# Two dimensional determination of source terms in linear parabolic equation from the final overdetermination 

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#### Abstract

A case of steady-case heat flow through a plane wall, which can be formulated as $u_{t}(x, y, t)-\operatorname{div}(k(x, y) \nabla u(x, y, t))=F(x, y, t)$ with Robin boundary condition $-k(1, y) u_{x}(1, y, t)=v_{1}\left[u(1, y, t)-T_{0}(t)\right],-k(x, 1) u_{y}(x, 1, t)=v_{2}\left[u(x, 1, t)-T_{1}(t)\right]$, where $\omega:=\left\{F(x, y, t) ; T_{0}(t) ; T_{1}(t)\right\}$ is to be determined, from the measured final data $\mu_{T}(x, y)=u(x, y, T)$ is investigated. It is proved that the Fréchet gradient of the cost functional $J(\omega)=\left\|\mu_{T}(x, y)-u(x, y, T ; \omega)\right\|^{2}$ can be found via the solution of the adjoint parabolic problem. Lipschitz continuity of the gradient is derived. The obtained results permit one to prove the existence of a quasi-solution of the inverse problem. A steepest descent method with line search, which produces a monotone iteration scheme based on the gradient, is formulated. Some convergence results are given.


Keywords: two dimensional determination; cost functional; steepest descent method

## 1 Introduction

Consider the one dimensional physical system in Figure 1, where the left of the solid is full of hot gas. Whenever a temperature gradient exists in the solid medium, heat will flow from the higher-temperature region to the lower-temperature region. According to Fourier's law, for a homogeneous, isotropic solid, the following equation holds:

$$
\begin{equation*}
q(x, t)=-k u_{x x}(x, t), \tag{1.1}
\end{equation*}
$$

where $q(x, t)$ represents heat flow per unit time, per unit area of the isothermal surface in the direction of decreasing temperature, $u(x, t)$ is the temperature distribution in the solid, and $k$ is called the thermal conductivity.

However, in practice, the thermal conductivity may depend on $x$, namely, $k:=k(x)$. Besides, there may be a heat source $g(x, t)$ in the solid. Under these conditions, the physical system can be formulated as

$$
\begin{cases}u_{t}=\left(k(x) u_{x}(x, t)\right)_{x}+g(x, t), & (x, t) \in \tilde{\Omega}_{T},  \tag{1.2}\\ u(x, 0)=\mu_{0}(x), & x \in\left(0, L_{1}\right), \\ u_{x}(0, t)=0, \quad-k\left(L_{1}\right) u_{x}\left(L_{1}, t\right)=\sigma\left[u\left(L_{1}, t\right)-T_{0}(t)\right], & t \in(0, T],\end{cases}
$$

Figure 1 The one dimensional physical model.

where $\tilde{\Omega}_{T}:=\left\{0<x<L_{1}, 0<t \leq T\right\}$, and $T_{0}(t)$ is the temperature of the cold gas on the right of the solid. $-k\left(L_{1}\right) u_{x}\left(L_{1}, t\right)=\sigma\left[u\left(L_{1}, t\right)-T_{0}(t)\right]$ represents the convection at $x=L_{1}$ according to Newton's law, and the constant $\sigma$ is called the convection coefficient.
It is desired to find the pair $\omega:=\left\{g(x, t) ; T_{0}(t)\right\}$ from the final state observation

$$
\begin{equation*}
\mu_{T}(x)=u(x, T) \tag{1.3}
\end{equation*}
$$

The mathematical model (1.2) can also arise in hydrology [1], material sciences [2], and transport problems [3]. In [4] a tsunami model based on shallow water theory is studied. The authors treat the inverse problem of determining an unknown initial tsunami source $q(x, y)$ by using measurements $f_{m}(t)$ of the height of a passing tsunami wave at a finite number of given points $\left(x_{m}, y_{m}\right), m=1,2, \ldots, M$, of the coastal area. These are nonlinear inverse problems, and it is well known that they are generally ill posed, i.e. the existence, uniqueness, and stability of their solutions are not always guaranteed [5]. There are many contributions for the linear parabolic equations with final overdetermination (see, for instance [6-12]). The time-dependent heat source $H(t)$ of the separable sources of the form $g(x, t)=F(x) H(t)$ is investigated in [13]. For $g(x, t)=p(x) u$, in [14], the author proved the existence and uniqueness of $p(x)$, and the local well-posedness of the inverse problem was discussed in [15]. The simultaneous reconstruction of the initial temperature and heat radiative coefficient was investigated in [16] by using the measurement of temperature given at a fixed time and the measurement of the temperature in a subregion of the physical domain. A determination of the unknown function $p(x)$ in the source term $g=p(x) f(u)$, via fixed point theory, was given in [17]. Based on the optimal control framework, [18] considered the determination of a pair $(p, u)$ in the nonlinear parabolic equation

$$
u_{t}-u_{x x}+p(x) f(u)=0,
$$

with initial and homogeneous Dirichlet boundary conditions:

$$
u(x, 0)=\phi(x),\left.\quad u\right|_{x=0}=\left.u\right|_{x=L_{1}}=0,
$$

from the overspecified data $u(x, T)=\mu_{T}(x)$. The local uniqueness and stability of the solution were proved. In [19], a weak solution approach for the pair $\omega(x, t):=\left\{g(x, t) ; T_{0}(t)\right\}$ was given via a steepest descent method on minimizing the cost functional

$$
\begin{equation*}
J(\omega)=\left\|\mu_{T}(x)-u(x, T ; \omega)\right\|_{L^{2}\left[0, L_{1}\right]}^{2} . \tag{1.4}
\end{equation*}
$$

Figure 2 The two dimensional physical model.


