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Stability and bifurcation analysis of an amensalism system with Allee effect

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Abstract

In this work, we propose and study a new amensalism system with Allee effect on the first species. First, we investigate the existence and stability of all possible coexistence equilibrium points and boundary equilibrium points of this system. Then, applying the Sotomayor theorem, we prove that there exists a saddle-node bifurcation under some suitable parameter conditions. Finally, we provide a specific example with corresponding numerical simulations to further demonstrate our theoretical results.

Keywords: Amensalism system; Allee effect; Stability; Saddle-node bifurcation

1 Introduction

The interaction between two or more species has been a central problem in ecology and biology since the famous Lotka–Volterra model was proposed. The interaction between different species generates a complicated dynamics of biological species and exhibits the complexity and diversity. Amensalism, a typical type of interaction between the species, has been intensively considered in the last decades. Amensalism describes a basic biological interaction in nature, where one species inflicts harm on another not affected by the former, which means that it does not receive any costs or benefits to itself. The first pioneer work for the investigations of amensalism model is due to Sun [1], who, in 2003, proposed the following two-species amensalism model:

$$\begin{cases} \frac{dx}{dt} = r_1x\left(\frac{k_1-x-Ty}{k_1}\right), \\ \frac{dy}{dt} = r_2y\left(\frac{k_2-y}{k_2}\right), \end{cases} \quad (1)$$

where $x = x(t)$ and $y = y(t)$ represent the population densities of two species at time t , respectively, and r_1, r_2, k_1, k_2 , and T are positive real numbers. In [1] the author investigated the stability properties of the equilibrium points of this system.

By rescaling we can see that model (1) can be rewritten as

$$\begin{cases} \frac{dx}{dt} = x(r_1 - a_{11}x - a_{12}y), \\ \frac{dy}{dt} = y(r_2 - a_{22}y), \end{cases} \quad (2)$$

where the parameters r_1, r_2, a_{11}, a_{12} , and a_{22} are positive constants.

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Since the first amensalism model was presented, there are numerous relevant works that focus on the complicated dynamics of amensalism models from different aspects [2–12]. For example, in [6] the author discussed the dynamical properties of a two-species amensalism system with nonmonotonic function response. Guan et al. [11] considered a two-species amensalism model with Beddington–DeAngelis functional response. The two-species amensalism model with Michaelis–Menten-type harvesting and a cover for the first species was proposed in [9], where the stability and bifurcation of this system were investigated.

In the nature world the dynamics of population is inevitably affected by difficulties in finding mate, predator avoidance, the evading natural enemies, and resource defense [13–16], where the Allee effect on species occurs. The Allee effect, as a significant phenomenon in population dynamics, was intensively investigated in the last decades. It was found that the Allee effect has a very important impact on the biological system and exhibits complex dynamics. Generally speaking, a population is said to have the Allee effect if the per capita population growth rate and population density have a positive correlation in a small population [17, 18]. Moreover, the Allee effect has two types, weak and strong [19, 20]. The weak Allee effect indicates that per capita growth rate is smaller at low species density, but not negative. The strong Allee effect means that the per capita growth rate is negative at low species density. Nowadays, a lot of efforts have been made to investigate the influence of the Allee effect on the dynamical behavior of biological systems [12, 21–26]. Especially, according to system (2), Wei et al. [12] recently proposed an amensalism model with weak Allee effect for the second species and studied its stability and bifurcation.

Inspired by the previous works, we naturally want to know: for an amensalism system, what about its dynamical properties when an Allee effect is introduced to the first species? Hence, based on system (2), we consider the following amensalism system with a weak Allee effect on the first species:

$$\begin{cases} \frac{dx}{dt} = x\left(\frac{r_1x}{m+x} - a_{11}x - a_{12}y\right), \\ \frac{dy}{dt} = y(r_2 - a_{22}y), \end{cases} \tag{3}$$

where all parameters $r_1, r_2, a_{11}, a_{12}, a_{22}$, and m are positive constants, and the term $K(x) = \frac{x}{m+x}$ represents a weak Allee effect, where m describes the intense of Allee effect on the first species; this function satisfies the following properties [24]:

- (i) $\lim_{x \rightarrow 0} K(x) = 0$, that is, there is no reproduction without partners;
- (ii) $K'(x) > 0$ for all $x \in (0, \infty)$, that is, the Allee effect decreases as population density increases;
- (iii) $\lim_{x \rightarrow \infty} K(x) = 1$, which means that the Allee effect disappears at high population densities.

Setting

$$\bar{t} = r_2 t, \quad \bar{x} = \frac{a_{11}}{r_2} x, \quad \bar{y} = \frac{a_{22}}{r_2} y$$

and dropping the bars, we transform system (3) into

$$\begin{cases} \frac{dx}{dt} = x\left(\frac{\alpha x}{\gamma+x} - x - \beta y\right), \\ \frac{dy}{dt} = y(1 - y), \end{cases} \tag{4}$$

where $\alpha = \frac{r_1}{r_2}$, $\beta = \frac{a_{12}}{a_{22}}$, and $\gamma = \frac{ma_{11}}{r_2}$.

