) d_q t + \frac{2\pi\sqrt{1-q}}{C \ln q} D_{q,x} \int_x^\infty f_e(t) d_q t \\ &\quad - D_{q,x} \int_x^\infty f_0(t) d_q t. \end{aligned} \tag{4.19}$$

Using the fundamental theorem of q -calculus, see [13], we obtain (4.16). To prove (4.17), we assume that f is q -regular at zero,

$$\begin{aligned} J_q \delta_q f(x) &= J_q(\delta_q f_e)(x) + J_q(\delta_q f_0)(x) \\ &= \frac{2}{C} \int_0^{\infty/\sqrt{1-q}} \delta_q f_e(t) \psi_1(x, t) d_q t + \frac{2}{C} \int_0^{\infty/\sqrt{1-q}} \delta_q f_0(t) \varphi_1(x, t) d_q t. \end{aligned}$$

But

$$\begin{aligned} \int_0^{\infty/\sqrt{1-q}} \delta_q f_e(t) \psi_1(x, t) d_q t &= \int_0^{\infty/\sqrt{1-q}} D_{q,t} f_e(t/q) \psi_1(x, t) d_q t \\ &= -\frac{C}{2} \int_{qx}^{\infty/\sqrt{1-q}} D_{q,t} f_e(t/q) d_q t = \frac{C}{2} f_e(x) \end{aligned}$$

and

$$\begin{aligned} \int_0^{\infty/\sqrt{1-q}} \delta_q f_0(t) \varphi_1(x, t) d_q t &= \int_0^{\infty/\sqrt{1-q}} D_{q,t} f_0(t) \varphi_1(x, t) d_q t \\ &= \left(\frac{\pi\sqrt{1-q}}{\ln q} + \frac{C}{2}\right) \int_0^x D_{q,t} f_0(t) d_q t + \frac{\pi\sqrt{1-q}}{\ln q} \int_x^\infty D_{q,t} f_0(t) d_q t \\ &= \frac{C}{2} f_o(x) - \left(\frac{\pi\sqrt{1-q}}{\ln q} + \frac{C}{2}\right) f_0(0) = \frac{C}{2} f_o(x) \end{aligned}$$

since $f_0(0) = 0$. This proves (4.17) and completes the proof of the theorem. \square

We can also prove that

$$\begin{aligned}\psi_2(x, t) &:= \begin{cases} -\frac{c}{2}(x - qt) + (x - t)\frac{\pi\sqrt{1-q}}{\ln q}, & x \leq q^{-1}t, \\ -\frac{c}{2}(t - qx) + (x - t)\frac{\pi\sqrt{1-q}}{\ln q}, & x \geq qt, \end{cases} \\ \psi_2(x, -t) &:= \begin{cases} \frac{c}{2}(x + qt) + (x + t)\frac{\pi\sqrt{1-q}}{\ln q}, & x \leq q^{-1}t, \\ \frac{c}{2}(t + qx) + (x + t)\frac{\pi\sqrt{1-q}}{\ln q}, & x \geq qt \end{cases}\end{aligned}$$

and

$$\psi_2(-x, t) := \begin{cases} \frac{c}{2}(x + qt) - (x + t)\frac{\pi\sqrt{1-q}}{\ln q}, & x \leq q^{-1}t, \\ \frac{c}{2}(t + qx) - (x + t)\frac{\pi\sqrt{1-q}}{\ln q}, & x \geq qt. \end{cases}$$

Hence,

$$\begin{aligned}J_q^2 f(x) &= \frac{\pi\sqrt{1-q}}{c\ln q} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} (x - t)f(t) d_q t \\ &\quad + \int_0^{q^{-1}x} (qxf_e(t) - tf_0(t)) d_q t + \int_{q^{-1}x}^{\infty/\sqrt{1-q}} (xf_e(t) - qtf_0(t)) d_q t.\end{aligned}$$

Definition 4.8 For $\alpha > 0$, we define a fractional q -difference operator δ_q^α on $L^2(\widetilde{\mathbb{R}}_q) \cap L^1(\widetilde{\mathbb{R}}_q)$ by

$$\delta_q^\alpha f(x) := J_q^{-\alpha} f(x) = \frac{1}{C} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) \Psi_{-\alpha}(x, t) d_q t. \quad (4.20)$$

Lemma 4.9 The operator δ_q^α coincides with Rubin's q -difference operator when α is a positive integer.