In this contribution, we consider the corresponding two dimensional problem (see Figure 2) as follows:

$$
\begin{cases}u_{t}(x, y, t)-\operatorname{div}(k(x, y) \nabla u(x, y, t))=F(x, y, t), & \text { in } \Omega_{T},  \tag{1.5}\\ u_{x}(0, y, t)=u_{y}(x, 0, t)=0, & x, y \in(0,1), t \in(0, T], \\ -k(1, y) u_{x}(1, y, t)=v_{1}\left[u(1, y, t)-T_{0}(t)\right], & x, y \in(0,1), t \in(0, T], \\ -k(x, 1) u_{y}(x, 1, t)=v_{2}\left[u(x, 1, t)-T_{1}(t)\right], & x, y \in(0,1), t \in(0, T], \\ u(x, y, 0)=u_{0}(x, y), & (x, y) \in \Omega,\end{cases}
$$

where $\Omega_{T}=\Omega \times(0, T)=(0,1) \times(0,1) \times(0, T)$ and $\nu_{1}, \nu_{2}>0$. The inverse problem here is to determine $\omega:=\left\{F(x, y, t) ; T_{0}(t) ; T_{1}(t)\right\}$ from the final state observation (the overspecified data)

$$
\begin{equation*}
u_{T}(x, y)=u(x, y, T) \tag{1.6}
\end{equation*}
$$

In this work, based on a weak solution approach we will show how the inverse problem can be formulated and solved for $\omega:=\left\{F(x, y, t) ; T_{0}(t) ; T_{1}(t)\right\}$. Moreover, we will prove that the gradient $J^{\prime}(w)$ of the cost (auxiliary) functional

$$
J(\omega)=\left\|\mu_{T}(x, y)-u(x, y, T ; \omega)\right\|^{2}
$$

can be expressed via the solution $\varphi=\varphi(x, y, t ; \omega)$ of the appropriate adjoint problem.
This paper is organized as follows. In Section 2, we give an analysis of the two dimensional problem and prove the Fréchet-differentiability of $J(\omega)$. In Section 3, we present the framework of a steepest descent iterate with line search of the two dimensional inverse problem, where $J^{\prime}(\omega)$ can be found via an adjoint parabolic problem. The convergence of the sequence is analyzed in Section 4. Conclusions are stated in Section 5.

## 2 The analysis of the two dimensional problem

The direct problem (1.5) is to get the solution from a given pair $w$. Firstly, we define $W:=$ $\mathcal{F} \times \mathcal{T}_{0} \times \mathcal{T}_{1}$, the set of admissible unknown sources $F(x, y, t), T_{0}(t), T_{1}(t)$ with

$$
\begin{align*}
& F(x, y, t) \in L^{2}\left(\Omega_{T}\right), \quad T_{0}(t), T_{1}(t) \in L^{2}[0, T]  \tag{2.1}\\
& 0<T_{0 *} \leq T_{0}(t), T_{1}(t) \leq T_{0}^{*}<+\infty \\
& k(x, y)>0, \quad k(x, y) \in L^{\infty}(\Omega), \quad u_{0}(x, y), u_{T}(x, y) \in L^{2}(\Omega) . \tag{2.2}
\end{align*}
$$

It is obvious that the set $W$ is a closed and convex subset in $L^{2}\left(\Omega_{T}\right) \times L^{2}[0, T] \times L^{2}[0, T]$. The scalar product in $W$ is defined as

$$
\begin{aligned}
\left(\omega_{1}, \omega_{2}\right)_{W}:= & \int_{\Omega_{T}} F_{1}(x, y, t) F_{2}(x, y, t) d x d y d t+\int_{0}^{T} T_{0}^{(1)}(t) T_{0}^{(2)}(t) d t \\
& +\int_{0}^{T} T_{1}^{(1)}(t) T_{1}^{(2)}(t) d t, \quad \forall \omega_{1}, \omega_{2} \in W,
\end{aligned}
$$

where $\omega_{m}:=\left\{F_{m} ; T_{0}^{(m)} ; T_{1}^{(m)}\right\}, m=1,2$. We can prove that this kind of problem has a quasisolution $u \in H^{1,0}\left(\Omega_{T}\right)$ satisfying the identity

$$
\begin{align*}
& \int_{\Omega} u(x, y, T) v(x, y, T) d x d y-\int_{\Omega} u_{0}(x, y) v(x, y, 0) d x d y \\
&-\int_{\Omega_{T}}\left(u v_{t}-k u_{x} v_{x}-k u_{y} v_{y}\right) d x d y d t \\
&= v_{1} \int_{0}^{T} \int_{0}^{1}\left[u(1, y, t)-T_{0}(t)\right] v(1, y, t) d y d t \\
&+v_{2} \int_{0}^{T} \int_{0}^{1}\left[u(x, 1, t)-T_{1}(t)\right] v(x, 1, t) d x d t+\int_{0}^{T} \int_{\Omega} F(x, y, t) v(x, y, t) d x d t \tag{2.3}
\end{align*}
$$

for all $v \in H^{1,0}\left(\Omega_{T}\right)$. Here $H^{1,0}\left(\Omega_{T}\right)$ is the Sobolev space with the norm

$$
\|u\|_{1}:=\left\{\int_{\Omega_{T}}\left[u^{2}+u_{x}^{2}+u_{y}^{2}\right] d x d y d t\right\}^{1 / 2}
$$

It is also known that, under conditions (2.1) and (2.2), the weak solution $u(x, y, t) \in$ $H^{1,0}\left(\Omega_{T}\right)$ of the direct problem (1.5) exists and is unique [20].

### 2.1 Method discussion

To solve the inverse problem (1.5)-(1.6), we introduce a cost functional

$$
J(\omega)=\int_{\Omega}\left[u(x, y, T ; \omega)-u_{T}(x, y)\right]^{2} d x d y
$$

We are going to give an iterative solution to this kind of problem. To begin, we study the derivative of the cost functional. Let us consider the first variation of the cost functional:

$$
\begin{align*}
\Delta J(\omega): & =J(\omega+\Delta \omega)-J(\omega) \\
= & \int_{\Omega}\left[u(x, y, T ; \omega+\Delta \omega)-u_{T}(x, y)\right]^{2} d x d y-\int_{\Omega}\left[u(x, y, T ; \omega)-u_{T}(x, y)\right]^{2} d x d y \\
= & \int_{\Omega} \Delta u(x, y, T ; \omega)\left(\Delta u(x, y, T ; \omega)+2\left(u(x, y, T ; \omega)-u_{T}(x, y)\right)\right) d x d y \\
= & 2 \int_{\Omega}\left(u(x, y, T ; \omega)-u_{T}(x, y)\right) \Delta u(x, y, T ; \omega) d x d y \\
& +\int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y, \tag{2.4}
\end{align*}
$$

where

$$
\begin{aligned}
& \omega+\Delta \omega:=\left\{F+\Delta F ; T_{0}+\Delta T_{0} ; T_{1}+\Delta T_{1}\right\} \in W \\
& \Delta u(x, y, t ; \omega):=u(x, y, t ; \omega+\Delta \omega)-u(x, y, t ; \omega) \in H^{1,0}\left(\Omega_{T}\right) .
\end{aligned}
$$