Our main purpose in this paper is investigating the local stability property of the possible equilibrium points and the saddle-node bifurcation of system (4). The rest of this paper is arranged as follows. In Sect. 2, we present the local dynamical behaviors, including the distribution of possible equilibrium points of system (4) and their stability. In Sect. 3, we prove the existence of saddle-node bifurcation when the coefficient α is chosen as a bifurcation parameter. In Sect. 4, we give an example with specific parameter values and the corresponding numerical simulations to further illustrate the validity of the main results. Finally, we end this paper with a brief conclusion in Sect. 5.

2 Existence and stability of equilibrium points

2.1 Existence of equilibrium points

The equilibrium points of system (4) satisfy the following equations:

$$\begin{cases} x\left(\frac{\alpha x}{\gamma+x} - x - \beta y\right) = 0, \\ y(1 - y) = 0. \end{cases} \tag{5}$$

By a simple computation we derive that system (4) has boundary equilibrium points $P_0(0, 0)$ and $P_1(0, 1)$. When $\alpha > \gamma$, there exists another boundary equilibrium point $P_2(\alpha - \gamma, 0)$. Moreover, if $x \neq 0$ and $y \neq 0$, then there exists a coexistence (positive) equilibrium point $P_3(x_3, 1)$, where x_3 is the root of the following equation:

$$f(x) = x^2 + (\beta + \gamma - \alpha)x + \beta\gamma = 0.$$

Set the discriminant of this equation,

$$\Delta(\alpha) = (\beta + \gamma - \alpha)^2 - 4\beta\gamma = \alpha^2 - 2(\beta + \gamma)\alpha + (\beta - \gamma)^2, \tag{6}$$

and denote the roots of $\Delta(\alpha) = 0$ by α_1 and α_2 . Then

$$\alpha_1 = \beta + \gamma - 2\sqrt{\beta\gamma}, \quad \alpha_2 = \beta + \gamma + 2\sqrt{\beta\gamma}. \tag{7}$$

Because the parameters β and γ are positive, using the mean value theorem, we can get that $\alpha_2 > \alpha_1 \geq 0$.

Considering all possible coexistence and boundary equilibria, we obtain the following results.

Theorem 2.1 *For system (4), there always are the boundary equilibrium points $P_0(0, 0)$ and $P_1(0, 1)$. Furthermore, we have:*

- (1) *if $\alpha > \gamma$, then system (4) has another boundary equilibrium point $P_2(\alpha - \gamma, 0)$;*
- (2) *For the possible coexistence equilibrium points,*

- (i) if $\alpha < \alpha_2$, then system (4) has no coexistence equilibrium point;
- (ii) if $\alpha = \alpha_2$, then there is a unique coexistence equilibrium point $P_{31}(x_{31}, 1)$, where $x_{31} = \sqrt{\beta\gamma}$;
- (iii) if $\alpha > \alpha_2$, then there are two coexistence equilibrium points $P_{32}(x_{32}, 1)$ and $P_{33}(x_{33}, 1)$, where $x_{32,33} = \frac{\alpha - \beta - \gamma \pm \sqrt{\Delta}}{2}$.

Proof It is obvious that (1) holds. If $\alpha_1 < \alpha < \alpha_2$, then $\Delta(\alpha) < 0$, which implies that $f(x)$ has no real roots; if $\alpha \leq \alpha_1$, then $f(0) = \beta\gamma > 0$, and the symmetry axis of $f(x)$ is $x = \frac{\alpha - \beta - \gamma}{2} \leq \frac{\alpha_1 - \beta - \gamma}{2} = -\sqrt{\beta\gamma} < 0$, so there is no positive real solution for $f(x) = 0$. Thus we complete the proof of (2)(i).

When $\alpha = \alpha_2$, then $\Delta(\alpha) = 0$, which means that $f(x) = 0$ has a unique positive solution $x_{31} = \sqrt{\beta\gamma}$. Hence the proof of (2)(ii) is completed.

When $\alpha > \alpha_2$, then $\Delta(\alpha) > 0$. Combing this with $f(0) = \beta\gamma > 0$ and the symmetry axis $x = \frac{\alpha - \beta - \gamma}{2} \geq \frac{\alpha_2 - \beta - \gamma}{2} = \sqrt{\beta\gamma} > 0$, we can get that $f(x) = 0$ has two positive roots $x_{32} = \frac{\alpha - \beta - \gamma - \sqrt{\Delta}}{2}$ and $x_{33} = \frac{\alpha - \beta - \gamma + \sqrt{\Delta}}{2}$. Thus the proof of (2)(iii) is finished, and we complete the proof of Theorem 2.1. □

2.2 Stability of the equilibrium points

The Jacobian matrix of system (4) evaluated at any equilibrium point is

$$J(x, y) = \begin{pmatrix} H(x, y) & -\beta x \\ 0 & 1 - 2y \end{pmatrix}, \tag{8}$$

where

$$H(x, y) = \frac{2\alpha\gamma x + \alpha x^2}{(\gamma + x)^2} - 2x - \beta y.$$

Consider the stability of boundary equilibria P_0, P_1 , and P_2 , we have the following results.