Proof Let $\alpha = k$ for some $k \in \mathbb{N}$, and let $f \in L^2(\widetilde{\mathbb{R}}_q) \cap L^1(\widetilde{\mathbb{R}}_q)$. Using (2.7), we conclude

$$f(x) = \frac{1}{C} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) \Psi_0(x, t) d_q t.$$

Then from Lemma 4.2 we obtain

$$\delta_q^k f(x) = \frac{1}{C} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) \Psi_{-k}(x, t) d_q t,$$

and the lemma follows. \square

Lemma 4.10 If $\alpha > 0$ and $f \in L^2(\widetilde{\mathbb{R}}_q)$, then

$$\delta_q^\alpha f(x) = \delta_q^k J_q^{k-\alpha} f(x) \quad (\alpha > 0; k = \lceil \alpha \rceil; x \in \mathbb{R}_q). \quad (4.21)$$

Now, if α is a positive integer, then $\lceil \alpha \rceil = \alpha$ and from (4.21)

$$\partial_q^\alpha = \partial_q^k.$$

Proof Since

$$\delta_q^k J_q^{k-\alpha} f(x) = \frac{1}{C} \delta_q^k \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(t) \Psi_{k-\alpha}(x, t) d_q t$$

and $\delta_{q,x}^k \Psi_{k-\alpha}(x, t) = \Psi_{-\alpha}(x, t)$, the proof follows. \square

Theorem 4.11 If $f \in L^1(\tilde{\mathbb{R}}_q) \cap L^2(\tilde{\mathbb{R}}_q)$ and α, β are complex numbers such that $\operatorname{Re}(\alpha) < 1$ and $\operatorname{Re}(\beta) < 1$ such that $\operatorname{Re}(\alpha + \beta) < 1$, then

$$J_q^\alpha J_q^\beta f = J_q^\beta J_q^\alpha f = J_q^{\alpha+\beta} f.$$

Proof Let $x \in \tilde{\mathbb{R}}_q$. Since

$$\begin{aligned} J_q^\alpha (J_q^\beta f)(x) &= \frac{1}{C^2} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(u) \Psi_\alpha(x, t) \Psi_\beta(t, u) d_q u d_q t \\ &= \frac{1}{C^2} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(u) \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} \Psi_\alpha(x, t) \Psi_\beta(t, u) d_q t d_q u, \end{aligned}$$

using the orthogonality relation (2.9), we obtain

$$\int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} \Psi_\alpha(x, t) \Psi_\beta(t, u) d_q t = C \Psi_{\alpha+\beta}(x; u).$$

Consequently,

$$J_q^\alpha (J_q^\beta f)(x) = \frac{1}{C} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f(u) \Psi_{\alpha+\beta}(x; u) d_q u = J_q^{\alpha+\beta} f(x). \quad \square$$

Example 4.12 Let $n \in \mathbb{Z}$ and let $f_n(x)$ be the even function defined on $\tilde{\mathbb{R}}_q$ by

$$f_n(x) := \begin{cases} 1, & x = \pm x_n, x_n := \frac{q^n}{\sqrt{1-q}}, \\ 0, & \text{otherwise.} \end{cases}$$

Hence,

$$f_n(x) = \sum_{k=-\infty}^{\infty} a_k \cos\left(x \frac{q^k}{\sqrt{1-q}}; q^2\right) + \sum_{k=-\infty}^{\infty} b_k \sin\left(x \frac{q^k}{\sqrt{1-q}}; q^2\right).$$