Furthermore, the function $\Delta u:=\Delta u(x, y, t ; \omega)$ is the solution of the following parabolic problem:

$$
\begin{cases}\Delta u_{t}-\operatorname{div}(k(x, y) \nabla \Delta u)=\Delta F(x, y, t), & \text { in } \Omega_{T},  \tag{2.5}\\ \Delta u_{x}(0, y, t)=\Delta u_{y}(x, 0, t)=0, & x, y \in(0,1), t \in(0, T], \\ -k(1, y) \Delta u_{x}(1, y, t)=v_{1}\left(\Delta u(1, y, t)-\Delta T_{0}(t)\right), & x, y \in(0,1), t \in(0, T] \\ -k(x, 1) \Delta u_{y}(x, 1, t)=v_{2}\left(\Delta u(x, 1, t)-\Delta T_{1}(t)\right), & x, y \in(0,1), t \in(0, T], \\ \Delta u(x, y, 0)=0, & (x, y) \in \Omega .\end{cases}
$$

Now, we are ready to estimate the derivative of the cost functional $J(\omega)$. Firstly, we estimate the first term in (2.4), i.e. $2 \int_{\Omega}\left(u(x, y, T ; \omega)-u_{T}(x, y)\right) \Delta u(x, y, T ; \omega) d x d y$.

Lemma 2.1 Let $\omega, \omega+\Delta \omega \in W$ be given elements. If $u(x, y, t ; \omega)$ is the corresponding solution of the direct problem (1.5), and $\varphi(x, y, t) \in H^{1,0}\left(\Omega_{T}\right)$ is the solution of the backward parabolic problem

$$
\begin{cases}\varphi_{t}(x, y, t)+\operatorname{div}(k(x, y) \nabla \varphi(x, y, t))=0, & \text { in } \Omega_{T},  \tag{2.6}\\ \varphi_{x}(0, y, t)=\varphi_{y}(x, 0, t)=0, & x, y \in(0,1), t \in(0, T], \\ -k(1, y) \varphi_{x}(1, y, t)=v_{1} \varphi(1, y, t), & x, y \in(0,1), t \in(0, T], \\ -k(x, 1) \varphi_{y}(x, 1, t)=v_{2} \varphi(x, 1, t), & x, y \in(0,1), t \in(0, T], \\ \varphi(x, y, T)=u(x, y, T ; \omega)-u_{T}(x, y), & (x, y) \in \Omega,\end{cases}
$$

then, for all $\omega \in W$, the following integral identity holds:

$$
\begin{align*}
\int_{\Omega} & \left(u(x, y, T ; \omega)-u_{T}(x, y)\right) \Delta u(x, y, T ; \omega) d x d y \\
\quad= & v_{1} \int_{0}^{T} \Delta T_{0}(t) \int_{0}^{1} \varphi(1, y, t ; \omega) d y d t+v_{2} \int_{0}^{T} \Delta T_{1}(t) \int_{0}^{1} \varphi(x, 1, t ; \omega) d x d t \\
& +\int_{\Omega_{T}} \Delta F(x, y, t) \varphi(x, y, t) d x d y d t . \tag{2.7}
\end{align*}
$$

Proof Taking the final condition at $t=T$ in (2.6) and the boundary conditions in (2.6) and (2.5) into account, we can deduce that

$$
\begin{aligned}
\int_{\Omega} & \left(u(x, y, T ; \omega)-u_{T}(x, y)\right) \Delta u(x, y, T ; \omega) d x d y \\
& =\int_{\Omega} \varphi(x, y, T ; \omega) \Delta u(x, y, T ; \omega) d x d y \\
& =\int_{\Omega} \int_{0}^{T}(\varphi(x, y, t ; \omega) \Delta u(x, y, t ; \omega))_{t} d t d x d y \\
& =\int_{\Omega_{T}} \varphi_{t}(x, y, t ; \omega) \triangle u(x, y, t ; \omega)+\varphi(x, y, t ; \omega) \triangle u_{t}(x, y, t ; \omega) d x d y d t
\end{aligned}
$$

$$
\begin{aligned}
= & \int_{\Omega_{T}}-\operatorname{div}(k(x, y) \nabla \varphi(x, y, t)) \Delta u(x, y, t ; \omega) \\
& +\varphi(x, y, t ; \omega) \operatorname{div}(k(x, y) \nabla \Delta u) d x d y d t+\int_{\Omega_{T}} \varphi(x, y, t) \Delta F(x, y, t) d x d y d t \\
= & -\int_{0}^{T} \int_{0}^{1} k(1, y) \varphi_{x}(1, y, t) \Delta u(1, y, t ; \omega) d y d t \\
& -\int_{0}^{T} \int_{0}^{1} k(x, 1) \varphi_{y}(x, 1, t) \Delta u(x, 1, t ; \omega) d x d t \\
& +\int_{0}^{T} \int_{0}^{1} k(1, y) \Delta u_{x}(1, y, t) \varphi(1, y, t ; \omega) d y d t \\
& +\int_{0}^{T} \int_{0}^{1} k(x, 1) \Delta u_{y}(x, 1, t) \varphi(x, 1, t ; \omega) d x d t+\int_{\Omega_{T}} \varphi(x, y, t) \Delta F(x, y, t) d x d y d t \\
= & v_{1} \int_{0}^{T} \Delta T_{0}(t) \int_{0}^{1} \varphi(1, y, t ; \omega) d y d t+\nu_{2} \int_{0}^{T} \Delta T_{1}(t) \int_{0}^{1} \varphi(x, 1, t ; \omega) d x d t \\
& +\int_{\Omega_{T}} \Delta F(x, y, t) \varphi(x, y, t) d x d y d t .
\end{aligned}
$$

This completes the proof of Lemma 2.1.

