# Bifurcation analysis of a discrete SIR epidemic model with constant recovery 

Hui Cao ${ }^{1 *}$ © , Huan Wu' and Xiaoqin Wang ${ }^{1}$

"Correspondence:
caohui@sust.edu.cn
${ }^{1}$ Department of Mathematics, Shaanxi University of Science and Technology, Xi'an, China


#### Abstract

We state and study a discrete SIR epidemic model with bilinear incidence rate and constant recovery. We obtain conditions for the existence of the disease-free equilibrium and endemic equilibria. Theoretic analysis shows that the disease-free equilibrium is globally asymptotically stable when the basic reproduction number is less than unity, and the numerical simulations illustrate that it is asymptotically stable when the number is greater than unity. We also obtain conditions for the stability of the endemic equilibria. More attention is paid to the existence of the saddle-node bifurcation, the flip bifurcation, the 1:1 resonance, and the Neimark-Sacker bifurcation. We obtain sufficient conditions for those bifurcations. Our numerical simulations demonstrate our theoretical results and the complexity of the model.


MSC: 39A10; 92D30
Keywords: The discrete epidemic model; Flip bifurcation; Neimark-Sacker bifurcation; 1: 1 resonance

## 1 Introduction

Mathematical models describing the population dynamics of infectious diseases have been playing an important role in better understanding epidemiological patterns and disease control for a long time [1]. In some cases a simple mathematical model can reveal the nature of the infectious disease transmission, which plays an important role in the control and prevention of the infectious diseases. In general, differential equations and difference equations are two typical mathematical approaches to modeling epidemic dynamical systems.

In recent years, there has been an increasing interest on discrete population dynamical systems because of their rich dynamics behavior and suitability. On the one hand, the discrete models exhibit richer dynamical behavior than the continuous models, which brings more challengeable problems for researches, and more interesting results can be obtained. For example, the simple logistic model $x_{n+1}=r x_{n}\left(1-x_{n} / K\right)$, the Ricker model $x_{n+1}=x_{n} e^{r-x_{n} / K}$, and the Hassell model $x_{n+1}=\lambda x_{n}\left(1+a x_{n}\right)^{-b}$ all exhibit rich dynamical behavior [2-5]. On the other hand, since epidemiological data are usually collected at discrete time intervals, it is easier to compare the numerical results of discrete epidemic model with real data. In addition, the straightforward recurrence relationship of the dif-

[^0]ference equation models is easier to be understood, which is also a prominent advantage over the difference equation models.
In theoretical epidemiology the Euler discretization has been extensively utilized to build the discrete epidemic models [1, 6-14]. The main idea is discretizing the existing continuous models. However, the disadvantage of this method is that it cannot guarantee the nonnegativity and boundedness of the solutions of the systems in some cases. As a result, some scholars establish discrete epidemic models by using the probability method [15-17]. Although the probability method can guarantee the nonnegativity and boundedness of the solutions, the model is more complex, and theoretical analysis is difficult. In this paper, we build a discrete SIR model by applying the Euler discretization. We are lucky that the solution of our model is nonnegative and bounded.
Theoretical results of discrete epidemic models mainly focus on these aspects, including the computation of the basic reproduction number $[8,18,19]$, the comparison between continuous-time epidemic models and discrete-time epidemic models [10, 20], the local stability and global stability of the disease-free equilibrium and endemic equilibrium [6, 7 , $9,11,15,16,21$ ], the extinction, persistence, and permanence of a disease [12, 22-25], periodic systems [8, 19, 26], bifurcations and chaos phenomena [1, 7, 9, 13, 14, 17, 25, 27, 28], and so on. In particular, these papers about the bifurcation analysis of discrete epidemic model mainly involve the conditions on the existence of codimension-one bifurcations, such as fold bifurcation, flip bifurcation, Neimark-Sacker bifurcation [1, 7, 9, 25, 27, 28], and so on, which were derived by using the center manifold theorem and bifurcation theory. Here [7] and [9] both focus on the influence of discretization step size on the dynamic behavior of the system. However, there are few studies on codimension-two bifurcations [13, 14], for example, bifurcations associated with $1: 1$ strong resonance. In [14], based on the analysis of various strong resonance, a controller was designed to eliminate the disease. In this paper, we are more interested in finding conditions of the existence of the saddle-node bifurcation, the flip bifurcation, the $1: 1$ resonance, and the Neimark-Sacker bifurcation.

Wang and Ruan [29] have used a continuous-time SIR model tell us that it is unnecessary to take such a large treatment capacity that endemic equilibria disappear to eradicate the disease, and the outcome of disease spread may depend on the position of the initial states for certain range of parameters. We put the continuous-time SIR model in [29] as follows:

$$
\begin{align*}
& \frac{d S(t)}{d t}=\Lambda-\beta S(t) I(t)-d S(t), \\
& \frac{d I(t)}{d t}=\beta S(t) I(t)-(d+\gamma) I(t)-h(I),  \tag{1.1}\\
& \frac{d R(t)}{d t}=\gamma I(t)+h(I)-d R(t),
\end{align*}
$$

where $S(t), I(t)$, and $R(t)$ denote the numbers of susceptible, infective, and recovered individuals at time $t$, respectively, $\Lambda$ is the recruitment rate, $\beta$ is the transmission rate, $d$ is the natural death rate, $\gamma$ is the spontaneous recovery rate of the infective individuals, and $h(I)$ is the removal rate of infective individuals due to the treatment of the form

$$
h(I)= \begin{cases}m>0 & \text { for } I>0  \tag{1.2}\\ m=0 & \text { for } I=0\end{cases}
$$

This means that if the number of infective individuals is greater than 0 , then it is possible for the infective individuals to be cured, leaving the infective class and entering the recovered class. If the number of infective individuals were 0 , then no infective individuals would enter the recovered class.
In this paper, we apply the forward Euler scheme to model (1.1) and obtain the following discrete-time SIR epidemic model:

$$
\begin{align*}
& S(t+1)=S(t)+\Lambda-\beta S(t) I(t)-d S(t) \\
& I(t+1)=I(t)+\beta S(t) I(t)-(d+\gamma) I(t)-h(I)  \tag{1.3}\\
& R(t+1)=R(t)+\gamma I(t)+h(I)-d R(t)
\end{align*}
$$

Here all parameters have the same meaning as in system (1.1), and $h(I)=m \geq 0$. We assume that all the parameters are positive constants except for $m$. Based on the epidemiological meaning, we require that $1-d>0,1-d-\gamma>0$, and $\Lambda>m$.

Due to the biological interpretation of system (1.3), only nonnegative solutions are meaningful to be considered. The following result reveals that the solutions $(S(t), I(t)$, $R(t)$ ) of system (1.3) with non-negative initial value ultimately remain nonnegative and bounded.

Theorem 1.1 All solutions of system (1.3) with nonnegative initial value remain nonnegative for all $t \geq 0$ and are ultimately bounded.

Proof Let $N(t)$ be the number of the total population at time $t$, that is, $N(t)=S(t)+I(t)+$ $R(t)$. Adding all equations in (1.3), we obtain the equation for the total population $N(t)$ :

$$
\begin{equation*}
N(t+1)=\Lambda+(1-d) N(t) . \tag{1.4}
\end{equation*}
$$

Substituting $S(t)=N(t)-I(t)-R(t)$ into the second equation of model (1.3) leads to the system

$$
\begin{align*}
& N(t+1)=\Lambda+(1-d) N(t) \\
& I(t+1)=I(t)+\beta(N(t)-I(t)-R(t)) I(t)-(d+\gamma) I(t)-h(I),  \tag{1.5}\\
& R(t+1)=R(t)+\gamma I(t)+h(I)-d R(t) .
\end{align*}
$$

Obviously, systems (1.5) and (1.3) are equivalent. In the following, we show that the solutions $(N(t), I(t), R(t))$ of system (1.5) with nonnegative initial value remain nonnegative and ultimately bounded.
Firstly, by the first equation of system (1.5) we have $N(t)=\Lambda\left[\frac{1-(1-d)^{t}}{d}\right]+(1-d)^{t} N_{0}$, which implies that $N(t) \geq 0$ for all $t \geq 0$ when the initial value $N_{0} \geq 0$.

Secondly, we show that $I(t)$ is nonnegative for $t \geq 0$. In fact, if there exists $t_{1}>0$ such that $I\left(t_{1}\right)=0$ and $I(t)>0$ for $t \in\left(0, t_{1}\right)$, then by the second equation of system (1.5) we have $I\left(t_{1}+1\right)=0$. This implies that $I(t) \geq 0$ for all $t \geq 0$.

