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Spectrum of the quadratic eigenparameter dependent discrete Dirac equations

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Abstract

Let us consider the Boundary Value Problem (BVP) for the discrete Dirac equations $a_{n+1}y_{n+1}^{(2)} + b_n y_n^{(2)} + p_n y_n^{(1)} = \lambda y_n^{(1)}, a_{n-1}y_{n-1}^{(1)} + b_n y_n^{(1)} + q_n y_n^{(2)} = \lambda y_n^{(2)}, n \in \mathbb{N},$ $(\gamma_0 + \gamma_1 \lambda + \gamma_2 \lambda^2)y_1^{(2)} + (\beta_0 + \beta_1 \lambda + \beta_2 \lambda^2)y_0^{(1)} = 0$, where $(a_n), (b_n), (p_n)$ and $(q_n), n \in \mathbb{N}$ are complex sequences, $\gamma_i, \beta_i \in \mathbb{C}, i = 0, 1, 2, \text{ and } \lambda$ is an eigenparameter. Discussing the eigenvalues and the spectral singularities, we prove that this BVP has a finite number of eigenvalues and spectral singularities with a finite number of multiplicities, if $\sum_{n=1}^{\infty} \exp(\varepsilon n^{\delta})(|1 - a_n| + |1 + b_n| + |p_n| + |q_n|) < \infty$, holds, for some $\varepsilon > 0$ and $\frac{1}{2} \le \delta \le 1$. **MSC:** 34L40; 47A10; 47A75

Keywords: discrete Dirac equations; eigenparameter; spectral analysis; discrete spectrum; spectral singularities

1 Introduction

In the last decade, discrete boundary value problems have been intensively studied. The spectral analysis of the difference equations have been treated by various authors in connection with the classical moment problem (see the monographs of Agarwal [1], Agarwal-Wong [2] and Kelley-Peterson [3] and the papers of Agarwal *et al.* [4, 5] and the references therein). Moreover, the modeling of certain linear and nonlinear problems in economies, optimal control theory, and other areas of study have led to the rapid development of the theory of difference equations. The spectral theory of the difference equations has also been applied to the solution of classes of nonlinear discrete Korteweg-de Vries equations and Toda lattices [6, 7].

Let us is consider the discrete boundary value problem (BVP)

$$a_{n-1}y_{n-1} + b_n y_n + a_n y_{n+1} = \lambda y_n, \quad n \in \mathbb{N} = \{1, 2, \ldots\},$$
(1.1)

$$y_0 = 0,$$
 (1.2)

where (a_n) and (b_n) are complex sequences, $a_0 \neq 0$ and λ is a spectral parameter. The spectral analysis of the BVP (1.1), (1.2) with continuous and point spectrum has been studied in [8]. In this article, the authors proved that the spectrum of the BVP (1.1), (1.2) consists of a continuous spectrum, eigenvalues and a spectral singularities. The spectral singularities are poles of the resolvent and are also imbedded in the continuous spectrum, but they are not eigenvalues. The effect of the spectral singularities in the spectral expansion of the BVP (1.1), (1.2) in terms of the principal vectors has been investigated in [9]. In [10,



©2014 Koprubasi; licensee Springer. This is an Open Access article distributed under the terms of the Creative Commons Attribution License (http://creativecommons.org/licenses/by/2.0), which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited. 11], the dependence of the structure of the spectral singularities of the BVP (1.1), (1.2) on the behavior of the sequences (a_n) and (b_n) was considered. Some problems related to the spectral analysis of difference equations with spectral singularities have been discussed in [12–15]. The spectral analysis of an eigenparameter dependent non-selfadjoint BVP for the system of difference equations of first order has been studied in [16, 17].