Remark We define the parabolic problem (2.6) as an adjoint problem corresponding to the inverse problem (1.5)-(1.6). It is easy to see that the parabolic problem (2.6) is a backward one, and this problem is well posed.

Next, we show that the second term in (2.4), i.e. $\int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y$, is of order $O\left(\|\Delta \omega\|_{W}^{2}\right)$.

Lemma 2.2 Let $\Delta u=\Delta u(x, t ; \omega) \in H^{1,0}\left(\Omega_{T}\right)$ be the solution of the parabolic problem (2.5) with respect to a given $\omega \in W$. Then the following estimate holds:

$$
\int_{\Omega}(\Delta u(x, y, T))^{2} d x d y \leq \frac{\max \left\{v_{1}, \nu_{2}, 1\right\}}{\epsilon}\|\Delta \omega\|_{W}^{2}
$$

where

$$
\|\Delta \omega\|_{W}:=\left(\int_{0}^{T}\left(\Delta T_{0}(t)\right)^{2} d t+\int_{0}^{T}\left(\Delta T_{1}(t)\right)^{2} d t+\int_{\Omega_{T}}(\Delta F(x, y, t))^{2} d x d y d t\right)^{1 / 2}
$$

is the $H^{0}$-norm of the function $\Delta \omega \in W$, and the constant $\epsilon$ is defined as follows:

$$
k_{*}=\min _{x, y \in[0,1]} k(x, y)>0, \quad \epsilon=\min \left\{k_{*} ; \frac{2 v_{1}}{v_{1}+2}\right\} .
$$

Proof Multiplying $\triangle u$ on both sides of (2.5) and integrating on $\Omega_{T}$, we obtain

$$
\begin{aligned}
0 & =\int_{\Omega_{T}}\left(\Delta u_{t}-\operatorname{div}(k(x, y) \nabla \Delta u)-\Delta F(x, y, t)\right) \Delta u d x d y d t \\
& =\int_{\Omega_{T}}\left(\frac{1}{2} \Delta u^{2}\right)_{t} d x d y d t-\int_{\Omega_{T}}\left(k(x, y) \Delta u_{x} u\right)_{x} d x d y d t
\end{aligned}
$$

$$
\begin{aligned}
& -\int_{\Omega_{T}}\left(k(x, y) \Delta u_{y} u\right)_{y} d x d y d t+\int_{\Omega_{T}} k(x, y)\left(\left(\Delta u_{x}\right)^{2}+\left(\Delta u_{y}\right)^{2}\right) d x d y d t \\
& -\int_{\Omega_{T}} \Delta F(x, y, t) \Delta u(x, y, t) d x d y d t \\
= & \frac{1}{2} \int_{\Omega}(\Delta u(x, y, T))^{2} d x d y-\int_{0}^{T} \int_{0}^{1} k(1, y) \Delta u_{x}(1, y, t) \Delta u(1, y, t) d y d t \\
& -\int_{0}^{T} \int_{0}^{1} k(x, 1) \Delta u_{y}(x, 1, t) \Delta u(x, 1, t) d x d t-\int_{\Omega_{T}} \Delta F(x, y, t) \Delta u(x, y, t) d x d y d t \\
& +\int_{\Omega_{T}} k(x, y)\left(\left(\Delta u_{x}(x, y, t)\right)^{2}+\left(\Delta u_{y}(x, y, t)\right)^{2}\right) d x d y d t \\
= & \frac{1}{2} \int_{\Omega^{2}}(\Delta u(x, y, T))^{2} d x d y-\int_{\Omega_{T}} \Delta F(x, y, t) \Delta u(x, y, t) d x d y d t \\
& +v_{1} \int_{0}^{T} \int_{0}^{1}(\Delta u(1, y, t))^{2} d y d t+v_{2} \int_{0}^{T} \int_{0}^{1}(\Delta u(x, 1, t))^{2} d x d t \\
& -v_{1} \int_{0}^{T} \int_{0}^{1} \Delta T_{0}(t) \Delta u(1, y, t) d y d t-v_{2} \int_{0}^{T} \int_{0}^{1} \Delta T_{1}(t) \Delta u(x, 1, t) d x d t \\
& +\int_{\Omega_{T}} k(x, y)\left(\left(\Delta u_{x}(x, y, t)\right)^{2}+\left(\Delta u_{y}(x, y, t)\right)^{2}\right) d x d y d t ;
\end{aligned}
$$

here, the initial and boundary conditions are used.
This implies the following identity:

$$
\begin{align*}
& \frac{1}{2} \int_{\Omega}(\Delta u(x, y, T))^{2} d x d y \\
& \quad+v_{1} \int_{0}^{T} \int_{0}^{1}(\Delta u(1, y, t))^{2} d y d t+v_{2} \int_{0}^{T} \int_{0}^{1}(\Delta u(x, 1, t))^{2} d x d t \\
& \quad+\int_{\Omega_{T}} k(x, y)\left(\left(\Delta u_{x}(x, y, t)\right)^{2}+\left(\Delta u_{y}(x, y, t)\right)^{2}\right) d x d y d t \\
& = \\
& v_{1} \int_{0}^{T} \int_{0}^{1} \Delta T_{0}(t) \Delta u(1, y, t) d y d t+v_{2} \int_{0}^{T} \int_{0}^{1} \Delta T_{1}(t) \Delta u(x, 1, t) d x d t  \tag{2.8}\\
& \quad+\int_{\Omega_{T}} \Delta F(x, y, t) \Delta u(x, y, t) d x d y d t
\end{align*}
$$