Theorem 2.2 *For system (4), the following statements are true.*

- (1) *For the equilibrium point P_0 :*
 - (i) *If $\alpha \neq \gamma$, then P_0 is a saddle node. That is, $S_\varepsilon(P_0)$ is divided into two parts by two separatrices that tend to P_0 along the upside and the underneath of P_0 , where $S_\varepsilon(P_0)$ is a neighborhood of P_0 with sufficient small radius ε . One part consists of two hyperbolic sectors, and the other part is a parabolic sector. Moreover, if $\alpha > \gamma$ (or $\alpha < \gamma$), then the parabolic sector is on the right (or left) half-plane.*
 - (ii) *If $\alpha = \gamma$, then P_0 is a nonhyperbolic saddle.*
- (2) *The equilibrium point P_1 is a hyperbolic stable node. Moreover, P_1 is globally asymptotically stable for $\alpha < \gamma$.*
- (3) *If $\alpha > \gamma$, then the equilibrium point P_2 is a hyperbolic saddle.*

Proof (1) The Jacobian matrix of system (4) evaluated at the equilibrium point P_0 is given by

$$J(P_0) = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}, \tag{9}$$

and the two eigenvalues of $J(P_0)$ are $\lambda_1(P_0) = 0$ and $\lambda_2(P_0) = 1 > 0$. Obviously, the equilibrium P_0 is nonhyperbolic, so it is hard to directly judge its type from the eigenvalues. We further discuss its stability properties by applying Theorem 7.1 in Chap. 2 in [27].

To change system (4) into a standard form, we expand system (4) in power series up to the fourth order around the origin:

$$\begin{cases} \frac{dx}{dt} = \left(\frac{\alpha}{\gamma} - 1\right)x^2 - \beta xy - \frac{\alpha}{\gamma^2}x^3 + \frac{\alpha}{\gamma^3}x^4 + Q_0(x) = P(x, y), \\ \frac{dy}{dt} = y - y^2 = y + Q(x, y), \end{cases} \tag{10}$$

where $Q_0(x)$ represents a power series with terms x^i ($i \geq 5$).

Combining the implicit function theorem with the second equation of (10), we get that there is a unique function $y = \varphi(x) = 0$ such that $\varphi(x) + Q(x, \varphi(x)) = 0$ and $\varphi(0) = \varphi'(0) = 0$. Then substituting $y = \varphi(x) = 0$ into the first equation of (10), we get that

$$\frac{dx}{dt} = \left(\frac{\alpha}{\gamma} - 1\right)x^2 - \frac{\alpha}{\gamma^2}x^3 + \frac{\alpha}{\gamma^3}x^4 + Q_0(x). \tag{11}$$

When $\alpha \neq \gamma$, we obtain that the coefficient at x^2 is $\frac{\alpha}{\gamma} - 1 \neq 0$. So, based on Theorem 7.1 in Chap. 2 in [27], the equilibrium point P_0 is a saddle node. This means that a neighborhood $S_\varepsilon(P_0)$ (ε is a sufficiently small radius) is divided into two parts by two separatrices that tend to P_0 along the upside and the underneath of P_0 . One part is a parabolic sector, and the other part consists of two hyperbolic sectors. Furthermore, if $\alpha > \gamma$ (or $\alpha < \gamma$), then the parabolic sector is on the right (or left) half-plane.

When $\alpha = \gamma$, (11) becomes

$$\frac{dx}{dt} = -\frac{1}{\gamma}x^3 + \frac{1}{\gamma^2}x^4 + Q_0(x). \tag{12}$$

Applying the notations of Theorem 7.1 in Chap. 2 in [27], we have $m = 3$ and $a_m = \frac{1}{\gamma} < 0$, so the equilibrium point P_0 is a nonhyperbolic saddle.

(2) For the equilibrium point P_1 , the Jacobian matrix is

$$J(P_1) = \begin{pmatrix} -\beta & 0 \\ 0 & -1 \end{pmatrix}. \tag{13}$$

The two eigenvalues of this matrix $J(P_1)$ are $\lambda_1(P_1) = -\beta < 0$ and $\lambda_2(P_1) = -1 < 0$, and thus P_1 is a hyperbolic stable node.

When $\alpha < \gamma$, there are two equilibrium points P_0 , which is unstable, and P_1 , which is locally asymptotically stable. To ensure that P_1 is globally asymptotically stable, we consider the Lyapunov function $V(x, y) = x$. Then

$$\frac{dV}{dt} = x \left(\frac{\alpha x}{\gamma + x} - x - \beta y \right) = \frac{\alpha - \gamma - x}{\gamma + x} x^2 - \beta xy.$$

Obviously, if $\alpha < \gamma$, then $\frac{dV}{dt} < 0$ for all $x > 0, y > 0$. So P_1 is globally asymptotically stable for $\alpha < \gamma$.