Since f_n is an even function, then $b_k = 0$ for all $k \in \mathbb{Z}$ and

$$a_k = \frac{q^k}{C} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} f_n(t) \cos\left(t \frac{q^k}{\sqrt{1-q}}; q^2\right) d_q t = \frac{2q^{n+k}\sqrt{1-q}}{C} \cos\left(\frac{q^{n+k}}{1-q}; q^2\right).$$

Consequently,

$$f_n(x) = \frac{2q^n\sqrt{1-q}}{C} \sum_{k=-\infty}^{\infty} q^k \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) \cos\left(\frac{q^k x}{\sqrt{1-q}}; q^2\right).$$

Hence,

$$\sum_{k=-\infty}^{\infty} q^k \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) \cos\left(\frac{q^k x}{\sqrt{1-q}}; q^2\right) = \begin{cases} \frac{Cq^{-n}}{2\sqrt{1-q}}, & x = x_n, \\ \text{zero}, & \text{otherwise.} \end{cases}$$

Since

$$\delta_q f_n(x) = \frac{1}{q} D_{q^{-1}} f_n(x) := \begin{cases} \frac{\pm 1}{x_n(1-q)}, & x = \pm x_n, \\ \frac{\mp 1}{qx_n(1-q)}, & x = \pm qx_n, \\ \text{zero}, & \text{otherwise} \end{cases}$$

and

$$\delta_q f_n(x) = -\frac{2q^n}{C} \sum_{k=-\infty}^{\infty} q^{2k} \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) \sin\left(\frac{q^k x}{\sqrt{1-q}}; q^2\right) \quad (x \in \widetilde{\mathbb{R}}_q).$$

We obtain

$$\sum_{k=-\infty}^{\infty} q^{2k} \sin\left(\frac{xq^k}{\sqrt{1-q}}; q^2\right) \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) = \begin{cases} \text{zero}, & x \notin \{\pm x_n, \pm qx_n\}, \\ \pm \frac{Cq^{-n}}{2x_n(1-q)}, & x = \pm x_n, \\ \pm \frac{Cq^{-n}}{2qx_n(1-q)}, & x = \pm qx_n. \end{cases} \quad (4.22)$$

Similarly,

$$\delta_q^2 f_n(x) = \begin{cases} \frac{-1}{qx_n^2(1-q)}, & x = \pm x_n, \\ \frac{-1}{q^2 x_n^2(1-q)^2}, & x = \pm qx_n, \\ \frac{-q}{x_n^2(1-q)^2}, & x = \pm q^{-1} x_n, \\ \text{zero}, & \text{otherwise.} \end{cases}$$

But

$$\delta_q^2 f_n(x) = -2 \frac{q^n}{C\sqrt{1-q}} \sum_{k=-\infty}^{\infty} q^{3k} \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) \cos\left(\frac{q^k x}{\sqrt{1-q}}; q^2\right) \quad (x \in \widetilde{\mathbb{R}}_q).$$

Hence,

$$\sum_{k=-\infty}^{\infty} q^{3k} \cos\left(\frac{q^{n+k}}{1-q}; q^2\right) \cos\left(\frac{q^k x}{\sqrt{1-q}}; q^2\right) = \begin{cases} \frac{Cq^{-n-1}}{2x_n^2\sqrt{1-q}}, & x = \pm x_n, \\ \frac{Cq^{-n-2}}{2x_n^2(\sqrt{1-q})^3}, & x = \pm qx_n, \\ \frac{Cq^{-n+1}}{2x_n^2(\sqrt{1-q})^3}, & x = \pm q^{-1} x_n, \\ 0, & \text{otherwise.} \end{cases} \quad (4.23)$$

In the following two examples, we show how we can use (4.20) when α is a positive integer to obtain new summation formulae.