By the $\epsilon$-Young inequality we make an estimate on the right-hand side integrals of (2.8):

$$
\begin{align*}
& v_{1} \int_{0}^{T} \int_{0}^{1} \Delta T_{0}(t) \Delta u(1, y, t) d y d t+v_{2} \int_{0}^{T} \int_{0}^{1} \Delta T_{1}(t) \Delta u(x, 1, t) d x d t \\
& \quad+\int_{\Omega_{T}} \Delta F(x, y, t) \Delta u(x, y, t) d x d y d t \\
& \leq \\
& \quad \frac{\nu_{1} \epsilon}{2} \int_{0}^{T} \int_{0}^{1}(\Delta u(1, y, t))^{2} d y d t+\frac{v_{1}}{2 \epsilon} \int_{0}^{T}\left(\Delta T_{0}(t)\right)^{2} d t \\
& \quad+\frac{\nu_{2} \epsilon}{2} \int_{0}^{T} \int_{0}^{1}(\Delta u(x, 1, t))^{2} d x d t+\frac{\nu_{2}}{2 \epsilon} \int_{0}^{T}\left(\Delta T_{1}(t)\right)^{2} d t  \tag{2.9}\\
& \quad+\frac{\epsilon}{2} \int_{\Omega_{T}}(\Delta u(x, y, t))^{2} d x d y d t+\frac{1}{2 \epsilon} \int_{\Omega_{T}}(\Delta F(x, y, t))^{2} d x d y d t
\end{align*}
$$

Besides, by the Cauchy-Schwarz inequality the term $(\Delta u(x, y, t))^{2}$ can be estimated as

$$
\begin{aligned}
(\Delta u(x, y, t))^{2} & =\left(\Delta u(1, y, t)-\int_{x}^{1} \Delta u_{\zeta}(\zeta, y, t) d \zeta\right)^{2} \\
& \leq 2(\Delta u(1, y, t))^{2}+2\left(\int_{x}^{1} \Delta u_{\zeta}(\zeta, y, t) d \zeta\right)^{2} \\
& \leq 2(\Delta u(1, y, t))^{2}+2 \int_{0}^{1} \Delta u_{x}(x, y, t)^{2} d x
\end{aligned}
$$

By integrating both sides of the above inequality on $\Omega_{T}$, we obtain

$$
\begin{align*}
\int_{\Omega_{T}}(\triangle u(x, y, t))^{2} d x d y d t \leq & 2 \int_{0}^{T} \int_{0}^{1}(\Delta u(1, y, t))^{2} d y d t \\
& +2 \int_{\Omega_{T}} \triangle u_{x}(x, y, t)^{2} d x d y d t \tag{2.10}
\end{align*}
$$

So, it follows from (2.8)-(2.10) that

$$
\begin{aligned}
& \frac{1}{2} \int_{\Omega}(\Delta u(x, y, T))^{2} d x d y+\left[v_{1}\left(1-\frac{\epsilon}{2}\right)-\epsilon\right] \\
& \quad \times \int_{0}^{T} \int_{0}^{1}(\Delta u(1, y, t))^{2} d y d t+\nu_{2}\left(1-\frac{\epsilon}{2}\right) \int_{0}^{T} \int_{0}^{1}(\Delta u(x, 1, t))^{2} d x d t \\
& \quad+\int_{\Omega_{T}}(k(x, y)-\epsilon)\left(\Delta u_{x}(x, y, t)\right)^{2}+k(x, y)\left(\Delta u_{y}(x, y, t)\right)^{2} d x d y d t \\
& \quad \leq \frac{v_{1}}{2 \epsilon} \int_{0}^{T}\left(\Delta T_{0}(t)\right)^{2} d t+\frac{\nu_{2}}{2 \epsilon} \int_{0}^{T}\left(\Delta T_{1}(t)\right)^{2} d t+\frac{1}{2 \epsilon} \int_{\Omega_{T}}(\Delta F(x, y, t))^{2} d x d y d t
\end{aligned}
$$

Therefore, if we set $k_{*}=\min _{x, y \in[0,1]} k(x, y)>0$ and $\epsilon=\min \left\{k_{*} ; \frac{2 v_{1}}{v_{1}+2}\right\}$, then we deduce that

$$
\begin{aligned}
& \frac{1}{2} \int_{\Omega}(\Delta u(x, y, T))^{2} d x d y \\
& \quad \leq \frac{\nu_{1}}{2 \epsilon} \int_{0}^{T}\left(\Delta T_{0}(t)\right)^{2} d t+\frac{\nu_{2}}{2 \epsilon} \int_{0}^{T}\left(\Delta T_{1}(t)\right)^{2} d t+\frac{1}{2 \epsilon} \int_{\Omega_{T}}(\Delta F(x, y, t))^{2} d x d y d t \\
& \quad \leq \frac{\max \left\{v_{1}, v_{2}, 1\right\}}{2 \epsilon}\|\Delta \omega\|_{W}^{2} .
\end{aligned}
$$

This completes the proof.

With the arguments above we are in a position to give the Fréchet derivative of $J(\omega)$.

Theorem 2.1 Let conditions (2.1)-(2.2) hold. Then the cost functional is Fréchet-differentiable, i.e., $J(\omega) \in C^{1}(W)$. Moreover, the Fréchet derivative of $J(\omega)$ at $\omega \in W$ can be defined, by the solution $\varphi \in H^{1,0}\left(\Omega_{T}\right)$ of the adjoint problem (2.6), as follows:

$$
\begin{equation*}
J^{\prime}(\omega)=\left\{\varphi(x, y, t ; \omega) ; v_{1} \int_{0}^{1} \varphi(1, y, t ; \omega) d y ; v_{2} \int_{0}^{1} \varphi(x, 1, t ; \omega) d x\right\} . \tag{2.11}
\end{equation*}
$$

Corollary 2.1 Let $J(\omega) \in C^{1}(W)$ and $W_{*} \subset W$ be the set of quasi-solutions of the inverse problem (1.5)-(1.6). Then $\omega_{*} \in W_{*}$ is a strict solution of the inverse problem (1.5)-(1.6) if and only if $\varphi\left(x, y, t ; \omega_{*}\right) \equiv 0$, a.e. on $\Omega_{T}$.

By the well-known theory of convex analysis [21], we can get the relationship between the minimization problem and the corresponding variational inequality in the following theorem.