Finally, we prove that $R(t) \geq 0$ for $t \geq 0$. The third equation of system (1.5) implies that $R(t)=(1-d)^{t} R(0)+(1-d)^{t-1}[\gamma I(0)+h(I)]+(1-d)^{t-2}[\gamma I(1)+h(I)]+\cdots+\gamma I(t-1)+h(I)>0$ with nonnegative initial value for $t \geq 0$.

Summarizing this analysis, we know that any solution of system (1.5) with nonnegative initial values remains nonnegative for all $t \geq 0$, that is, all solutions of system (1.3) with nonnegative initial value remain nonnegative for all $t \geq 0$.
In the following, we will show that the solutions of system (1.3) are ultimately bounded. It is clear that equation (1.4) has a unique equilibrium $N_{*}=\frac{\Lambda}{d}$, which is globally asymptotically stable, that is, $\lim _{t \rightarrow \infty} N(t)=N_{*}$. Therefore

$$
\limsup _{t \rightarrow+\infty}(S(t)+I(t)+R(t)) \leq \frac{\Lambda}{d}
$$

It follows that the omega limit set of system (1.3) is contained in the bounded feasible region

$$
\Gamma=\left\{(S, I, R): S, I, R \geq 0, S+I+R \leq \frac{\Lambda}{d}\right\} .
$$

Obviously, this region is positively invariant with respect to system (1.3), which implies that system (1.3) is ultimately bounded.

Since the first two equations of system (1.3) are independent of the third one, it suffices to consider the first two equations. Thus we replace $h(I)$ with $m$ in system (1.3) and restrict our attention to the following reduced model:

$$
\begin{align*}
& S(t+1)=S(t)+\Lambda-\beta S(t) I(t)-d S(t)  \tag{1.6}\\
& I(t+1)=I(t)+\beta S(t) I(t)-(d+\gamma) I(t)-m
\end{align*}
$$

where $m>0$ when $I>0$ and $m=0$ when $I=0$. Accordingly,

$$
\bar{\Gamma}=\left\{(S, I): S, I \geq 0, S+I<\frac{\Lambda}{d}\right\}
$$

is positively invariant with respect to system (1.6).
The purpose of this paper is to show the dynamics of discrete system (1.6), In particular, we mainly focus on the various bifurcations that might occur around the equilibria. The paper is organized as follows. The saddle-node bifurcation, flip bifurcation, NeimarkSacker bifurcation, and 1:1 resonance are shown in Sect. 2. In Sect. 3, a series of numerical simulations show that there are bifurcation and chaos in the discrete epidemic model. Finally, we give some conclusions.

## 2 Bifurcation analysis

Before bifurcation analysis, we first discuss the existence of the equilibria of system (1.6) and the global stability of the disease-free equilibrium. Letting $R_{0}=\frac{\Lambda \beta}{d(d+\gamma)}>0$ and $A=$ $\frac{m \beta}{d(d+\gamma)}>0$, we have the following result.

## Theorem 2.1

(i) System (1.6) always has the disease free equilibrium $E_{0}=\left(\frac{\Lambda}{d}, 0\right)$;
(ii) If $R_{0}=1+A+2 \sqrt{A}$, then system (1.6) has one endemic equilibrium $E_{1}=\left(S_{1}, I_{1}\right)$, where

$$
S_{1}=\frac{\Lambda}{d+\beta I_{1}}, \quad I_{1}=\frac{d \sqrt{A}}{\beta}
$$

(iii) If $R_{0}>1+A+2 \sqrt{A}$, then system (1.6) has two endemic equilibria $E_{2}=\left(S_{2}, I_{2}\right)$ and $E_{3}=\left(S_{3}, I_{3}\right)$, where $S_{i}=\frac{\Lambda}{d+\beta I_{i}}, i=2,3$, and

$$
\begin{aligned}
& I_{2}=\frac{d}{2 \beta}\left(R_{0}-1-A-\sqrt{\left(R_{0}-1-A\right)^{2}-4 A}\right) \\
& I_{3}=\frac{d}{2 \beta}\left(R_{0}-1-A+\sqrt{\left(R_{0}-1-A\right)^{2}-4 A}\right)
\end{aligned}
$$

Proof To discuss the existence of equilibria of system (1.6), we need discuss the following equations:

$$
\left\{\begin{array}{l}
\Lambda-\beta S I-d S=0  \tag{2.1}\\
\beta S I-(d+\gamma) I-m=0
\end{array}\right.
$$

When $I=0$, we have $m=0$. By the first equation of (2.1) we obtain $S=\frac{\Lambda}{d}$. This implies that system (1.6) always has the disease-free equilibrium $E_{0}=\left(\frac{\Lambda}{d}, 0\right)$.

When $I>0$, we have $m>0$. The first equation of (2.1) implies that $S=\frac{\Lambda}{\beta I+d}$. Taking $S=\frac{\Lambda}{\beta I+d}$ in the second equation of of (2.1), we obtain the following equation:

$$
\begin{equation*}
\frac{\beta}{d} I^{2}+\left(A+1-R_{0}\right) I+\frac{m}{d+\gamma}=0 \tag{2.2}
\end{equation*}
$$

where $R_{0}=\frac{\Lambda \beta}{d(d+\gamma)}$ and $A=\frac{m \beta}{d(d+\gamma)}$, that is, the following results hold:
(i) If $R_{0}=1+A+2 \sqrt{A}$, then (2.2) has one root $I_{1}=\frac{d \sqrt{A}}{\beta}$, which implies that system (1.6) has one endemic equilibrium $E_{1}=\left(S_{1}, I_{1}\right)$, where

$$
S_{1}=\frac{\Lambda}{d+\beta I_{1}}, \quad I_{1}=\frac{d \sqrt{A}}{\beta}
$$

(ii) If $R_{0}>1+A+2 \sqrt{A}$, then (2.2) has two roots $I_{2}=\frac{d}{2 \beta}\left(R_{0}-1-A-\sqrt{\left(R_{0}-1-A\right)^{2}-4 A}\right)$ and $I_{3}=\frac{d}{2 \beta}\left(R_{0}-1-A+\sqrt{\left(R_{0}-1-A\right)^{2}-4 A}\right)$, respectively. This means that system (1.6) has two endemic equilibria $E_{2}=\left(S_{2}, I_{2}\right)$ and $E_{3}=\left(S_{3}, I_{3}\right)$, where $S_{i}=\frac{\Lambda}{d+\beta I_{i}}, i=2,3$, and

$$
\begin{aligned}
& I_{2}=\frac{d}{2 \beta}\left(R_{0}-1-A-\sqrt{\left(R_{0}-1-A\right)^{2}-4 A}\right), \\
& I_{3}=\frac{d}{2 \beta}\left(R_{0}-1-A+\sqrt{\left(R_{0}-1-A\right)^{2}-4 A}\right) .
\end{aligned}
$$

Theorem 2.2 If $R_{0}<1$, then the disease-free equilibrium $E_{0}$ is globally asymptotically stable.

Proof We discuss the global stability of $E_{0}$ by constructing a Lyapunov function. Define $F:\left[0, N_{*}\right] \rightarrow\left[0, N_{*}\right]$ by

$$
F(I)=I(t)+\beta S(t) I(t)-(d+\gamma) I(t)-m .
$$

Obviously, $F$ is the mapping derived by the second equation of system (1.6), and 0 is a fixed point of $F$. The linear function $V(I)=I$ on $\left[0, N_{*}\right]$ is continuous and positive definite with respect to 0 . Therefore $V$ is a Lyapunov function on the domain of $F$. For any $I \in\left[0, N_{*}\right]$, we have

$$
\begin{aligned}
\Delta V(I) & =I(t)+\beta S(t) I(t)-(d+\gamma) I(t)-m-I(t) \\
& =\beta S(t) I(t)-(d+\gamma) I(t)-m \\
& \leq(\beta S(t)-(d+\gamma)) I(t) \\
& \leq\left(\frac{\Lambda}{d} \beta-(d+\gamma)\right) I(t) \\
& =(d+\gamma)\left(R_{0}-1\right) I(t) .
\end{aligned}
$$

Hence, if $R_{0}<1$, then $\Delta V(I)<0$ for $I \in\left[0, N_{*}\right]$, which implies that $\lim _{t \rightarrow \infty} I(t)=0$, that is, $\lim _{t \rightarrow \infty} S(t)=\frac{\Lambda}{d}$. It follows from Theorem 4.22 in [30] that $E_{0}$ is globally asymptotically stable.

In the following sections, we focus on bifurcations that may occur around the equilibria points.