(3) The Jacobian matrix of system (4) at equilibrium point P_2 is

$$J(P_2) = \begin{pmatrix} -\frac{\alpha^2 - 2\alpha\gamma + \gamma^2}{\alpha} & -\beta(\alpha - \gamma) \\ 0 & 1 \end{pmatrix}. \tag{14}$$

When $\alpha > \gamma$, it is obvious that $J(P_1)$ has two characteristic roots $\lambda_1(P_1) = -\frac{\alpha^2 - 2\alpha\gamma + \gamma^2}{\alpha} < 0$ and $\lambda_2(P_1) = 1 > 0$, so the equilibrium point P_2 is a hyperbolic saddle.

This ends the proof of Theorem 2.2. □

Discussing the stability of the coexistence equilibria P_{31} , P_{32} , and P_{33} , we have the following results.

Theorem 2.3 *For system (4), we have the following statements.*

- (1) *If $\alpha = \alpha_2$, then the equilibrium point P_{31} is a saddle node, that is, $S_\varepsilon(P_{31})$ is divided into two parts by two separatrices that tend to P_{31} along the upside and the underneath of P_{31} , where $S_\varepsilon(P_{31})$ is a neighborhood of P_{31} with sufficient small radius ε . One part consists of two hyperbolic sectors, and the other part is a parabolic sector. Moreover, the parabolic sector is on the right half-plane.*
- (2) *If $\alpha > \alpha_2$, then the equilibrium point P_{32} is a hyperbolic saddle, and the equilibrium point P_{33} is a hyperbolic stable node.*

Proof (1) The Jacobian matrix of system (4) at the equilibrium point P_{31} is given by

$$J(P_{31}) = \begin{pmatrix} H(x_{31}, 1) & -\beta\sqrt{\beta\gamma} \\ 0 & -1 \end{pmatrix}. \tag{15}$$

Note that $P_{31}(\sqrt{\beta\gamma}, 1)$ satisfies the equation $\frac{\alpha x_{31}}{\gamma + x_{31}} = x_{31} + \beta$. By some calculations we can derive that $H(x_{31}, 1) = \frac{\alpha_2 \gamma x_{31}}{(\gamma + x_{31})^2} - x_{31} = 0$. Then two characteristic roots of $J(P_{31})$ are $\lambda_1(P_{31}) = 0$ and $\lambda_2(P_{31}) = -1 < 0$, so the equilibrium point P_{31} is nonhyperbolic, and its stability cannot be given directly from the eigenvalues. We next analyze its stability behavior by using Theorem 7.1 in Chap. 2 in [27].

Letting $X = x - x_{31}$ and $Y = y - 1$, we translate the equilibrium point P_{31} to the origin and then expand the corresponding system in power series up to the third order around the origin, which makes system (4) to be the following form:

$$\begin{cases} \frac{dX}{dt} = c_0 + c_1X - \beta x_{31}Y + c_2X^2 - \beta XY + c_3X^3 + Q_1(X), \\ \frac{dY}{dt} = -Y - Y^2, \end{cases} \tag{16}$$

where

$$c_0 = \frac{\alpha_2 \gamma^2}{\gamma + x_{31}} + \alpha_2(x_{31} - \gamma) - \beta x_{31} - x_{31}^2 = 0,$$

$$c_1 = \frac{1}{\gamma + x_{31}}(\beta\gamma - x_{31}^2) = 0, \quad c_2 = \frac{\alpha_2 \gamma^2}{(\gamma + x_{31})^3} - 1, \quad c_3 = -\frac{\alpha_2 \gamma^2}{(\gamma + x_{31})^4},$$

and $Q_1(X)$ represents a power series with terms X^i ($i \geq 4$).

To transform the Jacobian matrix into a standard form, we use the invertible translation

$$\begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} 1 & -\beta x_{31} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} X \\ Y \end{pmatrix}. \tag{17}$$

Then system (16) becomes

$$\begin{cases} \frac{du}{dt} = d_{20}u^2 + d_{11}uv + d_{02}v^2 + d_{30}u^3 + d_{21}u^2v + d_{12}uv^2 + d_{03}v^3 \\ \quad + Q_2(u, v), \\ \frac{dv}{dt} = -v - v^2, \end{cases} \tag{18}$$

where

$$\begin{aligned} d_{20} &= -1 + \frac{\alpha_2 \gamma^2}{(\gamma + x_{31})^3}, & d_{11} &= -\beta \left(1 + 2x_{31} - \frac{2\alpha_2 \gamma^2 x_{31}}{(\gamma + x_{31})^3} \right), \\ d_{02} &= \beta x_{31} \left(1 - \beta - \beta x_{31} + \frac{\alpha_2 \beta \gamma^2 x_{31}}{(\gamma + x_{31})^3} \right), & d_{30} &= -\frac{\alpha_2 \gamma^2}{(\gamma + x_{31})^4}, \\ d_{21} &= -\frac{3\alpha_2 \beta x_{31} \gamma^2}{(\gamma + x_{31})^4}, & d_{12} &= -\frac{3\alpha_2 \gamma^2 \beta^2 x_{31}^2}{(\gamma + x_{31})^4}, & d_{03} &= -\frac{\alpha_2 \gamma^2 \beta^3 x_{31}^3}{(\gamma + x_{31})^4}, \end{aligned}$$

and $Q_2(u, v)$ is a power series with terms $u^i v^j$ ($i + j \geq 4$).