Example 4.13 Let f be an even function. Then, for each $k \in \mathbb{N}_0$, we have

$$\delta_q^{2k} f(x) = q^{k(k-1)} D_{q,x}^{2k} f(x/q^k), \quad \delta_q^{2k+1} f(x) = q^{k(k+1)} D_{q,x}^{2k+1} f(x/q^{k+1}).$$

Hence,

$$\delta_q^{2k} f(x) = q^{k(k-1)} (1-q)^{-2k} x^{-2k} \sum_{r=0}^{2k} q^r \frac{(q^{-2k};q)_r}{(q;q)_r} f(xq^{r-k}).$$

On the other hand,

$$\delta_q^{2k} f(x) = 4 \frac{(-1)^k}{C} \int_0^{\infty/\sqrt{1-q}} f(t) CC_{-2k}(x,t) d_q t.$$

Let $m \in \mathbb{Z}$ and let $f_m(x)$ be the even function defined on $\widetilde{\mathbb{R}}_q$ by

$$f_m(x) := \begin{cases} 1, & x = \pm x_m, x_m := \frac{q^n}{\sqrt{1-q}}, \\ 0, & \text{otherwise.} \end{cases}$$

Set $x = \frac{q^n}{\sqrt{1-q}}$. Hence,

$$\begin{aligned} \delta_q^{2k} f(x) &= 4 \frac{(-1)^k \sqrt{1-q}}{C} q^{-2mk} (1-q)^{-k} \frac{(q^{-2k}, q^{2n-2m+2}, q^2; q^2)_\infty}{(q^{2n-2m-2k}, q, q; q^2)_\infty} \\ &\quad \times {}_2\phi_1(q^{2n-2m-2k}, q^{-2k}; q^{2n-2m+2}; q^2, q^{2k+1}) \\ &= q^{k^2+m-n-2nk} (1-q)^{-k} \frac{(q^{-2k}; q)_{m-n+k}}{(q; q)_{m-n+k}}. \end{aligned}$$

Since

$$\frac{(q^{-2k}, q^{2n-2m+2}; q^2)_\infty}{(q^{2n-2m-2k}; q^2)_\infty} = 0 \quad \text{if } n \geq m+k \text{ or } n < m.$$

Hence, if $m \leq n \leq m+k-1$, we obtain

$$\begin{aligned} {}_2\phi_1(q^{2n-2m-2k}, q^{-2k}; q^{2n-2m+2}; q^2, q^{2k+1}) \\ = (-1)^k q^{k^2+(m-n)(2k+1)} \frac{(q^{-2k}; q)_{m-n+k}}{(q; q)_{m-n+k}} \frac{(q^2, q^{2n-2m-2k}; q^2)_\infty}{(q^{-2k}, q^{2n-2m+2}; q^2)_\infty}. \end{aligned}$$

Example 4.14 Let f be an odd function. Then, for each $k \in \mathbb{N}_0$, we have

$$\delta_q^{2k} f(x) = q^{k^2} D_{q,x}^{2k} f(x/q^k), \quad \delta_q^{2k+1} f(x) = q^{k^2} D_{q,x}^{2k+1} f(x/q^k).$$

Hence,

$$\delta_q^{2k} f(x) = q^{k^2} (1-q)^{-2k} x^{-2k} \sum_{r=0}^{2k} q^r \frac{(q^{-2k}; q)_r}{(q; q)_r} f(xq^{r-k}).$$

On the other hand,

$$\delta_q^{2k} f(x) = 4 \frac{(-1)^k}{C} \int_0^{\infty/\sqrt{1-q}} f(t) CS_{-2k}(x, t) d_q t.$$

Let $m \in \mathbb{Z}$ and let $f_m(x)$ be the odd function defined on $\tilde{\mathbb{R}}_q$ by

$$f_m(x) := \begin{cases} \pm 1, & x = \pm x_m, x_m := \frac{q^m}{\sqrt{1-q}}, \\ 0, & \text{otherwise.} \end{cases}$$