### 2.1 Saddle-node bifurcation and bifurcation with $1: 1$ resonance

The Jacobian matrix at the endemic equilibrium $E_{1}=\left(S_{1}, I_{1}\right)$ is

$$
J_{E_{1}}=\left(\begin{array}{cc}
1-d-\beta I_{1} & -\beta S_{1} \\
\beta I_{1} & 1-d-\gamma+\beta S_{1}
\end{array}\right) .
$$

The corresponding characteristic equation of $J_{E_{1}}$ is

$$
h_{1}(\lambda)=\lambda^{2}-(2-d+\gamma \sqrt{A}) \lambda+1-d+\gamma \sqrt{A}=0 .
$$

When $R_{0}=A+2 \sqrt{A}+1$, direct computation yields that

$$
h_{1}(1)=0, \quad h_{1}(-1)=2(2-d+\gamma \sqrt{A}), \quad h_{1}(0)=1-d+\gamma \sqrt{A},
$$

which implies that $h_{1}(\lambda)=0$ has eigenvalues $\lambda_{1}=1$ and $\lambda_{2}=1-d+\gamma \sqrt{A}$. If $\sqrt{A} \neq \frac{d}{\gamma}$, then $\left|\lambda_{2}\right| \neq 1$. The following theorem is the case where the endemic equilibrium $E_{1}$ is a saddlenode bifurcation point when $R_{0}=1+A+2 \sqrt{A}$. The existence of the endemic equilibria $E_{2}$ and $E_{3}$ is exhibited in Theorem 2.1, and the stability of both $E_{2}$ and $E_{3}$ is proved in the next section, Therefore we have the following:

Theorem 2.3 If $R_{0}=1+A+2 \sqrt{A}$ and $\sqrt{A} \neq \frac{d}{\gamma}$, then system (1.6) will undergo a saddlenode bifurcation at $E_{1}$. Moreover, when $R_{0}>1+A+2 \sqrt{A}$, then system (1.6) has two endemic equilibria $E_{2}$ and $E_{3}$.

In addition, when $R_{0}=1+A+2 \sqrt{A}$ and $\sqrt{A}=\frac{d}{\gamma}$, we have $h_{1}(1)=0$ and $h_{1}(0)=1$, which implies that

$$
\begin{equation*}
h_{1}(\lambda)=\lambda^{2}-2 \lambda+1=0, \tag{2.3}
\end{equation*}
$$

that is, in this case, $h_{1}(\lambda)=0$ has the eigenvalues $\lambda_{1}=\lambda_{2}=1$. The following theorem is the case where the endemic equilibrium $E_{1}$ is a $1: 1$ resonance point when $R_{0}=1+A+2 \sqrt{A}$ and $\sqrt{A}=\frac{d}{\gamma}$.

Theorem 2.4 If $R_{0}=1+A+2 \sqrt{A}$ and $\sqrt{A}=\frac{d}{\gamma}$, then system (1.6) will undergo a $1: 1$ resonance at $E_{1}$.

Proof Let $u_{1}(t)=S(t)-S_{1}, v_{1}(t)=I(t)-I_{1}$, and $\rho=\beta-\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}$. Then system (1.6) becomes

$$
\begin{align*}
u_{1}(t+1)= & \left(1-\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) I_{1}-d\right) u_{1}(t)-\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) S_{1} v_{1}(t) \\
& -\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) u_{1}(t) v_{1}(t) \\
v_{1}(t+1)= & \left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) I_{1} u_{1}(t)+\left(1-d-\gamma+\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) S_{1}\right) v_{1}(t)  \tag{2.4}\\
& +\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) u_{1}(t) v_{1}(t)
\end{align*}
$$

Furthermore, system (2.4) can be rewritten as

$$
\begin{equation*}
\binom{u_{1}(t+1)}{v_{1}(t+1)}=A_{0}\binom{u_{1}(t)}{v_{1}(t)}+\binom{g_{1}\left(u_{1}, v_{1}\right)}{g_{2}\left(u_{1}, v_{1}\right)}, \tag{2.5}
\end{equation*}
$$

where

$$
A_{0}=\left(\begin{array}{cc}
1-\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) I_{1}-d & -\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) S_{1} \\
\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) I_{1} & 1-d-\gamma+\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) S_{1}
\end{array}\right),
$$

and

$$
g_{1}\left(u_{1}, v_{1}\right)=-\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) u_{1}(t) v_{1}(t), \quad g_{2}\left(u_{1}, v_{1}\right)=\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) u_{1}(t) v_{1}(t) .
$$

Let

$$
P=\left(\begin{array}{cc}
-(d+\gamma) & \frac{\gamma-(d+\gamma)^{2}}{d} \\
d & d+\gamma
\end{array}\right) \quad \text { and } \quad P^{-1}=\left(\begin{array}{cc}
-\frac{d+\gamma}{\gamma} & \frac{\gamma-(d+\gamma)^{2}}{d \gamma} \\
\frac{d}{\gamma} & \frac{d+\gamma}{\gamma}
\end{array}\right) .
$$

The transformation

$$
\binom{y_{1}(t)}{y_{2}(t)}=P\binom{u_{1}(t)}{v_{1}(t)}
$$

for system (2.5) yields that

$$
\binom{y_{1}(t+1)}{y_{2}(t+1)}=\left(\begin{array}{ll}
1 & 1  \tag{2.6}\\
0 & 1
\end{array}\right)\binom{y_{1}(t)}{y_{2}(t)}+\binom{k_{1}\left(y_{1}, y_{2}\right)}{k_{2}\left(y_{1}, y_{2}\right)}
$$

where

$$
\begin{aligned}
k_{1}\left(y_{1}, y_{2}\right)= & \left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right)\left(d+\gamma+\frac{\gamma-(d+\gamma)^{2}}{d}\right)\left(-\frac{d(d+\gamma)}{\gamma^{2}} y_{1}^{2}\right. \\
& \left.+\frac{\gamma-2(d+\gamma)^{2}}{\gamma^{2}} y_{1} y_{2}+\frac{(d+\gamma)\left(\gamma-(d+\gamma)^{2}\right)}{d \gamma^{2}} y_{2}^{2}\right), \\
k_{2}\left(y_{1}, y_{2}\right)= & \left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) \gamma\left(-\frac{d(d+\gamma)}{\gamma^{2}} y_{1}^{2}+\frac{\gamma-2(d+\gamma)^{2}}{\gamma^{2}} y_{1} y_{2}\right. \\
& \left.+\frac{(d+\gamma)\left(\gamma-(d+\gamma)^{2}\right)}{d \gamma^{2}} y_{2}^{2}\right) .
\end{aligned}
$$

Using Sect. 9.5.2 in [31], we have

$$
\begin{aligned}
& a_{20}(\rho)=-2\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right)\left(d+\gamma+\frac{\gamma-(d+\gamma)^{2}}{d}\right)\left(\frac{d(d+\gamma)}{\gamma^{2}}\right), \\
& b_{20}(\rho)=-2\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) \gamma\left(\frac{d(d+\gamma)}{\gamma^{2}}\right), \\
& b_{11}(\rho)=\left(\rho+\frac{d(d+\gamma)^{3}}{\Lambda \gamma^{2}}\right) \gamma \frac{\gamma-2(d+\gamma)^{2}}{\gamma^{2}} .
\end{aligned}
$$

Furthermore, we have

$$
\begin{aligned}
& b_{20}(0)=-2 \frac{2 d^{2}(d+\gamma)^{4}}{\Lambda \gamma^{3}}<0, \\
& a_{20}(0)+b_{11}(0)-b_{20}(0)=-\frac{d(d+\gamma)^{3}(d \gamma(1-d-\gamma)+\gamma(d+\gamma)(1-d))}{\Lambda \gamma^{3}}<0 .
\end{aligned}
$$

It follows from Sect. 9.5.2 in [31] that system (1.6) will undergo a $1: 1$ resonance at $E_{1}$ when $R_{0}=1+A+2 \sqrt{A}$ and $\sqrt{A}=\frac{d}{\gamma}$.

In fact, the phenomenon of resonance with a double eigenvalue of 1 is similar to the Bogdanov-Takens bifurcation of the continuous system with double eigenvalues 0 .