Theorem 2.2 Let conditions of Theorem 2.1 hold; $W \in H^{0}\left(\Omega_{T}\right) \times H^{0}[0, T] \times H^{0}[0, T]$ is a closed convex set of unknown sources and $\varphi=\varphi(x, t ; \omega)$ is the solution of the adjoint problem (2.6), for a given $\omega \in W$. Then the element $w_{*}:=\left\{F_{*}(x, y, t) ; T_{0 *}(t) ; T_{1 *}(t)\right\} \in W$ is a quasisolution of the inverse problem (1.5)-(1.6) if and only if the following variational inequality holds:

$$
\left(J^{\prime}\left(\omega_{*}\right), \omega-\omega_{*}\right)_{W} \geq 0, \quad \forall \omega \in W
$$

where

$$
\begin{aligned}
\left(J^{\prime}\left(\omega_{*}\right), \omega-\omega_{*}\right)_{W}= & \left.\int_{\Omega_{T}} \varphi\left(x, y, t ; \omega_{*}\right)\left[F(x, y, t)-F_{( } x, y, t\right)\right] d x d y d t \\
& +v_{1} \int_{0}^{T}\left[T_{0}(t)-T_{0 *}(t)\right] \int_{0}^{1} \varphi(1, y, t ; \omega) d y d t \\
& +v_{2} \int_{0}^{T}\left[T_{1}(t)-T_{1 *}(t)\right] \int_{0}^{1} \varphi(x, 1, t ; \omega) d x d t
\end{aligned}
$$

## 3 A steepest descent method with line search

An iteration process, known as the steepest descent method in the optimization theory, can thus be implemented

$$
\begin{equation*}
\omega_{n+1}=\omega_{n}-\alpha_{n} J^{\prime}\left(\omega_{n}\right), \tag{3.1}
\end{equation*}
$$

with some appropriate chosen parameter $\alpha_{n}$.
The details of the algorithm are written as follows.

Algorithm 3.1 A steepest descent method with line search
Step 1 Initialization

- Choose an initial approximation $\omega_{0}=\left\{F_{(0)}(x, y, t) ; T_{0}^{(0)}(t) ; T_{1}^{(0)}(t)\right\}$.
- Set the stop tolerance $\varepsilon>0, n=0$.

Step 2 Stopping check

- Solve direct problem (1.2) with $\omega_{n}$ to get $u\left(x, y, T ; \omega_{n}\right)$, then solve (1.5), and get $J^{\prime}\left(\omega_{n}\right)$ as (2.11).
- If $\left\|J^{\prime}\left(\omega_{n}\right)\right\| \leq \varepsilon$ or $\left\|\omega_{n+1}-\omega_{n}\right\| \leq \varepsilon$, then stop and get $\omega_{n}$ as the solution.

Step 3 Update $\omega_{n}$.

- Solve

$$
\begin{equation*}
\min _{\alpha>0} J\left(\omega_{n}-\alpha J^{\prime}\left(\omega_{n}\right)\right) \tag{3.2}
\end{equation*}
$$

with line search and set $\alpha_{n}=\arg \min _{\alpha>0} J\left(\omega_{n}-\alpha J^{\prime}\left(\omega_{n}\right)\right)$.

- Set $\omega_{n+1}=\omega_{n}-\alpha_{n} J^{\prime}\left(\omega_{n}\right), n:=n+1$, go to Step 2 .

Different choices of the parameter $\alpha_{n}$ correspond to various gradient methods. Here we discuss both the exact and the inexact ones.

### 3.1 Exact line search

One of the exact line searches is the golden section method. It is adapted to solve (3.2) when the object function is a unimodal function. The main idea of this method is to construct a sequence of closed intervals $\left\{\left[a_{k}, b_{k}\right]\right\}$, which satisfies $\bar{\alpha} \in\left[a_{k}, b_{k}\right] \subset\left[a_{k-1}, b_{k-1}\right]$ and makes $\left[a_{k}, b_{k}\right]$ scaling down as $k$ increases. If the gradient $J^{\prime}(\omega)$ has Lipschitz continuity, the parameter can be estimated as follows (see Lemma 4.1):

$$
\begin{equation*}
0<\delta_{0} \leq \alpha_{n} \leq 2 /\left(L+2 \delta_{1}\right) \tag{3.3}
\end{equation*}
$$

where $\delta_{0}, \delta_{1}>0$ are arbitrary parameters. We can choose the initial interval as $a_{0}=\delta_{0}$, $b_{0}=2 /\left(L+2 \delta_{1}\right)$.

### 3.2 Inexact line search

There are two famous techniques, the Armijo line search and the Wolfe-Powell line search. If we denote $\phi(\alpha):=J\left(\omega_{n}-\alpha J^{\prime}\left(\omega_{n}\right)\right)$, the Armijo line search can be stated thus: to find $\alpha_{n}>0$ such that

$$
\begin{equation*}
\phi\left(\alpha_{n}\right) \leq \phi(0)+c_{1} \alpha_{n} \phi^{\prime}(0) \tag{3.4}
\end{equation*}
$$

where $0<c_{1}<1$ is constant. While the step length $\alpha_{n}$ found by (3.4) may be quite small or cannot converge to the exact minimum point of (3.2), this situation can be prevented by imposing the curvature condition,

$$
\begin{equation*}
\phi^{\prime}\left(\alpha_{n}\right) \geq c_{2} \phi^{\prime}(0) \tag{3.5}
\end{equation*}
$$

where $0<c_{1}<c_{2}<1$. Equations (3.4) and (3.5) are known collectively as the Wolfe-Powell line search.