Let us consider the non-selfadjoint BVP for the system of difference equations of first order

$$\begin{cases} a_{n+1}y_{n+1}^{(2)} + b_n y_n^{(2)} + p_n y_n^{(1)} = \lambda y_n^{(1)}, \\ a_{n-1}y_{n-1}^{(1)} + b_n y_n^{(1)} + q_n y_n^{(2)} = \lambda y_n^{(2)}, \quad n \in \mathbb{N}, \end{cases}$$
(1.3)

$$(\gamma_0 + \gamma_1 \lambda + \gamma_2 \lambda^2) y_1^{(2)} + (\beta_0 + \beta_1 \lambda + \beta_2 \lambda^2) y_0^{(1)} = 0,$$
(1.4)

where $\binom{y_n^{(1)}}{y_n^{(2)}}$, $n \in \mathbb{N}$ are vector sequences, $a_n \neq 0$, $b_n \neq 0$ for all n. Also $\gamma_2 \neq \frac{-\beta_1}{a_0}$, $|\gamma_2| + |\beta_2| \neq 0$, and $\gamma_0 \beta_1 - \gamma_1 \beta_0 \neq 0$ where $\gamma_i, \beta_i \in \mathbb{C}$, i = 0, 1, 2. If $a_n \equiv 1$ and $b_n \equiv -1$ for all $n \in \mathbb{N}$, then the system (1.3) reduces to

$$\begin{cases} \Delta y_n^{(2)} + p_n y_n^{(1)} = \lambda y_n^{(1)}, \\ -\Delta y_{n-1}^{(1)} + q_n y_n^{(2)} = \lambda y_n^{(2)}, \quad n \in \mathbb{N}, \end{cases}$$
(1.5)

where Δ is a forward difference operator. The system (1.5) is the discrete analog of the well-known Dirac system

$$\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} y_1' \\ y_2' \end{pmatrix} + \begin{pmatrix} p(x) & 0 \\ 0 & q(x) \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \lambda \begin{pmatrix} y_1 \\ y_2 \end{pmatrix}$$

(see [18], Chapter 2). Therefore the system (1.5) (also (1.3)) is called a discrete Dirac system. The specific feature of this paper is the presence of the spectral parameter not only in the difference equation and also in the boundary condition, which is quadratic.

In this article, we intend to investigate eigenvalues and spectral singularities of the BVP (1.3), (1.4), which has a finite number of eigenvalues and spectral singularities with a finite number of multiplicities, if the condition

$$\sum_{n=1}^{\infty} \exp\left(\varepsilon n^{\delta}\right) \left(\left|1-a_{n}\right|+\left|1+b_{n}\right|+\left|p_{n}\right|+\left|q_{n}\right|\right) < \infty$$

holds, for some $\varepsilon > 0$ and $\frac{1}{2} \le \delta \le 1$.

2 Jost solution of (1.3)

Suppose that the condition

$$\sum_{n=1}^{\infty} \exp(\varepsilon n^{\delta}) (|1-a_n| + |1+b_n| + |p_n| + |q_n|) < \infty$$
(2.1)

is satisfied for some $\varepsilon > 0$ and $\frac{1}{2} \le \delta \le 1$. It is well known [15] that (1.3) has the bounded solution

$$f_n(z) = \begin{pmatrix} f_n^{(1)}(z) \\ f_n^{(2)}(z) \end{pmatrix} = \alpha_n \left(I_2 + \sum_{m=1}^{\infty} A_{nm} e^{imz} \right) \begin{pmatrix} e^{i\frac{z}{2}} \\ -i \end{pmatrix} e^{inz}, \quad n \in \mathbb{N},$$
(2.2)

$$f_0^{(1)}(z) = \alpha_0^{11} \left\{ e^{i\frac{z}{2}} \left[1 + \sum_{m=1}^{\infty} A_{0m}^{11} e^{imz} \right] - i \sum_{m=1}^{\infty} A_{0m}^{12} e^{imz} \right\},$$
(2.3)

under the condition (2.1) for $\lambda = 2 \sin \frac{z}{2}$ and $z \in \overline{\mathbb{C}}_+ := \{z : z \in \mathbb{C}, \text{Im } z \ge 0\}$, where

$$\alpha_n = \begin{pmatrix} \alpha_n^{11} & \alpha_n^{12} \\ \alpha_n^{21} & \alpha_n^{22} \end{pmatrix}, \qquad I_2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \qquad A_{nm} = \begin{pmatrix} A_{nm}^{11} & A_{nm}^{12} \\ A_{nm}^{21} & A_{nm}^{22} \end{pmatrix}.$$