### 2.2 Flip bifurcation and Neimark-Sacker bifurcation

In this section, we first discuss the stability of the endemic equilibrium $E_{2}$. Obviously, the Jacobian matrix at the endemic equilibrium $E_{2}=\left(S_{2}, I_{2}\right)$ is as follows:

$$
J_{E_{2}}=\left(\begin{array}{cc}
1-d-\beta I_{2} & -\beta S_{2} \\
\beta I_{2} & 1-d-\gamma+\beta S_{2}
\end{array}\right)
$$

and the corresponding characteristic equation of $J_{E_{2}}$ is

$$
\begin{aligned}
h_{2}(\lambda)= & \lambda^{2}-\left(1-d-\beta I_{2}+1-d-\gamma+\beta S_{2}\right) \lambda+(1-d)(1-d-\gamma)+(1-d) \beta S_{2} \\
& -(1-d-\gamma) \beta I_{2} \\
= & 0 .
\end{aligned}
$$

Because of $h_{2}(1)=-d(d+\gamma)\left(\sqrt{\left(R_{0}-1-A\right)^{2}-4 A}\right)<0$ when $R_{0}>1+A+2 \sqrt{A}$, the endemic equilibrium $E_{2}$ is unstable, that is, we have the following:

Theorem 2.5 If $R_{0}>1+A+2 \sqrt{A}$, then the endemic equilibrium $E_{2}$ is unstable.
In the following, we discuss the stability of the endemic equilibrium $E_{3}$ and the bifurcation that may occur around $E_{3}$ when the stability of $E_{3}$ changes. Let

$$
R_{1}=\frac{-d \gamma+\sqrt{\Delta_{1}}}{2 d(1-d)(d+\gamma)(1-d-\gamma)}, \quad R_{2}=\frac{2(2-d) \gamma+2 \sqrt{\Delta_{2}}}{(2-d)(2-d-\gamma) d(d+\gamma)}
$$

with

$$
\begin{aligned}
& \Delta_{1}=d^{2} \gamma^{2}+4 d(1-d)(d+\gamma)(1-d-\gamma)\left(A(d(1-d-\gamma)+(1-d)(d+\gamma))^{2}+d^{2}\right) \\
& \Delta_{2}=(2-d)^{2} \gamma^{2}+(2-d)(2-d-\gamma) d(d+\gamma)\left(A(d+(1-d)(d+\gamma))^{2}+(2-d)^{2}\right)
\end{aligned}
$$

It is clear that $R_{1}>0$ and $R_{2}>0$.

Theorem 2.6 If $1+A+R_{1}<R_{0}<1+A+R_{2}$, then the endemic equilibrium $E_{3}$ is locally asymptotically stable, whereas if $R_{0}<1+A+R_{1}$, or $R_{0}>1+A+R_{2}$, then the endemic equilibrium $E_{3}$ is unstable.

Proof To study the stability of $E_{3}$, we give the Jacobian matrix at the endemic equilibrium $E_{3}=\left(S_{3}, I_{3}\right)$ :

$$
J_{E_{3}}=\left(\begin{array}{cc}
1-d-\beta I_{3} & -\beta S_{3} \\
\beta I_{3} & 1-d-\gamma+\beta S_{3}
\end{array}\right)
$$

and the corresponding characteristic equation of $J_{E_{3}}$ is

$$
\begin{aligned}
h_{3}(\lambda)= & \lambda^{2}-\left(1-d-\beta I_{3}+1-d-\gamma+\beta S_{3}\right) \lambda+(1-d)(1-d-\gamma)+(1-d) \beta S_{3} \\
& -(1-d-\gamma) \beta I_{3} \\
= & 0 .
\end{aligned}
$$

It is clear that

$$
\begin{aligned}
& h_{3}(1)=d(d+\gamma) \sqrt{\left(R_{0}-1-A\right)^{2}-4 A}>0 \\
& h_{3}(-1)=2(2-d)+\gamma\left(R_{0}-1-A\right)-(d+(1-d)(d+\gamma)) \sqrt{\left(R_{0}-1-A\right)^{2}-4 A}, \\
& h_{3}(0)=1-d+\frac{\gamma}{2}\left(R_{0}-1-A\right)-\frac{(1-d)(d+\gamma)+d(1-d-\gamma)}{2} \sqrt{\left(R_{0}-1-A\right)^{2}-4 A} .
\end{aligned}
$$

By direct computation we have $h_{3}(0)<1$ if $R_{0}>1+A+R_{1}, h_{3}(0)=1$ if $R_{0}=1+A+R_{1}$, and $h_{3}(0)>1$ if $R_{0}<1+A+R_{1}$. In addition, we have $h_{3}(-1)>0$ if $R_{0}<1+A+R_{2}, h_{3}(-1)=0$ if $R_{0}=1+A+R_{2}$, and $h_{3}(-1)<0$ if $R_{0}>1+A+R_{2}$.

Therefore we have $h_{3}(-1)>0$ and $h_{3}(0)<1$ if $1+A+R_{1}<R_{0}<1+A+R_{2}$, which implies that the endemic equilibrium $E_{3}$ is locally asymptotically stable, whereas if $R_{0}<1+A+R_{1}$ or $R_{0}>1+A+R_{2}$, then the endemic equilibrium $E_{3}$ is unstable.

In particular, when $R_{0}=1+A+R_{2}, h_{3}(\lambda)=0$ has the eigenvalues $\lambda_{1}=-1$ and

$$
\lambda_{2}=1-d-\frac{\gamma d(d+\gamma)}{2(d+(1-d)(d+\gamma))} R_{2}-\frac{d \gamma}{d+(1-d)(d+\gamma)} .
$$

Because of $0<d<1$, we have $\left|\lambda_{2}\right| \neq 1$. The following theorem considers the case where the endemic equilibrium $E_{3}$ is a flip bifurcation point when $R_{0}=1+A+R_{2}$.

Theorem 2.7 If $R_{0}=1+A+R_{2}$, then system (1.6) will undergo a flip bifurcation at $E_{3}$. Moreover, the periodic- 2 solution bifurcated from the endemic equilibrium $E_{3}$ is stable.

Proof In the case where $R_{0}=1+A+R_{2}$ the endemic equilibrium is $E_{3}=\left(S_{3}, I_{3}\right)$, where

$$
S_{3}=\frac{\Lambda}{d+\beta I_{3}}, \quad I_{3}=\frac{d(2-d)\left(2+(d+\gamma) R_{2}\right)}{2 \beta(d+(1-d)(d+\gamma))} .
$$

Let $\varepsilon(t)=S(t)-S_{3}, \delta(t)=I(t)-I_{3}$, and $\tau(t)=\Lambda-C$ with $C=\frac{d(d+\gamma)}{\beta}+m+\frac{d(d+\gamma)}{\beta} R_{2}$, where the parameter $\tau$ is a new and dependent variable. The system (1.6) becomes:

$$
\begin{align*}
& \varepsilon(t+1)=\left(1-d-\beta I_{3}\right) \varepsilon(t)-\frac{\beta C}{d+\beta I_{3}} \delta(t)-\beta \varepsilon(t) \delta(t)-\frac{\beta}{d+\beta I_{3}} \delta(t) \tau(t) \\
& \delta(t+1)=\beta I_{3} \varepsilon(t)+\left(1-d-\gamma+\frac{\beta C}{d+\beta I_{3}}\right) \delta(t)+\beta \varepsilon(t) \delta(t)+\frac{\beta}{d+\beta I_{3}} \delta(t) \tau(t),  \tag{2.7}\\
& \tau(t+1)=\tau(t)
\end{align*}
$$

Let

$$
M=\left(\begin{array}{ccc}
-1-\gamma & -(2-d-\gamma) & 0 \\
m_{21} & 2-d & 0 \\
0 & 0 & 1
\end{array}\right) \quad \text { and } \quad M^{-1}=\left(\begin{array}{ccc}
\frac{2-d}{\phi} & \frac{2-d-\gamma}{\phi} & 0 \\
-\frac{m_{21}}{\phi} & \frac{-1-\gamma}{\phi} & 0 \\
0 & 0 & 1
\end{array}\right)
$$

where

$$
\begin{aligned}
& m_{21}=-\frac{\gamma d(d+\gamma)}{2(d+(1-d)(d+\gamma))} R_{2}-\frac{d \gamma}{d+(1-d)(d+\gamma)}<0, \\
& \phi=(2-d-\gamma) m_{21}-(1+\gamma)(2-d)<0 .
\end{aligned}
$$