## 4 Convergence analysis

Lemma 4.1 Let the conditions of Theorem 2.1 hold. Then the functional $J(\omega)$ is of Hölder class $C^{1,1}(W)$ and

$$
\begin{equation*}
\left\|J^{\prime}(\omega+\Delta \omega)-J^{\prime}(\omega)\right\|_{W} \leq L\|\Delta \omega\|_{W}, \quad \forall \omega, \omega+\Delta \omega \in W \tag{4.1}
\end{equation*}
$$

where

$$
\begin{aligned}
\left\|J^{\prime}(\omega+\Delta \omega)-J^{\prime}(\omega)\right\|_{W}^{2}= & \int_{\Omega_{T}}(\Delta \varphi(x, y, t ; \omega))^{2} d x d y d t \\
& +v_{1}^{2} \int_{0}^{T}\left(\int_{0}^{1} \Delta \varphi(1, y, t ; \omega) d y\right)^{2} d t \\
& +v_{2}^{2} \int_{0}^{T}\left(\int_{0}^{1} \Delta \varphi(x, 1, t ; \omega) d x\right)^{2} d t
\end{aligned}
$$

and $L$ is defined as

$$
L=\sqrt{\frac{\max \left\{v_{1}, v_{2}, 1\right\}}{2 \epsilon}\left(\frac{1}{\min \left\{v_{1}, v_{2}\right\}}+\frac{1}{k_{*}}+v_{1}+v_{2}\right)} .
$$

Proof It is easy to check that the function $\Delta \varphi(x, y, t ; \omega):=\varphi(x, y, t ; \omega+\Delta \omega)-\varphi(x, y, t ; \omega)$ is the solution of the following backward parabolic problem:

$$
\begin{cases}\Delta \varphi_{t}(x, y, t)+\operatorname{div}(k(x, y) \nabla \Delta \varphi(x, y, t))=0, & \text { in } \Omega_{T},  \tag{4.2}\\ \Delta \varphi_{x}(0, y, t)=\Delta \varphi_{y}(x, 0, t)=0, & x, y \in(0,1), t \in(0, T] \\ -k(1, y) \Delta \varphi_{x}(1, y, t)=v_{1} \Delta \varphi(1, y, t), & x, y \in(0,1), t \in(0, T], \\ -k(x, 1) \Delta \varphi_{y}(x, 1, t)=v_{2} \Delta \varphi(x, 1, t), & x, y \in(0,1), t \in(0, T], \\ \Delta \varphi(x, y, T)=\Delta u(x, y, T ; \omega), & (x, y) \in \Omega .\end{cases}
$$

Multiplying both sides of (4.2) by $\Delta \varphi(x, y, t ; \omega)$, integrating on $\Omega_{T}$, and using the initial and boundary conditions, we can obtain the following energy identity:

$$
\begin{align*}
& \frac{1}{2} \int_{\Omega}(\Delta \varphi(x, y, 0 ; \omega))^{2} d x d y+v_{1} \int_{0}^{T} \int_{0}^{1}(\Delta \varphi(1, y, t ; \omega))^{2} d y d t \\
& \quad+v_{2} \int_{0}^{T} \int_{0}^{1}(\Delta \varphi(x, 1, t ; \omega))^{2} d x d t \\
& \quad+\int_{\Omega_{T}} k(x, y)\left[\left(\Delta \varphi_{x}(x, y, t ; \omega)\right)^{2}+\left(\Delta \varphi_{y}(x, y, t ; \omega)\right)^{2}\right] d x d y \\
& \quad=\frac{1}{2} \int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y . \tag{4.3}
\end{align*}
$$

This identity implies the following two inequalities:

$$
\left\{\begin{array}{l}
k_{*} \int_{\Omega_{T}}\left[\left(\Delta \varphi_{x}(x, y, t ; \omega)\right)^{2}+\left(\Delta \varphi_{y}(x, y, t ; \omega)\right)^{2}\right] d x d y  \tag{4.4}\\
\quad \leq \frac{1}{2} \int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y \\
\min \left\{v_{1}, v_{2}\right\}\left[\int_{0}^{T} \int_{0}^{1}(\Delta \varphi(1, y, t ; \omega))^{2} d y d t+\int_{0}^{T} \int_{0}^{1}(\Delta \varphi(x, 1, t ; \omega))^{2} d x d t\right] \\
\quad \leq \frac{1}{2} \int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y
\end{array}\right.
$$

Dealing with $(\Delta \varphi(x, y, t ; \omega))^{2}$ as (2.10) in the proof of Lemma 2.2, we have

$$
\begin{align*}
\int_{\Omega_{T}} & (\Delta \varphi(x, y, t ; \omega))^{2} d x d y d t \\
\leq & \int_{0}^{T} \int_{0}^{1}(\Delta \varphi(1, y, t ; \omega))^{2} d y d t+\int_{0}^{T} \int_{0}^{1}(\Delta \varphi(x, 1, t ; \omega))^{2} d x d t \\
& +\int_{\Omega_{T}}\left[\left(\Delta \varphi_{x}(x, y, t ; \omega)\right)^{2}+\left(\Delta \varphi_{y}(x, y, t ; \omega)\right)^{2}\right] d x d y . \tag{4.5}
\end{align*}
$$

The above three inequalities imply

$$
\int_{\Omega_{T}}(\Delta \varphi(x, y, t ; \omega))^{2} d x d y d t \leq \frac{1}{2}\left(\frac{1}{\min \left\{v_{1}, v_{2}\right\}}+\frac{1}{k_{*}}\right) \int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y
$$

According to the energy identity (4.3), using the Cauchy-Schwarz inequality, we can also conclude

$$
\left\{\begin{align*}
v_{2}^{2} \int_{0}^{T}\left(\int_{0}^{1} \Delta \varphi(x, 1, t ; \omega) d x\right)^{2} d t & \leq v_{2}^{2} \int_{0}^{T} \int_{0}^{1}(\Delta \varphi(x, 1, t ; \omega))^{2} d x d t  \tag{4.6}\\
& \leq \frac{1}{2} v_{2} \int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y \\
v_{1}^{2} \int_{0}^{T}\left(\int_{0}^{1} \Delta \varphi(1, y, t ; \omega) d y\right)^{2} d t & \leq v_{1}^{2} \int_{0}^{T} \int_{0}^{1}(\Delta \varphi(1, y, t ; \omega))^{2} d y d t \\
& \leq \frac{1}{2} \nu_{1} \int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y
\end{align*}\right.
$$

Thus, from the above estimates, we deduce that

$$
\begin{aligned}
\int_{\Omega_{T}} & (\Delta \varphi(x, y, t ; \omega))^{2} d x d y d t+v_{1}^{2} \int_{0}^{T}\left(\int_{0}^{1} \Delta \varphi(1, y, t ; \omega) d y\right)^{2} d t \\
& +v_{2}^{2} \int_{0}^{T}\left(\int_{0}^{1} \Delta \varphi(x, 1, t ; \omega) d x\right)^{2} d t \\
\leq & \frac{1}{2}\left(\frac{1}{\min \left\{v_{1}, v_{2}\right\}}+\frac{1}{k_{*}}+v_{1}+v_{2}\right) \int_{\Omega}(\Delta u(x, y, T ; \omega))^{2} d x d y
\end{aligned}
$$

It follows from this inequality and Lemma 2.2 that we can set

$$
L=\sqrt{\frac{\max \left\{v_{1}, v_{2}, 1\right\}}{2 \epsilon}\left(\frac{1}{\min \left\{v_{1}, v_{2}\right\}}+\frac{1}{k_{*}}+v_{1}+v_{2}\right)},
$$

which completes the proof.