Note that α_n^{ij} and A_{nm}^{ij} (i, j = 1, 2) are expressed in terms of (a_n) , (b_n) , (p_n) , and (q_n) , $n \in \mathbb{N}$. Also

$$\left|A_{nm}^{ij}\right| \le C \sum_{k=n+[\lfloor \frac{m}{2} \rfloor]}^{\infty} \left(|1-a_k|+|1+b_k|+|p_k|+|q_k|\right)$$
(2.4)

holds, where C > 0 is a constant and $[|\frac{m}{2}|]$ is the integer part of $\frac{m}{2}$. Therefore f_n is a vectorvalued analytic function with respect to z in $\mathbb{C}_+ := \{z : z \in \mathbb{C}, \operatorname{Im} z > 0\}$ and continuous in $\overline{\mathbb{C}}_+$ [15]. The solution $f(z) = (f_n(z)) = \begin{pmatrix} f_n^{(1)}(z) \\ f_n^{(2)}(z) \end{pmatrix}$ is called the Jost solution of (1.3).

3 Eigenvalues and spectral singularities of (1.3)-(1.4)

We define

$$f(z) = \left(\gamma_0 + 2\gamma_1 \sin\frac{z}{2} + 4\gamma_2 \sin^2\frac{z}{2}\right) f_1^{(2)}(z) + \left(\beta_0 + 2\beta_1 \sin\frac{z}{2} + 4\beta_2 \sin^2\frac{z}{2}\right) f_0^{(1)}(z),$$
(3.1)

using (1.4), (2.2), and (2.3). It follows from (2.2) and (2.3) that the function f is analytic in \mathbb{C}_+ , continuous up to the real axis and

$$f(z+4\pi)=f(z).$$

Let

$$T_0 := \{ z : z \in \mathbb{C}, z = x + iy, -2\pi \le x \le 2\pi, y > 0 \}, \qquad T := T_0 \cup [-2\pi, 2\pi].$$

We will denote the set of all eigenvalues and spectral singularities of the BVP (1.3), (1.4) by σ_d and σ_{ss} , respectively. It is clear that

$$\sigma_d = \left\{ \lambda : \lambda = 2 \sin \frac{z}{2}, z \in T_0, f(z) = 0 \right\},\tag{3.2}$$

$$\sigma_{ss} = \left\{ \lambda : \lambda = 2 \sin \frac{z}{2}, z \in [-2\pi, 2\pi], f(z) = 0 \right\}.$$
(3.3)

From (2.2), (2.3), and (3.1) we obtain

$$\begin{split} f(z) &= -\alpha_0^{11}\beta_2 e^{-i\frac{z}{2}} + i \big(\alpha_1^{22}\gamma_2 + \alpha_0^{11}\beta_1 \big) \\ &+ \big[\alpha_1^{22}\gamma_1 - \alpha_1^{21}\gamma_2 + \alpha_0^{11}(\beta_0 + 2\beta_2) \big] e^{i\frac{z}{2}} \end{split}$$