Using the transformation

$$
\left(\begin{array}{l}
x(t) \\
y(t) \\
z(t)
\end{array}\right)=M^{-1}\left(\begin{array}{l}
\varepsilon(t) \\
\delta(t) \\
\tau(t)
\end{array}\right)
$$

for system (2.7) yields that

$$
\left(\begin{array}{l}
x(t+1)  \tag{2.8}\\
y(t+1) \\
z(t+1)
\end{array}\right)=\left(\begin{array}{ccc}
\lambda_{2} & 0 & 0 \\
0 & -1 & 0 \\
0 & 0 & 1
\end{array}\right)\left(\begin{array}{l}
x(t) \\
y(t) \\
z(t)
\end{array}\right)+\left(\begin{array}{c}
g_{1} \\
g_{2} \\
0
\end{array}\right),
$$

where

$$
\begin{aligned}
g_{1}= & \frac{\gamma}{\phi} \beta(1+\gamma) m_{21} x^{2}(t)+\frac{((1+\gamma)(2-d)+2-d-\gamma) \beta \gamma}{\phi} x(t) y(t) \\
& +\frac{\gamma \beta(2-d-\gamma)(2-d)}{\phi} y^{2}(t)-\frac{\gamma \beta m_{21}}{\phi\left(d+\beta I_{3}\right)} x(t) z(t)-\frac{\gamma \beta(2-d)}{\phi\left(d+\beta I_{3}\right)} y(t) z(t), \\
g_{2}= & -\frac{\left(m_{21}-1-\gamma\right) \beta(1+\gamma)}{\phi} x^{2}(t)-\frac{\left(m_{21}-1-\gamma\right) \beta}{\phi}((1+\gamma)(2-d)+2-d-\gamma) x(t) y(t) \\
& -\frac{\left(m_{21}-1-\gamma\right) \beta(1+\gamma)}{\phi}(2-d)(2-d-\gamma) y^{2}(t)+\frac{\left(m_{21}-1-\gamma\right) \beta m_{21}}{\phi\left(d+\beta I_{3}\right)} x(t) z(t) \\
& +\frac{\left(m_{21}-1-\gamma\right) \beta(2-d)}{\phi\left(d+\beta I_{3}\right)} y(t) z(t) .
\end{aligned}
$$

From the center manifold theory of discrete system we know that there exists a local manifold of system (2.8) [32]. The local manifold has the expansion

$$
x(t)=b_{1} y^{2}(t)+b_{2} y(t) z(t)+b_{3} z^{2}(t)+o\left((|y(t)|+z(t))^{3}\right) .
$$

After substituting the expansion into system (2.8) and using the invariance property of the local manifold, the straightforward and careful calculation gives

$$
b_{1}=-\frac{\beta \gamma(2-d)(2-d-\gamma)}{\phi\left(\lambda_{2}-1\right)}, \quad b_{2}=\frac{\beta \gamma(2-d)}{\phi\left(d+\beta I_{3}\right)\left(\lambda_{2}+1\right)}, \quad b_{3}=0 .
$$

From the third equation of system (2.8) we have that $z(t)$ is always constant. Therefore the one-dimensional model induced by the center manifold is

$$
y(t+1)=G_{1}(y(t), z)
$$

where

$$
\begin{aligned}
G_{1}(y(t), z)= & -y(t)-\frac{\left(m_{21}-1-\gamma\right) \beta(1+\gamma)}{\phi}(2-d)(2-d-\gamma) y^{2}(t) \\
& +\frac{\left(m_{21}-1-\gamma\right) \beta(2-d)}{\phi\left(d+\beta I_{3}\right)} y(t) z(t) \\
& +\frac{\left(m_{21}-1-\gamma\right) \beta^{2} \gamma}{\phi^{2}\left(\lambda_{2}-1\right)}((1+\gamma)(2-d)+2-d-\gamma)(2-d)(2-d-\gamma) y^{3}(t) \\
& -\frac{\left(m_{21}-1-\gamma\right) \beta^{2} \gamma(2-d)}{\phi^{2}\left(d+\beta I_{3}\right)}\left(\frac{(1+\gamma)(2-d)+2-d-\gamma}{\lambda_{2}+1}\right. \\
& \left.+\frac{m_{21}(2-d-\gamma)}{\lambda_{2}-1}\right) y^{2}(t) z(t) \\
& +\frac{\left(m_{21}-1-\gamma\right) \beta^{2} \gamma m_{21}(2-d)}{\phi^{2}\left(d+\beta I_{3}\right)^{2}\left(\lambda_{2}+1\right)} y(t) z^{2}(t)+o\left((|v(t)|+\delta(t))^{4}\right) .
\end{aligned}
$$

It is clear that $G_{1}(0, z(t))=0, \frac{\partial G_{1}(0,0)}{\partial y}=-1$. In addition, $\frac{\partial^{2} G(0,0)}{\partial y \partial z}=\frac{\left(m_{21}-1-\gamma\right) \beta(2-d)}{\phi\left(d+\beta I_{3}\right)}>0$ since $m_{21}-1-\gamma<0$ and $\phi<0$, and $\frac{\partial^{3} G^{2}(0,0)}{\partial y^{3}}=-12\left(\frac{\left(m_{21}-1-\gamma\right)^{2} \beta^{2}(1+\gamma)^{2}(2-d)^{2}(2-d-\gamma)^{2}}{\phi^{2}}+\right.$ $\left.\frac{\left(m_{21}-1-\gamma\right) \beta^{2} \gamma((1+\gamma)(2-d)+2-d-\gamma)(2-d)(2-d-\gamma)}{\phi^{2}\left(\lambda_{2}-1\right)}\right)<0$ since $m_{21}-1-\gamma<0$ and $\lambda_{2}-1<0$. Therefore system (1.6) undergoes a flip bifurcation at $E_{3}$, and the bifurcation solution of period 2 is stable [32]. The proof is completed.

When $R_{0}=1+A+R_{1}$ and $\frac{2 d}{\gamma}<R_{1}<\frac{4 d(1-d-\gamma)+2 d^{2}(d+\gamma)+4(1-d)(d+\gamma)}{d \gamma(d+\gamma)}$, we have $h_{3}(0)=1$ and $\Delta=\left(1-d-\beta I_{3}+1-d-\gamma+\beta S_{3}\right)^{2}-4<0$. This implies that $h_{3}(\lambda)=0$ has a pair of conjugate complex roots, and their molds both are equal to 1 . Furthermore the endemic equilibrium $E_{3}$ may be a Neimark-Sacker bifurcation point. Therefore we have the following:

Theorem 2.8 If $R_{0}=1+A+R_{1}$ and $\frac{2 d}{\gamma}<R_{1}<\frac{4 d(1-d-\gamma)+2 d^{2}(d+\gamma)+4(1-d)(d+\gamma)}{d \gamma(d+\gamma)}$, then system (1.6) will undergo a Neimark-Sacker bifurcation at $E_{3}$.

Proof Let $x_{1}(t)=S(t)-S_{3}$ and $y_{1}(t)=I(t)-I_{3}$. Then system (1.6) becomes

$$
\begin{align*}
& x_{1}(t+1)=\left(1-d-\beta I_{3}\right) x_{1}(t)-\beta S_{3} y_{1}(t)-\beta x_{1}(t) y_{1}(t)  \tag{2.9}\\
& y_{1}(t+1)=\beta I_{3} x_{1}(t)+\left(1-d-\gamma+\beta S_{3}\right) y_{1}(t)+\beta x_{1}(t) y_{1}(t)
\end{align*}
$$

The eigenvalues of the linearized matrix are

$$
\lambda_{1,2}=\frac{1}{2}\left(B_{1} \pm \sqrt{4 B_{2}-B_{1}^{2}} i\right),
$$

where

$$
\begin{aligned}
& B_{1}=2-d+\frac{\gamma}{2}\left(R_{0}-1-A\right)-\frac{2 d+\gamma}{2} \sqrt{\left(R_{0}-1-A\right)^{2}-4 A}, \\
& B_{2}=1-d+\frac{\gamma}{2}\left(R_{0}-1-A\right)-\frac{2 d+\gamma-d(d+\gamma)}{2} \sqrt{\left(R_{0}-1-A\right)^{2}-4 A},
\end{aligned}
$$

and $\left|\lambda_{i}\right|=1-d+(1-d)(d+\gamma)\left(R_{0}-1-A\right)-\frac{d(1-d-\gamma)+(1-d)(d+\gamma)}{2}\left(R_{0}-1-A+\right.$ $\left.\sqrt{\left(R_{0}-1-A\right)^{2}-4 A}\right), i=1,2$.