Set $x = \frac{q^n}{\sqrt{1-q}}$. Hence,

$$\begin{aligned} \delta_q^{2k} f(x) &= 4 \frac{(-1)^k \sqrt{1-q}}{C} q^{-2mk} (1-q)^{-k-\frac{1}{2}} \frac{(q^{-2k}, q^{2n-2m+2}, q^2; q^2)_\infty}{(q^{2n-2m-2k}, q, q; q^2)_\infty} \\ &\quad \times {}_2\phi_1(q^{2n-2m-2k}, q^{-2k-2}; q^{2n-2m+2}, q^2, q^{2k+3}) \\ &= q^{k^2+k+m-n} (1-q)^{-k-\frac{1}{2}} \frac{(q^{-2k-1}; q)_{m-n+k}}{(q; q)_{m-n+k}}. \end{aligned}$$

Since

$$\frac{(q^{-2k}, q^{2n-2m+2}; q^2)_\infty}{(q^{2n-2m-2k}; q^2)_\infty} = 0 \quad \text{if } n \geq m+k \text{ or } n < m.$$

Hence, if $m \leq n \leq m+k-1$, we obtain

$$\begin{aligned} &{}_2\phi_1(q^{2n-2m-2k}, q^{-2k-2}; q^{2n-2m+2}, q^2, q^{2k+3}) \\ &= (-1)^k q^{k^2+k+(m-n)(2k+1)} \frac{(q^{-2k-1}; q)_{m-n+k}}{(q; q)_{m-n+k}} \frac{(q^2, q^{2n-2m-2k}; q^2)_\infty}{(q^{-2k}, q^{2n-2m+2}; q^2)_\infty}. \end{aligned}$$

5 Application of the q^2 -analogue of the Fourier transform to solve q -fractional difference equations

In [18], Ho explored the possibility of using the classical Fourier and Mellin integral transforms to solve the class of q -difference differential equations

$$D_{q,t}^n u(x, t) = \frac{\partial^2}{\partial x^2} u(x, t), \quad x \in \mathbb{R}, t > 0, n \geq 1, \tag{5.1}$$

with the initial conditions

$$y(x, 0) = f(x), \quad D_{q,t}^k y(x, t)|_{t=0^+} = g_k(x) \quad (k = 1, \dots, n-1),$$

where the functions $f(x)$ and $g_k(x)$ are assumed to vanish as $x \rightarrow \pm\infty$. In [19] Brahim and Quanes used the q^2 -Fourier transform and the q -Mellin transform to solve equation (5.1) in case of $n = 1, 2$, and only for q satisfying condition (3.2). In this section, we use the q^2 -Fourier transform with the ${}_q L_s$ transform to solve the q -fractional diffusion equation

$$\begin{aligned} D_{q,t}^\alpha u(x, t) &= \lambda \partial_{q,x}^2 u(x, t), \\ x \in \tilde{\mathbb{R}}_q, t \in \mathbb{R}, 0 &\leq \alpha < 1, 0 < q < 1, \end{aligned} \tag{5.2}$$

with the initial conditions

$$I_{q,t}^{1-\alpha} u(x,t)|_{t=0^+} = \phi(x), \quad \partial_{q,x}^k u(x,t) \in L_q^1(\tilde{\mathbb{R}}_q) \quad (k = 0, 1), \quad (5.3)$$

$$\phi \in (L_q^1 \cap L_q^2)(\tilde{\mathbb{R}}_q). \quad (5.4)$$

Theorem 5.1 *The solution of q -fractional diffusion equation (5.2) subject to the initial conditions (5.3)-(5.4) is given by*

$$u(x,t) = \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} [T_y G(x,t)](x) \phi(y) d_q y,$$