By introducing a new time variable $T = -t$ we get

$$\begin{cases} \frac{du}{dT} = -d_{20}u^2 - d_{11}uv - d_{02}v^2 - d_{30}u^3 - d_{21}u^2v - d_{12}uv^2 - d_{03}v^3 \\ \quad - Q_2(u, v) = P(u, v), \\ \frac{dv}{dT} = v + v^2 = v + Q(u, v). \end{cases} \tag{19}$$

Based on the implicit function theorem, from $\frac{dv}{dT} = 0$ we can deduce a unique function $v = \phi(u) = 0$ satisfying $\phi(0) = \phi'(0) = 0$ and $\phi(u) + Q(u, \phi(u)) = 0$. Then substituting it into the first equation of (19), we have that

$$\frac{du}{dT} = -d_{20}u^2 - d_{30}u^3 + Q_3(u),$$

where $Q_3(u)$ is a power series with terms u^i ($i \geq 4$).

Thus we can derive that the coefficient at the term u^2 is

$$-d_{20} = -\frac{\alpha_2 \gamma^2}{(\gamma + x_{31})^3} + 1 = -\frac{\gamma^2}{(\gamma + x_{31})^3} \left(\frac{\beta}{x_{31}} - \frac{2x_{31}}{\gamma} - \frac{x_{31}^2}{\gamma^2} \right) > 0.$$

Then, according to Theorem 7.1 in Chap. 2 in [27], we obtain that the equilibrium point P_{31} is a saddle node. This means that a neighborhood $S_\varepsilon(P_{31})$ (ε is a sufficiently small radius) of P_{31} is divided into two parts by two separatrices that tend to P_{31} along the upside and the underneath of P_{31} . One part consists of two hyperbolic sectors, and the other part is a parabolic sector. Moreover, the parabolic sector is on the right half-plane because the coefficient at the term u^2 is greater than zero.

(2) If $\alpha > \alpha_2$, then we have $f(\sqrt{\beta\gamma}) = \beta\gamma + (\beta + \gamma - \alpha)\sqrt{\beta\gamma} + \beta\gamma < 0$. Hence $x_{32} < \sqrt{\beta\gamma}$ and $x_{33} > \sqrt{\beta\gamma}$.

The Jacobian matrix of system (4) evaluated at the equilibrium P_{32} is calculated as

$$J(P_{32}) = \begin{pmatrix} H(x_{32}, 1) & -\beta\sqrt{\beta\gamma} \\ 0 & -1 \end{pmatrix}. \tag{20}$$

Then from $\frac{\alpha x_{32}}{\gamma + x_{32}} = x_{32} + \beta$ it follows that

$$H(x_{32}, 1) = \frac{2\alpha\gamma x_{32} + \alpha x_{32}^2}{(\gamma + x_{32})^2} - 2x_{32} - \beta = \frac{\beta\gamma - x_{32}^2}{\gamma + x_{32}} > 0.$$

Therefore we get that two characteristic roots of the matrix $J(P_{32})$ are $\lambda_1(P_{32}) = \frac{\beta\gamma - x_{32}^2}{\gamma + x_{32}} > 0$ and $\lambda_2(P_{32}) = -1 < 0$, implying that P_{32} is a hyperbolic saddle.

Similarly to the proof of the former P_{32} , we can deduce that the two eigenvalues of $J(P_{33})$ are $\lambda_1(P_{33}) = \frac{\beta\gamma - x_{33}^2}{\gamma + x_{33}} < 0$ and $\lambda_2(P_{33}) = -1 < 0$, which means that P_{33} is a hyperbolic stable node.

This completes the proof of Theorem 2.3. □

3 Saddle-node bifurcation

In Sect. 2.1, we have derived the parameter conditions for the existence of the coexistence equilibrium point P_{31} . We next show that system (4) experiences a saddle-node bifurcation at equilibrium P_{31} as the parameter α passes through the bifurcation value $\alpha = \alpha_{SN} = \beta + \gamma + 2\sqrt{\beta\gamma}$ by applying the Sotomayor theorem [28]. For saddle-node bifurcation, we formulate the following results.