Let $R_{0}$ be the bifurcation parameter, and let $R_{0}^{*}=1+A+R_{1}$. A direct calculation yields that

$$
\begin{aligned}
& \lambda_{1,2}\left(R_{0}^{*}\right)=\left.\frac{1}{2}\left(B_{1} \pm \sqrt{4 B_{2}-B_{1}^{2}} i\right)\right|_{R_{0}=1+A+R_{1}}=\frac{1}{2}\left(B \pm \sqrt{4-B^{2}} i\right), \quad\left|\lambda_{1,2}\left(R_{0}^{*}\right)\right|=1, \\
& \left.\frac{d\left|\lambda_{i}\right|}{d R_{0}}\right|_{R_{0}=R_{0}^{*}}=-\frac{d \gamma+2 d(d+\gamma)(1-d)(1-d-\gamma) R_{1}}{\left(\gamma R_{1}-2 d\right)}<0,
\end{aligned}
$$

with $B=\frac{2 d+2(1-d)^{2}(d+\gamma)-\gamma d(d+\gamma) R_{1}}{d(1-d-\gamma)+(1-d)(d+\gamma)}$.
In addition, $\lambda_{1,2}^{m} \neq 1$ for $m=1,2,3,4$. Let $T=\left(\begin{array}{cc}1 & 0 \\ t_{21} & t_{22}\end{array}\right)$ with $t_{22}=\frac{\left(d+\beta I_{3}\right) \sqrt{4 B_{2}-B_{1}^{2}}}{\Lambda \beta}>0$ and $t_{21}=-\frac{\left(d+\beta I_{3}\right)\left(\frac{1}{2} B_{1}-1+d+\beta I_{3}\right)}{\Lambda \beta}<0$. By the transformation

$$
\binom{x_{1}(t)}{y_{1}(t)}=T\binom{\bar{x}(t)}{\bar{y}(t)}
$$

we obtain

$$
\binom{\bar{x}(t+1)}{\bar{y}(t+1)}=\left(\begin{array}{cc}
\frac{1}{2} B_{1} & -\frac{1}{2} \sqrt{4 B_{2}-B_{1}^{2}} \\
\frac{1}{2} \sqrt{4 B_{2}-B_{1}^{2}} & \frac{1}{2} B_{1}
\end{array}\right)\binom{\bar{x}(t)}{\bar{y}(t)}+\binom{P_{1}(\bar{x}, \bar{y})}{P_{2}(\bar{x}, \bar{y})},
$$

where

$$
\begin{aligned}
& P_{1}(\bar{x}, \bar{y})=-\beta t_{21} \bar{x}^{2}-\beta t_{22} \bar{x} \bar{y} \\
& P_{2}(\bar{x}, \bar{y})=\frac{t_{21}+1}{t_{22}}\left(\beta t_{21} \bar{x}^{2}+\beta t_{22} \bar{x} \bar{y}\right) .
\end{aligned}
$$

From Theorem 3.5.3 of [31] we see that the existence of Neimark-Sacker bifurcation can be determined by the quantity $\theta$, where

$$
\theta=-\operatorname{Re}\left[\frac{(1-2 \lambda) \bar{\lambda}^{2}}{1-\lambda} l_{11} l_{20}\right]-\frac{1}{2}\left|l_{11}\right|^{2}-\left|l_{02}\right|^{2}+\operatorname{Re}\left(\bar{\lambda} l_{21}\right),
$$

and

$$
\begin{aligned}
& l_{11}=\frac{1}{4}\left[P_{1 \bar{x} \bar{x}}+P_{1 \bar{y} \bar{y}}+i\left(P_{2 \bar{x} \bar{x}}+P_{2 \bar{y} \bar{y}}\right)\right], \\
& l_{20}=\frac{1}{8}\left[\left(P_{1 \bar{x} \bar{x}}-P_{1 \bar{y} \bar{y}}+2 P_{2 \bar{x} \bar{y}}\right)+i\left(P_{2 \bar{x} \bar{x}}-P_{2 \bar{y} \bar{y}}-2 P_{1 \bar{x} \bar{y}}\right)\right], \\
& l_{02}=\frac{1}{8}\left[\left(P_{1 \bar{x} \bar{x}}-P_{1 \bar{y} \bar{y}}-2 P_{2 \bar{x} \bar{y}}\right)+i\left(P_{2 \bar{x} \bar{x}}-P_{2 \bar{y} \bar{y}}+2 P_{1 \bar{x} \bar{y}}\right)\right], \\
& l_{21}=\frac{1}{16}\left[\left(P_{1 \bar{x} \bar{x} \bar{x}}+P_{1 \bar{x} \bar{y} \bar{y}}+P_{2 \bar{x} \bar{x} \bar{y}}+P_{2 \bar{y} \bar{y} \bar{y}}\right)+i\left(P_{2 \bar{x} \bar{x} \bar{x}}+P_{2 \bar{x} \bar{y} \bar{y}}-P_{1 \bar{x} \bar{x} \bar{y}}-P_{1 \bar{y} \bar{y} \bar{y}}\right)\right] .
\end{aligned}
$$

By a straightforward and tedious calculation we obtain that

$$
\theta=\theta_{1}+\theta_{2}+\theta_{3}+\theta_{4}<0
$$

with

$$
\begin{aligned}
\theta_{1}= & \frac{B_{1}^{2}-2 B_{2}+B_{1} B_{2}-B_{1}^{3}-4 B_{2}^{2}+2 B_{1}^{2} B_{2}}{2\left(1-B_{1}+B_{2}\right)} \times \frac{\beta^{2} t_{21}^{2}}{8 t_{21}^{2}} \times\left(t_{21}+1-t_{22}\right)\left(t_{21}+1+t_{22}\right) \\
& +\frac{B_{1}-B_{2}-B_{1}^{2}+2 B_{1} B_{2}}{2\left(1-B_{1}+B_{2}\right)} \times \frac{\beta^{2} t_{21}}{8 t_{22}} \times \frac{1-t_{21}^{2}-t_{22}^{2}}{t_{22}^{2}}<0, \\
\theta_{2}= & -\frac{1}{8} \beta^{2} t_{21}^{2}\left(1+\frac{\left(1+t_{21}\right)^{2}}{t_{22}^{2}}\right)<0 \\
\theta_{3}= & -\frac{1}{16} \beta^{2}\left(\left(1+2 t_{21}\right)^{2}+\left(\frac{t_{21}\left(1+t_{21}\right)}{t_{22}}-t_{22}\right)^{2}\right)<0, \\
\theta_{4}= & 0
\end{aligned}
$$

Using the Neimark-Sacker bifurcation theorem in [31], we obtain that there exists a Neimark-Sacker bifurcation when $R_{0}$ passes through $R_{0}^{*}$, which competes the proof.


Figure 1 The stability analysis of disease-free equilibrium $E_{0}$ when $R_{0}>1$, where $\Lambda=0.1, d=0.006, r=0.1$, and $m=0.00001$

Figure 2 The existence of the endemic equilibrium $E_{1}, E_{2}$, and $E_{3}$, where $\Lambda=200, d=0.07, r=0.05$, and $m=0.1$


## 3 Numerical simulation

To provide some numerical evidence for the qualitative dynamic behavior of the system (1.6), we will display some bifurcation diagrams to illustrate the analytical results and explore a new dynamics behavior as the parameters change.

We first show the stability of a disease-free equilibrium $E_{0}$ when $R_{0}>1$. Letting $\Lambda=$ $0.1, d=0.006, r=0.1, m=0.00001$, and $\beta=0.0064$, we have $R_{0}=1.0063>1$ and $1+A+$ $2 \sqrt{A}=1.0202$. Accordingly, we have $1<R_{0}<1+A+2 \sqrt{A}$. In this case, Fig. 1(a) shows that the solutions of system (1.6) with different initial values tend to $E_{0}$ as $t$ tends to infinity. Keeping the other parameters the same as in Fig. 1(a) except for $\beta=0.01$, we have $R_{0}=$ $1.0535>1$ and $1+A+2 \sqrt{A}=1.0206$, that is, $R_{0}>1+A+2 \sqrt{A}$. In this case, Fig. 1(b) shows that the solutions of system (1.6) also tend to $E_{0}$ as $t$ tends to infinity. The numerical simulations illustrate that the disease-free equilibrium $E_{0}$ is stable when $R_{0}>1$.
Secondly, we show some examples about the existence of the endemic equilibrium $E_{i}$, $i=1,2,3$, Let $\Lambda=200, d=0.07, r=0.05, m=0.1$. Taking $\beta$ as the variable parameter, we see in Fig. 2 the existence of the endemic equilibrium $E_{i}, i=1,2,3$. Then we display the saddle-node bifurcation around endemic equilibrium $E_{1}$. Leaving the other parameters as in Fig. 2 except for $m=0.00001$ and $\beta=0.0001$, we have $R_{0}=2.38095$ and $1+A+$ $2 \sqrt{A}=1.00069$. Accordingly, we have $R_{0}>1+A+2 \sqrt{A}, E_{1}=(2856.157389,0.24152), E_{2}=$ (2857.1426, 0.00006), and $E_{3}=(1200.000103,966.66652)$. Figure 3 shows that the solutions


Figure 3 The saddle-node bifurcation of system (1.6) at $E_{1}$ when $R_{0}=1+A+2 \sqrt{A}$, where $\Lambda=200, d=0.07$, $r=0.05, m=0.00001$, and $\beta=0.0001$


Figure 4 The flip bifurcation of system (1.6) when $R_{0}=1+A+R_{2}$, where $\Lambda=200, d=0.007$, $\beta=0.0011308665, \gamma=0.1$, and $m=0.1$
of system (1.6) with different initial values tend to $E_{3}$ as $t$ tends to infinity, which implies that $E_{2}$ is unstable and $E_{3}$ is asymptotically stable, that is, $E_{1}$ is the saddle-node.