where

$$[T_y G(x,t)](x) = \frac{(1+q)^{1/2}}{2\Gamma_q(\frac{1}{2})} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} e(-iy\xi; q^2) g(\xi, t) e(ix\xi; q^2) d_q \xi$$

and

$$g(\xi, t) := \begin{cases} t^{\alpha-1} e_{\alpha,\alpha}(-\lambda \xi^2 t^\alpha; q), & |\lambda \xi^2 t^\alpha| < \frac{1}{(1-q)^\alpha}, \\ 0, & \text{otherwise,} \end{cases}$$

where, in general, $e_{\alpha,\beta}(x; q)$ is the q -analogue of the q -Mittag-Leffler function defined for $\operatorname{Re}(\alpha > 0)$ and $\beta \in \mathbb{C}$ by

$$e_{\alpha,\beta}(x; q) = \sum_{k=0}^{\infty} \frac{x^k}{\Gamma_q(\alpha k + \beta)}, \quad |x(1-q)^\alpha| < 1.$$

Proof First we calculate the q^2 -Fourier transform of (5.2) with respect to the variable x . Hence, applying (3.12) yields

$$D_{q,t}^\alpha U(\xi, t) = -\lambda \xi^2 U(\xi, t), \quad (5.5)$$

where

$$U(\xi, t) := \mathcal{F}_{q,x}(u(x, t))(\xi).$$

Now we calculate the ${}_q L_s$ transform of (5.5) with respect to the variable t . Using (1.16) we obtain

$$(P^\alpha + \lambda \xi^2) V(\xi, s) = \frac{\mathcal{F}_q(\phi)(\xi)}{1-q}, \quad (5.6)$$

where

$$V(\xi, s) = {}_{q,t} L_s(U(\xi, t))(s), \quad p = \frac{s}{1-q}.$$

One can verify that

$${}_{q,t}L_s(t^{\alpha-1}e_{\alpha,\alpha}(-\lambda\xi^2t^\alpha;q)) = \frac{1}{1-q} \frac{1}{p^\alpha + \lambda\xi^2}.$$

Consequently,

$$U(\xi, t) = \mathcal{F}_q(\phi)(\xi) t^{\alpha-1} e_{\alpha,\alpha}(-\lambda\xi^2 t^\alpha), \quad |\lambda\xi^2 t^\alpha| < \frac{1}{(1-q)^\alpha}.$$

It follows from the inversion formula of the q^2 -Fourier transform that

$$\begin{aligned} t^{\alpha-1} e_{\alpha,\alpha}(-\lambda\xi^2 t^\alpha; q) &= \mathcal{F}_{q,x}(G(x, t))(\xi), \\ G(x, t) &= \frac{\sqrt{1+q}}{2\Gamma_{q^2}(1/2)} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} t^{\alpha-1} e_{\alpha,\alpha}(-\lambda\xi^2 t^\alpha; q) e(i\xi x; q^2) d_q \xi, \end{aligned}$$

where the variable of the q -integration ξ runs only over all $\xi \in \widetilde{\mathbb{R}}_q$ such that

$$|\lambda\xi^2 t^\alpha| < \frac{1}{(1-q)^\alpha}.$$

Consequently,

$$u(x, t) = \phi(x) * G(x, t).$$

Applying the q^2 -Fourier convolution formula gives

$$u(x, t) = \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} [T_y G(x, t)](x) \phi(y) d_q y,$$

where

$$[T_y G(x, t)](x) = \frac{\sqrt{1+q}}{2\Gamma_{q^2}(1/2)} \int_{-\infty/\sqrt{1-q}}^{\infty/\sqrt{1-q}} e(-iy\xi; q^2) t^{\alpha-1} e_{\alpha,\alpha}(-\lambda\xi^2 t^\alpha; q) e(ix\xi; q^2) d_q \xi. \quad \square$$

Competing interests

The author declares that they have no competing interests.

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