$$\begin{aligned} &+i\left[-\alpha_{1}^{22}(\gamma_{0}+2\gamma_{2})+\alpha_{1}^{21}\gamma_{1}-\alpha_{0}^{11}\beta_{1}\right]e^{iz} \\ &+\left[-\alpha_{1}^{22}\gamma_{1}+\alpha_{1}^{21}(\gamma_{0}+2\gamma_{2})-\alpha_{0}^{11}\beta_{2}\right]e^{i\frac{3z}{2}} \\ &+i\left(\alpha_{1}^{22}\gamma_{2}-\alpha_{1}^{21}\gamma_{1}\right)e^{2iz}-\alpha_{1}^{21}\gamma_{2}e^{i\frac{5z}{2}} \\ &+i\sum_{m=1}^{\infty}\alpha_{0}^{11}\beta_{2}A_{0m}^{12}e^{i(m-1)z} +\sum_{m=1}^{\infty}\alpha_{0}^{11}\left(\beta_{1}A_{0m}^{12}-\beta_{2}A_{0m}^{11}\right)e^{i(m-\frac{1}{2})z} \\ &+i\sum_{m=1}^{\infty}\left\{\gamma_{2}\left(\alpha_{1}^{22}A_{1m}^{22}+\alpha_{1}^{21}A_{1m}^{12}\right) \\ &+\alpha_{0}^{11}\left[\beta_{1}A_{0m}^{11}-\left(\beta_{0}+2\beta_{2}\right)A_{0m}^{12}\right]\right\}e^{imz} \\ &+\sum_{m=1}^{\infty}\left\{\alpha_{1}^{22}\left(\gamma_{1}A_{1m}^{22}-\gamma_{2}A_{1m}^{21}\right)+\alpha_{1}^{21}\left(\gamma_{1}A_{1m}^{12}-\gamma_{2}A_{1m}^{11}\right) \\ &+\alpha_{0}^{11}\left[\left(\beta_{0}+2\beta_{2}\right)A_{0m}^{11}-\beta_{1}A_{0m}^{12}\right]\right]e^{i(m+\frac{1}{2})z} \\ &+i\sum_{m=1}^{\infty}\left\{-\alpha_{1}^{22}\left[\left(\gamma_{0}+2\gamma_{2}\right)A_{1m}^{22}-\gamma_{1}A_{1m}^{21}\right] \\ &-\alpha_{1}^{21}\left[\left(\gamma_{0}+2\gamma_{2}\right)A_{1m}^{12}-\gamma_{1}A_{1m}^{12}\right] \\ &+\alpha_{0}^{11}\left(\beta_{2}A_{0m}^{22}-\beta_{1}A_{0m}^{11}\right)\right]e^{i(m+1)z} \\ &+\sum_{m=1}^{\infty}\left\{-\alpha_{1}^{22}\left[\left(\gamma_{2}A_{2m}^{22}-\gamma_{1}A_{1m}^{21}\right)+\alpha_{1}^{21}\left(\gamma_{2}A_{1m}^{12}-\gamma_{1}A_{1m}^{11}\right)\right]e^{i(m+2)z} \\ &+i\sum_{m=1}^{\infty}\left\{-\alpha_{1}^{22}\left[\gamma_{1}A_{1m}^{22}-\gamma_{1}A_{1m}^{21}\right] \\ &-\alpha_{0}^{11}\beta_{2}A_{0m}^{11}\right]e^{i(m+\frac{3}{2})z} \\ &+i\sum_{m=1}^{\infty}\left[\alpha_{1}^{22}\left(\gamma_{2}A_{1m}^{22}-\gamma_{1}A_{1m}^{21}\right)+\alpha_{1}^{21}\left(\gamma_{2}A_{1m}^{12}-\gamma_{1}A_{1m}^{11}\right)\right]e^{i(m+2)z} \\ &+\sum_{m=1}^{\infty}\left[-\gamma_{2}\left(\alpha_{1}^{22}A_{1m}^{21}+\alpha_{1}^{21}\gamma_{2}A_{1m}^{21}\right)e^{i(m+\frac{5}{2})z}. \end{aligned}$$

$$(3.4)$$

Let

$$F(z) := f(z)e^{i\frac{z}{2}},$$
(3.5)

then the function F is analytic in $\mathbb{C}_+,$ continuous in $\overline{\mathbb{C}}_+,$

$$F(z) = -\alpha_0^{11} \beta_2 + i (\alpha_1^{22} \gamma_2 + \alpha_0^{11} \beta_1) e^{i\frac{z}{2}} + [\alpha_1^{22} \gamma_1 - \alpha_1^{21} \gamma_2 + \alpha_0^{11} (\beta_0 + 2\beta_2)] e^{iz} + i [-\alpha_1^{22} (\gamma_0 + 2\gamma_2) + \alpha_1^{21} \gamma_1 - \alpha_0^{11} \beta_1] e^{i\frac{3z}{2}} + [-\alpha_1^{22} \gamma_1 + \alpha_1^{21} (\gamma_0 + 2\gamma_2) - \alpha_0^{11} \beta_2] e^{2iz} + i (\alpha_1^{22} \gamma_2 - \alpha_1^{21} \gamma_1) e^{i\frac{5z}{2}} - \alpha_1^{21} \gamma_2 e^{3iz} + i \sum_{m=1}^{\infty} \alpha_0^{11} \beta_2 A_{0m}^{12} e^{i(m-\frac{1}{2})z} + \sum_{m=1}^{\infty} \alpha_0^{11} (\beta_1 A_{0m}^{12} - \beta_2 A_{0m}^{11}) e^{imz}$$

and

$$F(z+4\pi)=F(z).$$

Using (3.2)-(3.5),

$$\sigma_d = \left\{ \lambda : \lambda = 2 \sin \frac{z}{2}, z \in T_0, F(z) = 0 \right\},\tag{3.7}$$

$$\sigma_{ss} = \left\{ \lambda : \lambda = 2 \sin \frac{z}{2}, z \in [-2\pi, 2\pi], F(z) = 0 \right\}.$$
(3.8)

Definition 3.1 The multiplicity of a zero of F in T is called the multiplicity of the corresponding eigenvalue or spectral singularity of the BVP (1.3), (1.4).