Thirdly, flip bifurcation is shown in Fig. 4. We take $\Lambda=200, d=0.007, \beta=0.0011308665$, $\gamma=0.1$, and $m=0.1$; then flip bifurcation occurs at endemic equilibrium $E_{3}$.
Then we chose $\Lambda=200, d=0.00001, \beta=0.0050003, \gamma=0.99999$, and $m=0.01$; then the endemic equilibrium $E_{3}$ is unstable, and Neimark-Sacker bifurcation happens at $E_{3}$ (see Fig. 5).

Finally, we chose $\Lambda=4, d=0.007$, and $\gamma=0.01$ to display the bifurcation diagrams at $E_{2}$ and $E_{3}$ in Fig. 6(a) and Fig. 6(b), respectively. Here $m$ and $\beta$ are the bifurcation parameters.

In addition, the numerical simulations show that the model possesses more complicated dynamics. For example, if we take $\Lambda=10, d=0.007, m=0.1$, and $\gamma=0.5115$, then Fig. 7 shows that system (1.6) undergoes a process from periodic doubling to chaos. As $\beta$ increases, firstly, the positive equilibrium loses its stability, and a stable two-cycle period appears, then the two-cycle period looses its stability, and a stable four-cycle period appears, and, finally, the period doubling process continues to chaos (see Fig. 7).


Figure 5 The Neimark-Sacker bifurcation of system (1.6) when $R_{0}=1+A+R_{1}$, where $\Lambda=200, d=0.00001$, $\beta=0.0050003, \gamma=0.99999$, and $m=0.01$

(a) the bifurcation diagram at $E_{2}$

(b) the bifurcation diagram at $E_{3}$

Figure 6 The bifurcation diagram of system (1.6) at the endemic equilibrium, where $\Lambda=4, d=0.007$, and $\gamma=0.01$


Figure 7 The chaos behavior of system (1.6), where $\Lambda=10, d=0.007, m=0.1$, and $\gamma=0.5115$

## 4 Discussion and conclusion

The discrete epidemic model is more suitable to describe the spread of diseases since the epidemiological data are usually collected in discrete time units, such as daily, weekly, or monthly. In this paper, we have studied the dynamical behavior of the discrete SIR epidemic model (1.3) with bilinear incidence rate and constant recovery by analyzing the discrete model (1.6). We have obtained sufficient conditions for the existence of the diseasefree equilibrium and multiple positive equilibria, the saddle-node bifurcation, the flip bifurcation, the 1:1 resonance, and the Neimark-Sacker bifurcation. Theoretic analysis shows that the disease-free equilibrium is globally asymptotically stable when the basic reproduction number is less than unity, and the numerical simulations illustrate that it is still asymptotically stable when the number is greater than unity.
Model (1.3) is suitable to describe rotavirus, hand, foot, and mouth disease, influenza, and so on. These infectious diseases have the following common characteristics: (1) Most cases are mild; after isolation, some need take a symptomatic treatment, and some can recover without any treatment. (2) Only the tiny fractions of the cases have serious complications that need take treatment in hospital. Therefore, like rotavirus, hand, foot, and mouth disease, and influenza, the recovery form of the infective individuals consists of two parts, the spontaneous recovery and the recovery due to treatment. In this paper, we use $\gamma I$ to describe the spontaneous recovery of the infective individuals without any treatment, and use $h(I)$ to describe the recovery of the infective individuals with treatment. Naturally, the recovery function $\gamma I+h(I)$ that we use here is a good description of the recovery of these diseases.

The dynamics of the discrete SIR epidemic model with bilinear incidence rate and constant recovery is much more complicated. We have investigated the stability of the diseasefree and positive equilibria and the bifurcation of the model analytically or numerically. In fact, the dynamic behavior of (1.3) has interesting practical significance.
Firstly, $R_{0}=\frac{\Lambda \beta}{d(d+\gamma)}$ is just the reproduction number of system (1.3) in the absence of the removal rate of the infective individuals. So, it is not the classic threshold of the epidemic dynamical models, and it cannot determines whether the disease will die out or not. However, $1<R_{0}<1+A+2 \sqrt{A}$ is a significant interval, which characterizes the importance of treatment for the disease control.
Secondly, the dynamic behavior of the model is different under different conditions when $R_{0}>1+A+2 \sqrt{A}$. When $1+A+R_{1}<R_{0}<1+A+R_{2}$, the endemic equilibrium $E_{3}$ is locally asymptotically stable. When $R_{0}=1+A+R_{1}$ and $\frac{2 d}{\gamma}<R_{1}<\frac{4 d(1-d-\gamma)+2 d^{2}(d+\gamma)+4(1-d)(d+\gamma)}{d \gamma(d+\gamma)}$, the system will undergo a Neimark-Sacker bifurcation. When $R_{0}=1+A+R_{2}$, the system will undergo a flip bifurcation. Namely, in the case where the disease persists in population, the disease may tend to a certain equilibrium, or it may exist in the form of periodic shocks.

At last, the theoretical analysis show that in the case $R_{0}>1$ the disease-free equilibrium and endemic equilibrium not only coexist, but are both stable. This implies that the disease may still be extinct even if there is a stable endemic equilibrium, that is, if the initial state of the disease is within the stable domain of the disease free equilibrium, then the disease will disappear, whereas if the initial state of the disease is within the stable domain of endemic equilibrium, then the disease will persist.
There are still many challenging problems on the dynamics of the discrete system (1.3). The numerical simulations demonstrate that the disease-free equilibrium is stable when
$R_{0}>1$. However, how to prove mathematically strictly the stability of the disease-free equilibrium $E_{0}$ when $R_{0}>1$ is a meaningful work. In addition, we also have no effective method to prove the global stability of the disease-free equilibrium $E_{0}$ when $1<R_{0}<1+A+2 \sqrt{A}$. We think they may be related to the switching system theory and expect that some analytical results can be obtained on those problems in the future.

## Appendix

Firstly, we prove that $R_{1} \geq 2 \sqrt{A}$, that is, $\frac{-d \gamma+\sqrt{\Delta_{1}}}{2 d(1-d)(d+\gamma)(1-d-\gamma)} \geq 2 \sqrt{A}$ with

$$
\Delta_{1}=d^{2} \gamma^{2}+4 d(1-d)(d+\gamma)(1-d-\gamma)\left(A(d(1-d-\gamma)+(1-d)(d+\gamma))^{2}+d^{2}\right)
$$

To this end, we need verify the following inequality:

$$
-d \gamma+\sqrt{\Delta_{1}} \geq 4 d(1-d)(d+\gamma)(1-d-\gamma) \sqrt{A}
$$

In fact, we have $\left(-d \gamma+\sqrt{\Delta_{1}}\right)^{2}-16 d^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2} A \geq 0$. Namely,

$$
\begin{aligned}
d^{2} \gamma^{2} & -2 d \gamma \sqrt{\Delta_{1}}+\Delta_{1}-16 d^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2} A \\
= & 2 d^{2} \gamma^{2}-2 d \gamma \sqrt{\Delta_{1}}+4 d^{3}(1-d)(d+\gamma)(1-d-\gamma) \\
& -8 d^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2} A \\
& +4 d^{3}(1-d)(d+\gamma)(1-d-\gamma)^{3} A+4 d(1-d)^{3}(d+\gamma)^{3}(1-d-\gamma) A \\
= & 2 d^{2} \gamma^{2}-2 d \gamma \sqrt{\Delta_{1}}+4 d^{3}(1-d)(d+\gamma)(1-d-\gamma) \\
& +4 d(1-d)(d+\gamma)(1-d-\gamma) A \gamma^{2} \\
= & 2 d^{2} \gamma^{2}-2 d \gamma \sqrt{\Delta_{1}}+4 d(1-d)(d+\gamma)(1-d-\gamma)\left(A \gamma^{2}+d^{2}\right) .
\end{aligned}
$$