Theorem 3.1 *System (4) undergoes a saddle-node bifurcation at the equilibrium point P_{31} if $\alpha = \alpha_{SN} = \beta + \gamma + 2\sqrt{\beta\gamma}$. Moreover, two equilibrium points bifurcate from P_{31} for $\alpha > \alpha_{SN}$, coalesce as the equilibrium point P_{31} for $\alpha = \alpha_{SN}$, and disappear for $\alpha < \alpha_{SN}$.*

Proof Now we will prove the transversality condition for the occurrence of a saddle-node bifurcation at $\alpha = \alpha_{SN}$ by utilizing the Sotomayor theorem. From Sect. 2.1 we see that the two eigenvalues of $J(P_{31})$ are $\lambda_1(P_{31}) = 0$ and $\lambda_2(P_{31}) = -1 < 0$. Denote by V and W the eigenvectors corresponding to the eigenvalue $\lambda_1(P_{31})$ for the matrices $J(P_{31})$ and $J(P_{31})^T$, respectively. Some computations yield

$$V = \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}; \quad W = \begin{pmatrix} w_1 \\ w_2 \end{pmatrix} = \begin{pmatrix} 1 \\ -\beta\sqrt{\beta\gamma} \end{pmatrix}.$$

Furthermore, we obtain

$$F_\alpha(P_{31}; \alpha_{SN}) = \begin{pmatrix} \frac{x^2}{\gamma+x} \\ 0 \end{pmatrix}_{(P_{31}; \alpha_{SN})} = \begin{pmatrix} \frac{x_{31}^2}{\gamma+x_{31}} \\ 0 \end{pmatrix} \tag{21}$$

and

$$\begin{aligned}
 D^2F(P_{31}; \alpha_{SN})(V, V) &= \begin{pmatrix} \frac{\partial^2 F_1}{\partial x^2} v_1^2 + 2 \frac{\partial^2 F_1}{\partial x \partial y} v_1 v_2 + \frac{\partial^2 F_1}{\partial y^2} v_2^2 \\ \frac{\partial^2 F_2}{\partial x^2} v_1^2 + 2 \frac{\partial^2 F_2}{\partial x \partial y} v_1 v_2 + \frac{\partial^2 F_2}{\partial y^2} v_2^2 \end{pmatrix}_{(P_{31}; \alpha_{SN})} \\
 &= \begin{pmatrix} \frac{2\alpha_{SN}\gamma^2}{(\gamma + x_{31})^3} - 2 \\ 0 \end{pmatrix}. \tag{22}
 \end{aligned}$$

Obviously, the vectors V and W satisfy

$$W^T F_\alpha(P_{31}; \alpha_{SN}) = \frac{x_{31}^2}{\gamma + x_{31}} \neq 0$$

and

$$W^T [D^2F(P_{31}; \alpha_{SN})(V, V)] = \frac{2\alpha_{SN}\gamma^2}{(\gamma + x_{31})^3} - 2 = -\frac{2\sqrt{\beta\gamma}}{\gamma + \sqrt{\beta\gamma}} \neq 0.$$

Hence, by the Sotomayor theorem, when $\alpha = \alpha_{SN}$, system (4) undergoes a saddle-node bifurcation at nonhyperbolic critical point P_{31} . The number of positive equilibrium points of system (4) changes from zero to two as α passes from the left of $\alpha = \alpha_{SN}$ to the right. More concretely, there are no equilibrium points when $\alpha < \alpha_{SN}$, there is one equilibrium point P_{31} when $\alpha = \alpha_{SN}$, and there are two equilibrium points P_{32} and P_{33} when $\alpha > \alpha_{SN}$. Thus the proof of Theorem 3.1 is completed. \square

4 Numerical simulations

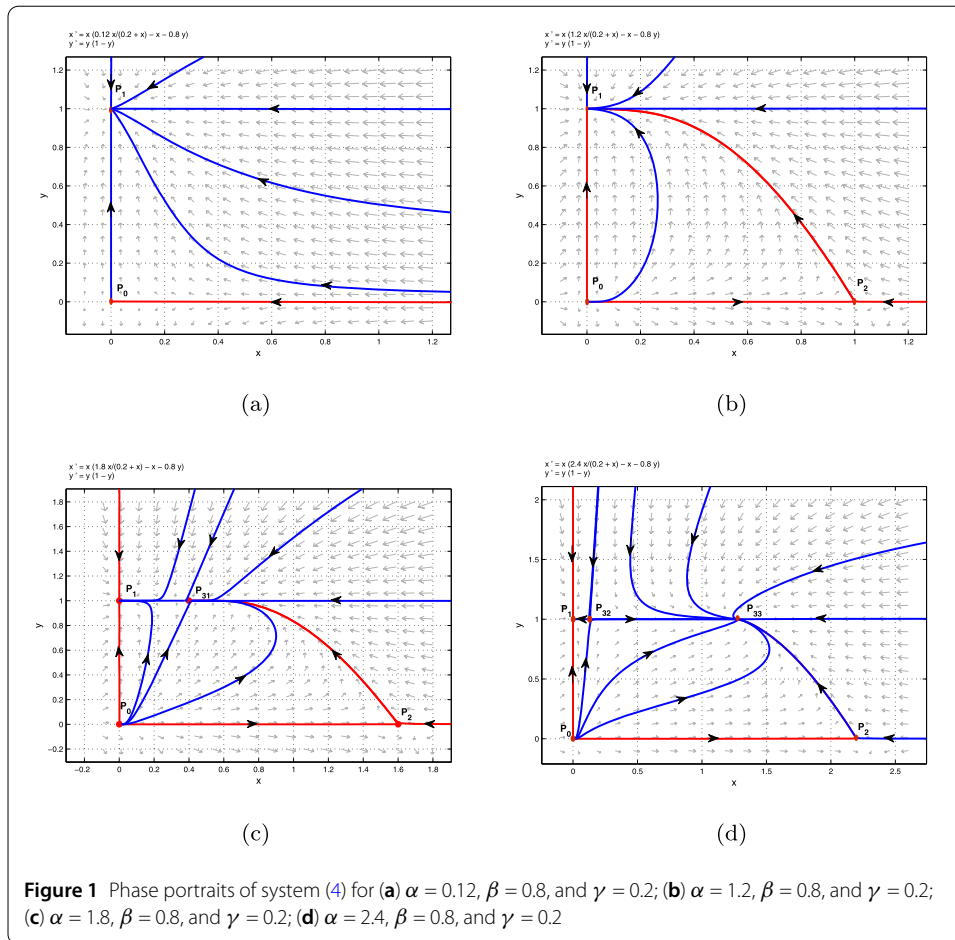
In this section, by numerical simulations we give some phase portraits of the amensalism model (4) to further illustrate the previous theoretical analysis. Without loss of generality, we consider the following example.