It follows from (3.2) and (3.3) that, in order to investigate the quantitative properties of the eigenvalues and the spectral singularities of the BVP (1.3), (1.4), we need to discuss the quantitative properties of the zeros of F in T.

Let

$$S_{1} := \{ z : z \in T_{0}, F(z) = 0 \},$$

$$S_{2} := \{ z : z \in [-2\pi, 2\pi], F(z) = 0 \}.$$
(3.9)

We also denote the set of all limit points of S_1 by S_3 and the set of all zeros of F with infinite multiplicity by S_4 .

From (3.2), (3.3), and (3.9) we get

$$\sigma_{d} = \left\{ \lambda : \lambda = 2 \sin \frac{z}{2}, z \in S_{1} \right\},$$

$$\sigma_{ss} = \left\{ \lambda : \lambda = 2 \sin \frac{z}{2}, z \in S_{2} \right\}.$$
(3.10)

Theorem 3.1 If (2.1) holds, then:

- (i) The set S_1 is bounded and countable.
- (ii) $S_1 \cap S_3 = \emptyset$, $S_1 \cap S_4 = \emptyset$.
- (iii) The set S_2 is compact and $\mu(S_2) = 0$, where μ denotes the Lebesgue measure in the real axis.
- (iv) $S_3 \subset S_2, S_4 \subset S_2; \mu(S_3) = \mu(S_4) = 0.$ (v) $S_3 \subset S_4.$

Proof Using (1.4), (2.4), and (3.6), we have

$$F(z) = \begin{cases} -\alpha_0^{11}\beta_2 + o(e^{-y}), & \beta_2 \neq 0, z \in T, y \to \infty, \\ i(\alpha_1^{22}\gamma_2 + \alpha_0^{11}\beta_1)e^{i\frac{z}{2}} + o(e^{-y}), & \beta_2 = 0, z \in T, y \to \infty. \end{cases}$$
(3.11)

Equation (3.11) shows that S_1 is bounded. Since F is analytic in \mathbb{C}_+ and is a 4π periodic function, we find that S_1 has at most a countable number of elements. This proves (i).

From the uniqueness theorems of analytic functions we obtain (ii)-(iv) [19]. Using the continuity of all derivatives of *F* on $[-2\pi, 2\pi]$ we get (v).

From (3.10) and Theorem 3.1, we have the following.

Theorem 3.2 Under the condition (2.1):

- (i) The set of eigenvalues of the BVP (1.3), (1.4) is bounded and countable and its limit points can lie only in [-2, 2].
- (ii) $\sigma_{ss} \subset [-2, 2]$, $\sigma_{ss} = \overline{\sigma}_{ss}$ and $\mu(\sigma_{ss}) = 0$.

For $\delta = 1$ the condition (2.1) reduces to

$$\sum_{n=1}^{\infty} \exp(\varepsilon n) \left(|1 - a_n| + |1 + b_n| + |p_n| + |q_n| \right) < \infty.$$
(3.12)

Theorem 3.3 Under the condition (3.12) the BVP (1.3), (1.4) has a finite number of eigenvalues and spectral singularities and each of them is of finite multiplicity.

Proof Using (2.4) we find that

$$\left|A_{nm}^{ij}\right| \le C \exp\left[-\frac{\varepsilon}{4}(n+m)\right], \quad i,j=1,2,n,m\in\mathbb{N},$$
(3.13)

where C > 0 is a constant. From (3.6) and (3.13) we observe that the function F has an analytic continuation to the half-plane Im $z > -\frac{\varepsilon}{4}$. Since F is a 4π periodic function, the limit points its zeros in T cannot lie in $[-2\pi, 2\pi]$. Using Theorem 3.1 we have the bounded

It is seen that the condition (3.12) guarantees the analytic continuation of F from the real axis to lower half-plane. So the finiteness of eigenvalues and spectral singularities of the BVP (1.3), (1.4) are obtained as a result of this analytic continuation.