We further show that $d^{2} \gamma^{2}+2 d(1-d)(d+\gamma)(1-d-\gamma)\left(A \gamma^{2}+d^{2}\right) \geq d \gamma \sqrt{\Delta_{1}}$, that is,

$$
\begin{aligned}
\left(d^{2} \gamma^{2}\right. & \left.+2 d(1-d)(d+\gamma)(1-d-\gamma)\left(A \gamma^{2}+d^{2}\right)\right)^{2}-d^{2} \gamma^{2} \Delta_{1} \\
= & d^{4} \gamma^{4}+4 d^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2}\left(A \gamma^{2}+d^{2}\right)^{2} \\
& \quad+4 d^{3} \gamma^{2}(1-d)(d+\gamma)(1-d-\gamma)\left(A \gamma^{2}+d^{2}\right) \\
& \quad-d^{4} \gamma^{4}-4 d^{3} \gamma^{2}(1-d)(d+\gamma)(1-d-\gamma)\left(A(d(1-d-\gamma)+(1-d)(d+\gamma))^{2}+d^{2}\right) \\
= & 4 d^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2}\left(A \gamma^{2}+d^{2}\right)^{2} \\
& \quad+4 d^{3} \gamma^{2}(1-d)(d+\gamma)(1-d-\gamma) A\left[\gamma^{2}-(d(1-d-\gamma)+(1-d)(d+\gamma))^{2}\right] \\
= & 4 d^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2}\left(A \gamma^{2}+d^{2}\right)^{2} \\
& \quad-16 d^{4} \gamma^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2} A . \\
= & 4 d^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2}\left[\left(A \gamma^{2}+d^{2}\right)^{2}-4 d^{2} \gamma^{2} A\right] \\
= & 4 d^{2}(1-d)^{2}(d+\gamma)^{2}(1-d-\gamma)^{2}\left(A \gamma^{2}-d^{2}\right)^{2} \\
\geq & 0 .
\end{aligned}
$$

## A summary of these calculations shows that $R_{1} \geq 2 \sqrt{A}$.

## Acknowledgements

We would like to thank very much the referees for a careful review and valuable comments to this manuscript, which improved it greatly. The authors also would like to thank Professor Jianquan Li for his suggestions, which also helped us to improve this paper.

## Funding

This work is supported by National Natural Science Foundation of China (grant 11971281, 11671243), by Natural Science Basic Research Plan in Shaanxi Province of China grant 2019JM-081, and by Natural Science Foundation of Shaanxi Provincial Department of Education grant 18JK0092.

## Availability of data and materials

This paper focuses on theoretical analysis and does not involve experiments and data.

## Competing interests

We declare that we have no competing interests.

## Authors' contributions

HC is responsible for the model formulation and study planning. HW has done the calculations and the proof. XW is responsible for the simulation section. All authors have read and approved the final manuscript.

## Publisher's Note

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Received: 15 May 2019 Accepted: 9 January 2020 Published online: 29 January 2020

## References

1. Li, L., Sun, G., Jin, Z.: Bifurcation and chaos in an epidemic model with nonlinear incidence rates. Appl. Math. Comput. 216, 1226-1234 (2010)
2. May, R.M.: Biological population obeying difference equations: stable points, stable cycles, and chaos. J. Theor. Biol. 51, 511-524 (1975)
3. May, R.M.: Deterministic models with chaotic dynamics. Nature 256, 165-166 (1975)
4. May, R.M.: Simple mathematical models with very complicated dynamics. Nature 261, 459-467 (1976)
5. Hassell, M.P.: Density dependence in single-species populations. J. Anim. Ecol. 44, 283-289 (1975)
6. Zhou, Y.C., Ma, Z.E.: Global stability of a class of discrete age-structured SIS models with immigration. Math. Biosci. Eng. 6, 409-425 (2009)
7. Hu, Z., Teng, Z.D., Jiang, H.: Stability analysis in a class of discrete SIRS epidemic models. Nonlinear Anal., Real World Appl. 13(5), 2017-2033 (2012)
8. Cao, H., Zhou, Y.C.: The basic reproduction number of discrete SIR and SEIS models with periodic parameters. Discrete Contin. Dyn. Syst., Ser. B 18(1), 37-56 (2013)
9. Hu, Z., Teng, Z.D., Zhang, L.: Stability and bifurcation analysis in a discrete SIR epidemic model. Math. Comput. Simul. 97, 80-93 (2014)
10. Suryanto, A., Kusumawinahyu, W.: Dynamically consistent discrete epidemic model with modified saturated incidence rate. Comput. Appl. Math. 32(2), 373-383 (2013)
11. Enatsu, Y., Nakata, Y., Muroya, Y.: Global stability for a class of discrete SIR epidemic models. Math. Biosci. Eng. 7, 347-361 (2010)
12. Sekiguchi, M.: Permanence of a discrete SIRS epidemic model with time delays. Appl. Math. Lett. 23, 1280-1285 (2010)
13. Chen, Q., Teng, Z., Wang, L., Jiang, H.: The existence of codimension-two bifurcation in a discrete SIS epidemic model with standard incidence. Nonlinear Dyn. 71, 55-73 (2013)
14. Yi, N., Zhang, Q., Liu, P., Lin, Y.: Codimension-two bifurctions analysis and tracking control on a discrete epidemic model. J. Syst. Sci. Complex. 24, 1033-1056 (2011)
15. Castillo-Chavez, C., Yakubu, A.A.: Discrete-time SIS models with complex dynamics. Nonlinear Anal. 47, 4753-4762 (2001)
16. Castillo-Chavez, C., Yakubu, A.A.: Discrete-time SIS models with simple and complex population dynamics. In: Castillo-Chavez, C., Blower, S., van den Driessche, P., Kirschner, D., Yakubu, A.A. (eds.) Mathematical Approaches for Emerging and Reemerging Infectious Diseases: An Introduction, pp. 153-163. Springger, New York (2002)
17. Cao, H., Zhou, Y.C., Song, B.J.: Complex dynamics of discrete SEIS models with simple demography. Discrete Dyn. Nat. Soc. 2011, 1-21 (2011).
18. Allena, L.J.S., van den Driessche, P.: The basic reproduction number in some discrete-time epidemic models. J. Differ. Equ. Appl. 14, 1127-1147 (2008)
19. Bacaër, N.: Periodic matrix population models: growth rate, basic reproduction number, and entropy. Bull. Math. Biol. 71, 1781-1792 (2009)
20. Allena, L.J.S.: Some discrete-time SI, SIR, and SIS epidemic models. Math. Biosci. 124, 83-105 (1994)
21. Zhou, L., Wang, Y., Xiao, Y., Li, M.Y.: Global dynamics of a discrete age-structured SIR epidemic model with applications to measles vaccination strategies. Math. Biosci. 308, 27-37 (2019)
22. Salceanu, P.L., Smith, H.L.: Persistence in a discrete-time stage-structured fungal disease model. J. Biol. Dyn. 3, 271-285 (2009)
23. Salceanu, P.L., Smith, H.L.: Persistence in a discrete-time stage-structured epidemic model. J. Differ. Equ. Appl. 16, 73-103 (2010)
24. van den Driessche, P., Yakubu, A.A.: Disease extinction versus persistence in discrete-time epidemic models. Bull. Math. Biol. (2018). https://doi.org/10.1007/s11538-018-0426-2
25. Franke, J.E., Yakubu, A.A.: Disease-induced mortality in density-dependent discrete-time SIS epidemic models. J. Math. Biol. 57(6), 755-790 (2008)
26. Franke, J.E., Yakubu, A.A.: Discrete-time SIS epidemic model in a seasonal environment. SIAM J. Appl. Math. 66, 1563-1587 (2006)
27. Cao, H., Zhou, Y.C., Ma, Z.E.: Bifurcation analysis of a discrete SIS model with bilinear incidence depending on new infection. Math. Biosci. Eng. 5\&6, 1399-1417 (2013)
28. Abdelaziz, M.A.M., Ismail, A.I., Abdullah, F.A., et al.: Bifurcations and chaos in a discrete SI epidemic model with fractional order. Adv. Differ. Equ. 2018, 44, 1-19 (2018).
29. Wang, W.D., Ruan, S.G.: Bifurcations in an epidemic model with constant removal rate of the infectives. J. Math. Anal. Appl. 291, 775-793 (2004)
30. Elaydi, S.N.: An Introduction to Difference Equations, 2nd edn. Undergraduate Texts in Mathematics. Springer, New York (1999)
31. Kuznetsov, Y.A.: Elements of applied bifurcation theory. In: Antman, S.S., Marsden, J.E., Sirovich, L. (eds.) Applied Mathematical Sciences. Springer, New York (2004)
32. Wiggins, S.: Introduction to Applied Nonlinear Dynamical Systems and Chaos, 2nd edn. Springer, New York (2003)

## Submit your manuscript to a SpringerOpen ${ }^{\bullet}$ journal and benefit from:

- Convenient online submission
- Rigorous peer review
- Open access: articles freely available online
- High visibility within the field
- Retaining the copyright to your article

Submit your next manuscript at $>$ springeropen.com


[^0]:    © The Author(s) 2020. This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit http://creativecommons.org/licenses/by/4.0/.