Example 4.1 Consider the system

$$\begin{cases} \frac{dx}{dt} = x\left(\frac{\alpha x}{0.2+x} - x - 0.8y\right), \\ \frac{dy}{dt} = y(1 - y). \end{cases} \tag{23}$$

In system (4), we give specific parameter values $\beta = 0.8$ and $\gamma = 0.2$. Then we get system (23). By some calculations we have $\alpha_1 = \beta + \gamma - 2\sqrt{\beta\gamma} = 0.2$ and $\alpha_2 = \beta + \gamma + 2\sqrt{\beta\gamma} = 1.8$. According to Theorems 2.1 and 2.2, for system (4), there exist two boundary equilibrium points $P_0(0, 0)$ and $P_1(0, 1)$ for all positive parameters; $P_0(0, 0)$ is a saddle node, and $P_1(0, 1)$ is a stable node. Next, we consider the following four cases by choosing various parameters α .

Case 1. Fix $\alpha = 0.12$. Then for system (23), there is no other equilibrium point, as shown in Fig. 1(a). Furthermore, we choose the initial condition $(x(0), y(0)) = (1, 0.6)$. From Fig. 3(a) we can see that as t increases, the trajectories tend to P_1 , which is a stable node.

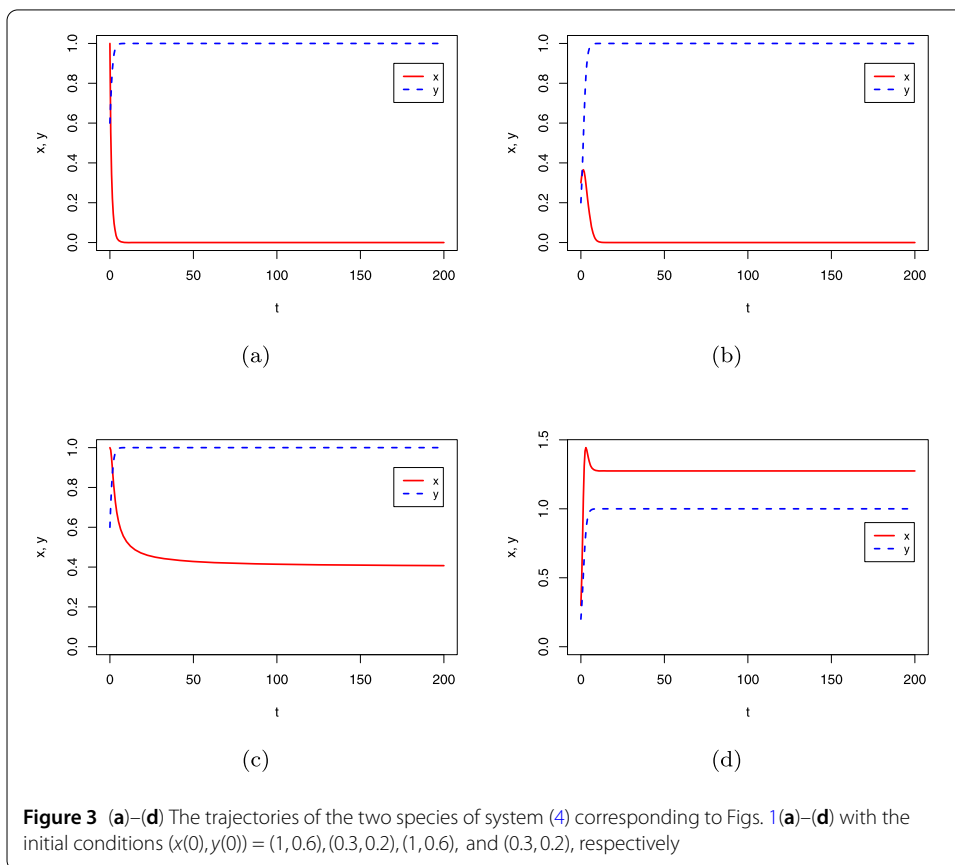
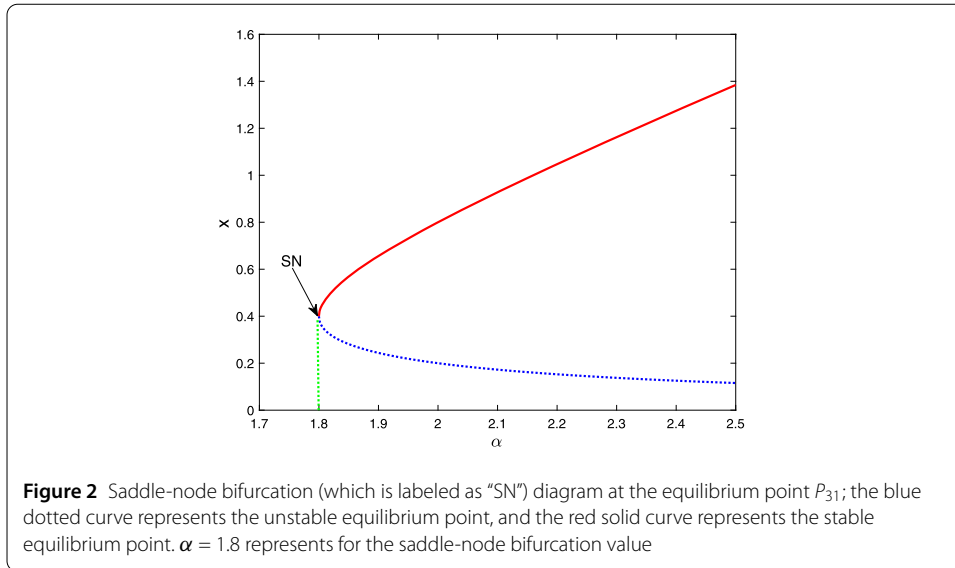