Now let us suppose that

$$\sum_{n=1}^{\infty} \exp(\varepsilon n^{\delta}) \left(|1 - a_n| + |1 + b_n| + |p_n| + |q_n| \right) < \infty, \quad \varepsilon > 0, \frac{1}{2} \le \delta < 1,$$
(3.14)

which is weaker than (3.12). It is evident that under the condition (3.14) the function F is analytic in \mathbb{C}_+ and infinitely differentiable on the real axis. But F does not have an analytic continuation from the real axis to lower half-plane. Therefore under the condition (3.14) the finiteness of eigenvalues and spectral singularities of the BVP (1.3), (1.4) cannot be shown in a way similar to Theorem 3.3.

Under the condition (3.14), to prove that the eigenvalues and the spectral singularities of the BVP (1.3), (1.4) are of finite number we will use the following.

Theorem 3.4 ([14]) Let us assume that the 4π periodic function g is analytic in \mathbb{C}_+ , all of its derivatives are continuous in $\overline{\mathbb{C}}_+$ and

$$\sup_{z\in T} \left| g^{(k)}(z) \right| \le A_k, \quad k\in \mathbb{N}\cup\{0\}.$$

If the set $G \subset [-2\pi, 2\pi]$ with Lebesgue measure zero is the set of all zeros the function g with infinite multiplicity in T, if

$$\int_0^\infty \ln K(s) \, d\mu(G_s) = -\infty,\tag{3.15}$$

where $K(s) = \inf_k \frac{A_k s^k}{k!}$ and $\mu(G_s)$ is the Lebesgue measure of s-neighborhood of G and $\omega \in (-2\pi, 2\pi)$ is an arbitrary constant, then $g \equiv 0$ in $\overline{\mathbb{C}}_+$.

Under the condition (3.14) from (2.4) and (3.6) we find

$$\left|F^{(k)}(z)\right| \leq A_k, \quad k \in \mathbb{N} \cup \{0\},$$

where

$$A_k = 2^k C \sum_{m=1}^{\infty} m^k \exp(-\varepsilon m^{\delta})$$

and C > 0 is a constant. We can obtain the following estimate:

$$A_k \le 2^k C \int_0^\infty x^k \exp\left(-\varepsilon x^\delta\right) dx \le D d^k k! k^{k\frac{1-\delta}{\delta}},$$
(3.16)

where *D* and *d* are constants depending on *C*, ε , and δ .

Theorem 3.5 If (3.14) holds, then $S_4 = \emptyset$.

Proof The function F satisfies all conditions of Theorem 3.4 except (3.15). But F is not identically equal to zero. In this case the function F satisfies the condition

$$\int_0^{\omega} \ln K(s) \, d\mu(S_{4,s}) > -\infty \tag{3.17}$$

instead of (3.15), where $K(s) = \inf_k \frac{A_k s^k}{k!}$, $k \in \mathbb{N} \cup \{0\}$, and $\mu(S_{4,s})$ is the Lebesgue measure of the *s*-neighborhood of S_4 and A_k defined by (3.16). Substituting (3.16) in the definition of K(s), we get

$$K(s) = D \exp\left\{-\frac{1-\delta}{\delta}e^{-1}d^{-\frac{\delta}{1-\delta}}s^{-\frac{\delta}{1-\delta}}\right\}.$$
(3.18)

It follows from (3.17) and (3.18) that

$$\int_0^{\omega} s^{-\frac{\delta}{1-\delta}} d\mu(S_{4,s}) < \infty.$$
(3.19)

Since $\frac{\delta}{1-\delta} \ge 1$, (3.19) holds for arbitrary *s* if and only if $\mu(S_{4,s}) = 0$ or $S_4 = \emptyset$.

Theorem 3.6 Under the condition (3.14) the BVP (1.3), (1.4) has a finite number of eigenvalues and spectral singularities and each of them is of finite multiplicity.

Proof To be able to prove the theorem we have to show that the function F has a finite number of zeros with finite multiplicities in T.

From Theorem 3.1 and Theorem 3.5 we see that $S_3 = \emptyset$. So the bounded sets S_1 and S_2 have no limit points, *i.e.*, the function *F* has only a finite number of zeros in *T*. Since $S_4 = \emptyset$, these zeros are of finite multiplicity.

Competing interests

The author declares that they have no competing interests.

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