Now we analyze the convergence of the sequence $\left\{J\left(\omega_{n}\right)\right\}$, where the iterations $\omega_{n} \in W$ ( $n=0,1,2, \ldots$ ) are produced by Algorithm 3.1.

Theorem 4.1 Let $W$ be a closed convex set and $J(\omega) \in C^{1,1}(W)$. If $\omega_{n} \in W(n=0,1,2, \ldots)$ is generated by Algorithm 3.1, then $J\left(\omega_{n}\right)$ is a monotone decreasing convergent sequence and

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|J^{\prime}\left(\omega_{n}\right)\right\|=0 \tag{4.7}
\end{equation*}
$$

Proof Set $d_{n}:=-J^{\prime}\left(\omega_{n}\right)$. By using Lemma 4.1, for all $\alpha>0$, we have

$$
\begin{aligned}
J\left(\omega_{n}+\alpha d_{n}\right)-J\left(\omega_{n}\right) & =\int_{0}^{1}\left(J^{\prime}\left(\omega_{n}+\theta \alpha d_{n}\right), \alpha d_{n}\right) d \theta \\
& =\int_{0}^{1}\left(J^{\prime}\left(\omega_{n}+\theta \alpha d_{n}\right)-J^{\prime}\left(\omega_{n}\right), \alpha d_{n}\right) d \theta+\alpha\left(J^{\prime}\left(\omega_{n}\right), d_{n}\right) \\
& \leq \int_{0}^{1} \theta \alpha^{2}\left\|d_{n}\right\|^{2} d \theta+\alpha\left(J^{\prime}\left(\omega_{n}\right), d_{n}\right) \\
& \leq\left(\frac{\alpha^{2}}{2} L-\alpha\right)\left\|J^{\prime}\left(\omega_{n}\right)\right\|^{2}
\end{aligned}
$$

Especially, this inequality holds for $\widehat{\alpha}=1 / L$, i.e.

$$
J\left(\omega_{n}\right)-J\left(\omega_{n}+\widehat{\alpha} d_{n}\right) \geq \frac{1}{2 L}\left\|J^{\prime}\left(\omega_{n}\right)\right\|^{2}
$$

If an exact line search is used in (3.2), one has $J\left(\omega_{n+1}\right)=J\left(\omega_{n}+\alpha_{n} d_{n}\right) \leq J\left(\omega_{n}+\widehat{\alpha} d_{n}\right)$. Thus, the following inequality holds:

$$
\begin{equation*}
J\left(\omega_{n}\right)-J\left(\omega_{n+1}\right) \geq \frac{1}{2 L}\left\|J^{\prime}\left(\omega_{n}\right)\right\|^{2} \tag{4.8}
\end{equation*}
$$

This implies that $J\left(\omega_{n}\right)$ is a monotone decreasing convergent sequence and then (4.7) holds.
For the case of the inexact line search, we consider the Wolfe-Powell line search. By using Lemma 4.1, (3.5) implies

$$
-\left(1-c_{2}\right)\left(J^{\prime}\left(\omega_{n}\right), d_{n}\right) \leq\left(J^{\prime}\left(\omega_{n}+\alpha_{n} d_{n}\right)-J^{\prime}\left(\omega_{n}\right), d_{n}\right) \leq \alpha_{n} L\left\|d_{n}\right\|^{2} .
$$

Thus

$$
\alpha_{n} \geq \frac{1-c_{2}}{L} .
$$

Besides, it follows from (3.4) that

$$
J\left(\omega_{n}\right)-J\left(\omega_{n+1}\right) \geq \alpha_{n}\left\|J^{\prime}\left(\omega_{n}\right)\right\|^{2} \geq \frac{1-c_{2}}{L}\left\|J^{\prime}\left(\omega_{n}\right)\right\|^{2}
$$

This completes the proof.

Denote by

$$
J_{*}:=J\left(\omega_{*}\right)=\lim _{n \rightarrow \infty} J\left(\omega_{n}\right), \quad \omega_{*} \in W
$$

the limit of the sequence $J\left(\omega_{n}\right)$. Let us remark that if $W$ is a closed convex set in $L^{2}(\Omega) \times$ $L^{2}[0, T] \times L^{2}[0, T]$ and the conditions (2.1)-(2.2) hold, then, for any initial data $\omega_{0} \in W$, the sequence of iteration $\left\{\omega_{n}\right\} \subset W$, given by Algorithm 3.1, weakly converges to a quasisolution $\omega_{*} \in W$ of the inverse problem (1.5)-(1.6).

## 5 Conclusions

This paper presents a theoretical study of a case of steady-case heat flow through a plane wall with the two dimensional Robin boundary condition. The inverse problem consists of determining the source terms $\omega:=\left\{F(x, y, t) ; T_{0}(t) ; T_{1}(t)\right\}$ by using observational measurements of the final state $u_{T}(x, y)=u(x, y, T)$. The proposed approach is based on the weak solution theory for parabolic PDEs and the adjoint problem method for minimization of the corresponding cost functional. The adjoint problem is defined to obtain an explicit gradient formula for the cost functional $J(\omega)=\left\|\mu_{T}(x, y)-u(x, y, T ; \omega)\right\|^{2}$. A steepest descent algorithm based on an explicit gradient formula is presented and its convergence is analyzed.

## Competing interests

The authors declare that they have no competing interests.

## Authors' contributions

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