Case 2. Fix $\alpha = 1.2$. Then for system (23), there is a boundary equilibrium point $P_2 = (1, 0)$ that is a saddle depicted in Fig. 1(b). Furthermore, we take the initial value $(x(0), y(0)) = (0.3, 0.2)$, and from Fig. 3(b) it follows that the orbits converge to the stable node P_1 as t increases.

Case 3. Fix $\alpha = 1.8$. Then system (23) has one boundary equilibrium point $P_2 = (1.6, 0)$, which is a saddle, and one coexistence equilibrium point $P_{31} = (0.4, 1)$. From Theorem 2.3(1) we know that P_{31} is a saddle node, and the above dynamical behaviors are shown in Fig. 1(c). From Theorem 3.1 we have that system (4) undergoes a saddle-node bifurcation at the equilibrium P_{31} for $\alpha = \alpha_{SN} = 1.8$; the corresponding illustrative bifurcation diagram is depicted in Fig. 2. Moreover, we take the initial value $(x(0), y(0)) = (1, 0.6)$, it follows from Fig. 3(c) that the trajectories approach the saddle node P_{31} as t increases.

Case 4. Fix $\alpha = 2.4$. Then system (23) has one boundary equilibrium point $P_2 = (2.2, 0)$, which is a saddle, and two positive equilibrium points $P_{32} \approx (0.1255, 1)$ and $P_{33} \approx (1.2745, 1)$. According to Theorem 2.3(2), we have that P_{32} is a saddle, P_{33} is a stable node, as displayed in Fig. 1(d). Moreover, we choose the initial value $(x(0), y(0)) = (0.3, 0.2)$. From Fig. 3(d) we can observe that the orbits tend to the stable node P_{33} as t increases.

Figs. 3(a)–(d) display the trajectories of system (4) corresponding to Figs. 1(a)–(d) with the different initial values.



5 Conclusions

In this paper, we proposed a new amensalism system with Allee effect on the first species and studied the stability and bifurcation of this model. We discussed the distribution and stability of all boundary equilibrium points and the coexistence equilibrium points. The dynamical properties of the amensalism model become complex when the amensalism

system subjects to the Allee effect. By comparing with system (1), we can see the system (4) possesses some new dynamical phenomena as follows.

- (i) There exist at most five equilibrium points for system (4), including two distinct interior points, whereas for system (1), there are at most four equilibrium points, including one interior point.
- (ii) As the number of equilibrium points increases, the dynamical properties of the system (4) become more complicated; for example, system (4) undergoes a saddle-node bifurcation at equilibrium point P_{31} .
- (iii) The stability of coexistence equilibrium points is more complex. In [1] the coexistence equilibrium point (if it exists) of system (1) is globally stable, whereas for system (4), the equilibrium point P_{31} is a saddle node when $\alpha = \alpha_2$.

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Competing interests

The authors declare that they have no competing interests.

Authors' contributions

Both authors contributed equally to the writing of this paper. Both authors read and approved the final manuscript.

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